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## Modelling inter-yarn friction in woven fabric armour

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### SUMMARY

7 The effects of inter-yarn friction on the ballistic performance of woven fabric armour are investigated  
9 in this paper. Frictional sliding between yarns is implemented in a computational model of the fabric  
11 that takes the form of a network. Yarn crimp and its viscoelastic properties are taken into account.  
13 Ballistic experiments are performed to verify the predictions of the model. Parametric studies show  
15 that the ballistic response of woven fabric is very sensitive to yarn friction when the friction coefficient  
is low but insensitive beyond a certain level. The results also show that very high inter-yarn friction  
can lead to premature yarn rupture, thus reducing the ability of the fabric to absorb impact energy.  
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15 KEY WORDS: woven fabric armour; yarn friction; ballistic

### 1. INTRODUCTION

17 Advancements in fabric armour have been driven mainly by improvements in the strength and  
19 flexibility of the yarns from which they are made. While important, the dynamic response  
of fabric systems when they are subjected to ballistic impact is not only dependent on the  
response of individual yarns, but also on the way yarns interact with one another because they  
21 are woven together into a fabric system. In addition to materials selection, the design parameters  
of flexible fabric armour include weave architecture, weave density, surface treatment of the  
23 constituent yarns and yarn count—these affect the way yarns interact with one another.

Many studies on the ballistic resistance of fabric systems have been experimental in nature.  
25 Ballistic impact tests provide useful information such as relationships between energy absorbed  
by a fabric and impact energy, high-speed photographic sequences of fabric deformation and  
27 post impact analyses of fabric perforation mechanisms. The majority of the experiments are

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1 aimed at identifying the effects of key parameters governing the ballistic performance of a  
2 system [1–4].

3 Computational simulation is increasingly used to investigate the ballistic performance of high-  
4 strength fabric as computational cost continues to decrease with the availability of affordable  
5 high-performance computing systems. Simulations can yield important information not measur-  
6 able from experiments when the fabric system is modelled with sufficient detail and accuracy.  
7 Many models of woven fabric have been developed. Early models are generally very simplified  
8 representations of the woven fabric. For example, Taylor and Vinson [5] treated fabric under-  
9 going ballistic impact as a homogeneous conical shell under quasi-static loading. The model can  
10 be considered empirical in that many of the parameters involved require inputs from results of  
11 ballistic tests on fabric. Although such membrane models may lack certain details, membrane  
12 model representations of woven fabric continue to be used for ballistic impact simulations.  
13 Recently, Lim *et al.* [6] studied the ballistic response of fabric using *finite element analysis* by  
14 modelling the fabric with isotropic membrane elements. Their emphasis was on developing a  
15 constitutive relation to represent the viscoelastic behaviour of the material. Roylance and Wang  
16 [7] introduced a more structurally accurate model in the form of a mesh of flexible linear ele-  
17 ments pin-jointed at yarn crossover points, where the mass is lumped. Shim *et al.* [8] adopted  
18 this fabric model and incorporated yarn crimp and the material viscoelasticity to investigate  
19 the ballistic response of cross-woven aramid fabric. The importance of accounting for crimp  
20 and dynamic material response was shown through parametric studies and comparison with  
21 experimental data.

22 Recent efforts at simulating the ballistic response of woven fabric have continued in the  
23 direction of including more structural features unique to woven fabrics into fabric models  
24 and formulating more accurate material models. Johnson *et al.* [9] introduced a FE model  
25 using a combination of bar and shell elements. The bar elements modelled the structural  
26 response of warp and weft yarns while shell elements provided low-level shear stiffness and  
27 grid stability. Many physical parameters like crimp modulus and locking strain were included.  
28 Tabiei and Ivanov [10, 11] formulated reduced and fully integrated membrane shell elements  
29 for cross-woven fabric and tested the material model using a commercial finite element analysis  
30 code (LS-DYNA). The material model was derived by a homogenization technique to describe  
31 fabric micro-mechanical models via an effective stiffness matrix. The material model took  
32 into account certain special characteristics of a fabric, such as reorientation and locking of  
33 yarns during impact. Shockey *et al.* [12] created finite element models of individual yarns  
34 using 3-dimensional brick elements and assembled them to form a woven fabric. Each yarn  
35 was discretized into eight hexahedral elements over its cross section and 12 elements along  
36 a crimp wavelength. While the model could include many details of the woven architecture,  
37 it also increased the computation time significantly because of the large number of degrees  
38 of freedom (DOF). Duan *et al.* [13, 14] presented a similar FE model in LS-DYNA, which  
39 was able to account for yarn sliding and frictional effects between interwoven yarns. Their  
40 parametric study showed that friction between yarns facilitates energy dissipation within the  
41 fabric target.

42 While it is now possible to perform computational simulations of ballistic impact on fabric  
43 systems using models where the yarns are fully discretized into fine solid finite elements,  
44 such models still remain computationally expensive. The focus of this paper is to imple-  
45 ment sliding between yarns in the fabric models presented by Roylance and Wang [7] and  
Shim *et al.* [8] while maintaining computational efficiency. Yarn–yarn interactions at crossover

1 points are responsible for transmitting energy from yarns in contact with the projectile (primary  
 2 yarns) to yarns that are not (secondary yarns) [15]. The interactions also affect the way the  
 3 fabric is perforated by the projectiles. Various effects of yarn–yarn friction have been reported.  
 4 Prosser [16] analysed fabrics impacted by 0.22 calibre fragment simulating projectiles (FSP) and  
 5 found that if slippage could be reduced by just four yarns at the impact point, an almost 50%  
 6 increase in protection could be achieved. Bazhenov [17] reported that in multiple-ply fabric  
 7 systems, the ravelling of yarns at the unclamped edges is related to the energy transferred to  
 8 yarns in each fabric layer. It is believed that yarn slippage generated additional energy dissipa-  
 9 tion via friction associated with slippage at crossover points. Impact experiments [1, 2, 18]  
 10 have been designed to investigate the effects of yarn friction on fabric ballistic resistance by  
 11 using chemical treatment to vary the yarn surface roughness. It was found that fabrics with a  
 12 high friction coefficient dissipated larger amounts of energy. Lee *et al.* [19] performed static  
 13 penetration tests on dry fabric and on a low resin content fabric-reinforced composite and  
 14 showed that the composite can absorb more energy than the dry fabric specimen because yarn  
 15 sliding is reduced.

16 In this paper, the interwoven yarn structure of fabrics is physically modelled to allow for  
 17 yarn slippage between yarns with friction. Fabric models which take yarn slippage into account  
 18 are currently limited to FE models where the yarns are fully discretized into solid elements  
 19 such as those of Shockey *et al.* [12] and Duan *et al.* [14]. The proposed model is able to  
 20 include yarn slippage using significantly smaller number of degrees of freedom.

## 21 2. NUMERICAL MODEL OF WOVEN FABRIC

22 The plain weave is one of the most common weave patterns for armour-grade fabric. Each  
 23 yarn in the warp direction passes alternately under and over an orthogonal yarn in the weft  
 24 direction. In the proposed fabric model, each yarn in the fabric is modelled by a zigzag  
 25 chain of linear elements pin-jointed together at nodal points. Each node is assigned a lumped  
 26 mass to represent the areal density of the fabric. Figure 1 shows the fabric model before  
 27 it is impacted by the projectile. Initially, the nodes coincide with yarn crossover points. For  
 28 a balanced plain-woven fabric, the geometry of the fabric model is completely described by  
 29 two parameters—the distance  $L_0$  between neighbouring nodes on a yarn and the distance  $L_1$   
 30 between a pair of nodes at a crossover point.  $L_1$  can be determined by the fabric thickness  
 31 and  $L_0$  can be determined by the number of yarns per unit length along warp/weft direction.  
 32 These parameters are normally available from the fabric specifications or are easily measured.

33 The explicit computational simulation proceeds at incremental time steps and is adopted  
 34 from the work of Shim *et al.* [8]. Nodal positions and velocities are updated after each time  
 35 step increment using a central time difference integration scheme. The nodal velocity at time  
 36  $t + \Delta t$  is computed from,

$$37 \quad \mathbf{v}_{t+\Delta t} = \mathbf{v}_t + \frac{\Delta t}{m} \sum \mathbf{f} \quad (1)$$

38 where  $\sum \mathbf{f}$  is the resultant force acting on the node arising from tension in the two connected  
 39 yarn elements, and  $m$  is the mass of the node. The nodal position is then updated using

$$40 \quad \mathbf{x}_{t+\Delta t} = \mathbf{x}_t + \mathbf{v}_{t+\Delta t} \Delta t \quad (2)$$

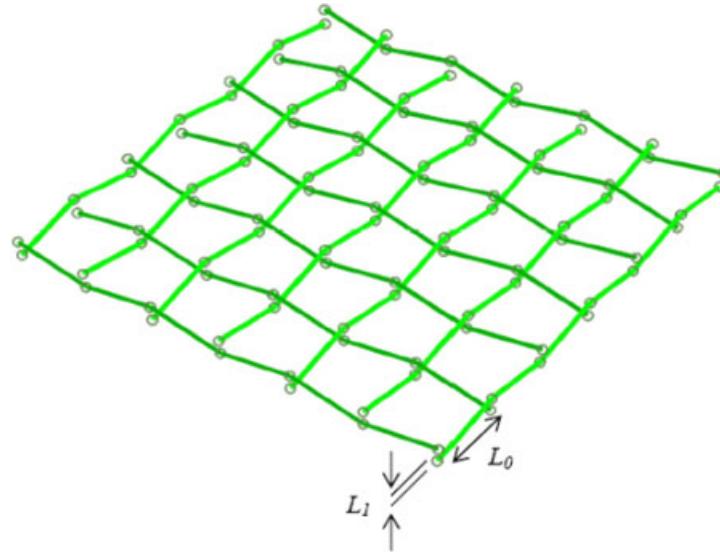


Figure 1. Network model for cross-woven fabric.

1 The force  $\mathbf{f}$  is computed via a viscoelastic constitutive equation, written generically as

$$\mathbf{f}_{t+\Delta t} = \varphi(\boldsymbol{\sigma}_t, \boldsymbol{\varepsilon}_t, \boldsymbol{\varepsilon}_{t+\Delta t}) \quad (3)$$

3 where the stress,  $\boldsymbol{\sigma}$ , and strain,  $\boldsymbol{\varepsilon}$ , are calculated for each linear yarn element.

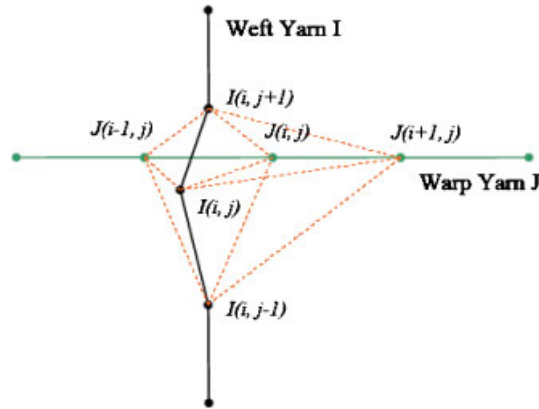
5 Equations (1)–(3) are applied to all the nodes first without consideration of possible interpenetration of nodes into the surface of the projectile or interpenetration of yarn elements. The nodal co-ordinates and velocities are then adjusted to exclude interpenetration in the subsequent step through two contact algorithms—one for yarn-to-yarn contact and another for fabric-to-projectile contact.

### 9 3. ALGORITHM FOR CONTACT BETWEEN YARN ELEMENTS

11 At each time step, a check is made to determine if any yarn element has interpenetrated another element. Instead of implementing a general contact algorithm, it was observed from post-impact specimens of cross-woven fabric that a simplified contact algorithm tailored for yarn-to-yarn contact would suffice because,

- 13 • any pair of warp and weft yarns are in contact with each other at no more than one
- 15 location, except for the case of ravelled yarns along the fabric edges, and
- 17 • weft yarns are prevented from contacting adjacent weft yarns by the interwoven warp yarns and *vice versa* because armour-grade fabrics are tightly woven.

19 After the nodal positions are updated by Equation (2), the nearest points between each pair of warp and weft yarns are determined. Their relative position will dictate if interpenetration has occurred. The nearest nodes between each warp and weft pair are first identified. A small



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Figure 2. Orthogonal yarn elements (dashed lines indicate probable pairs of nearest nodes).

1 time step size for the explicit computation means that the configuration of the fabric model  
 2 changes only incrementally between time steps. The maximum time step increment  $\Delta t$  required  
 3 for numerical stability is described by Shim *et al.* [8]— $\Delta t$  must be smaller than the time for  
 stress wave to propagate through one single linear element, i.e.

$$5 \quad \Delta t \leq \alpha \frac{L}{c} \quad (4)$$

7 where  $L$  is the length of one linear segment in the model,  $c = \sqrt{E/\rho}$  (the wave speed in a  
 single yarn) and  $\alpha$  a positive constant less than unity.

9 The use of a small time step means that the nearest pair of nodes can change by at most  
 only one node along the yarns between time steps. Hence, if nodes  $I(i, j)$  and  $J(i, j)$  in  
 Figure 2 are the nearest pair in the previous time step for weft yarn  $I$  and warp yarn  $J$ , it  
 11 is sufficient to compare the distances between nodes  $I(i, j - 1)$ ,  $I(i, j)$ ,  $I(i, j + 1)$  and nodes  
 $J(i - 1, j)$ ,  $J(i, j)$ ,  $J(i + 1, j)$  to determine the nearest nodes for the current time step.

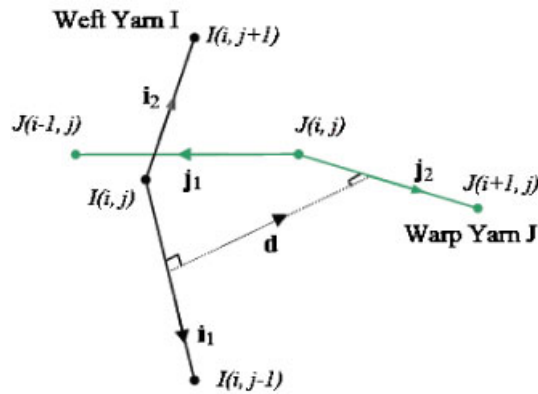
13 Assuming nodes  $I(i, j)$  and  $J(i, j)$  are the closest pair of nodes, as shown in Figure 3,  
 then the point on yarn  $I$  closest to yarn  $J$  is assumed to lie between nodes  $I(i, j - 1)$  and  
 15  $I(i, j + 1)$ . Similarly, the point on yarn  $J$  closest to yarn  $I$  is assumed to lie between nodes  
 $J(i - 1, j)$  and  $J(i + 1, j)$ . The four yarn segments in Figure 3 are denoted by vectors  $\mathbf{i}_\alpha$  and  
 17  $\mathbf{j}_\beta$  ( $\alpha, \beta = 1, 2$ ) as shown. The relative position of a point on a yarn segment of yarn  $J$  from  
 a point on a yarn segment of yarn  $I$  is given by

$$19 \quad \mathbf{d} = \vec{IJ} - a\mathbf{i}_\alpha + b\mathbf{j}_\beta \quad (5)$$

20 where  $\vec{IJ}$  is the vector from node  $I(i, j)$  to node  $J(i, j)$  and ( $0 \leq a, b \leq 1$ ).

21 The shortest distance between the two yarns is given by the minima of  $|\mathbf{d}|$  with respect to  $a$   
 and  $b$ , i.e.

$$23 \quad \begin{bmatrix} \partial|\mathbf{d}|/\partial a \\ \partial|\mathbf{d}|/\partial b \end{bmatrix} = \mathbf{0} \quad (6)$$



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Figure 3. Definition of contact vector ( $\mathbf{d}$ ) and yarn vectors ( $\mathbf{i}, \mathbf{j}$ ) for nearest node pair  $I(i, j)$  and  $J(i, j)$ .

1 Combining Equations (5) and (6) gives

$$\begin{bmatrix} -\mathbf{i}_\alpha \cdot (\vec{IJ} - a\mathbf{i}_\alpha + b\mathbf{j}_\beta) \\ \mathbf{j}_\beta \cdot (\vec{IJ} - a\mathbf{i}_\alpha + b\mathbf{j}_\beta) \end{bmatrix} = \mathbf{0} \quad (7)$$

3 Equation (7) can be expressed as

$$\begin{bmatrix} \mathbf{i}_\alpha \cdot \mathbf{i}_\alpha & -\mathbf{i}_\alpha \cdot \mathbf{j}_\beta \\ -\mathbf{i}_\alpha \cdot \mathbf{j}_\beta & \mathbf{j}_\beta \cdot \mathbf{j}_\beta \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} \vec{IJ} \cdot \mathbf{i}_\alpha \\ \vec{IJ} \cdot \mathbf{j}_\beta \end{bmatrix} \quad (8)$$

5 Equation (8) is solved for  $a$  and  $b$  to locate the points on the two yarn segments closest to  
 7 each other after which,  $\mathbf{d}$  can be determined from (5). The pair of yarn segments with the  
 lowest value of  $|\mathbf{d}|$  is then checked for interpenetration. It should be noted that  $\mathbf{d}$  is normal to  
 a plane parallel to  $\mathbf{i}_\alpha$  and  $\mathbf{j}_\beta$  and is referred to as the contact vector.

9 Figure 4 shows a warp yarn,  $I$ , going over and under adjacent weft yarns,  $J$  and  $J + 1$ .  
 Node numbers along all warp yarns increase from bottom to top and node numbers along weft  
 11 yarns increase from left to right. For every warp yarn element,  $\mathbf{i}^+$  is a vector along the yarn  
 element in the direction of increasing node number. A vector  $\mathbf{j}^+$  is similarly defined for weft  
 13 yarn elements. It can be seen from Figure 4 that a vector pointing from yarn  $I$  to yarn  $J$  at  
 their crossover point is given by  $\mathbf{n} = -\mathbf{i}^+ \times \mathbf{j}^+$ , where  $\mathbf{i}^+$  and  $\mathbf{j}^+$  are associated with the yarn  
 15 segments in contact. The crossover point may change as yarns  $I$  and  $J$  slide along one another,  
 however,  $-\mathbf{i}^+ \times \mathbf{j}^+$  will continue to indicate the relative position of yarn  $J$  with respect to  $I$ .  
 17 Hence, if contact vector  $\mathbf{d}$  points in the opposite direction to  $\mathbf{n}$ , it indicates that interpenetration  
 has occurred. Interpenetration of other weft–warp yarn pairs is detected in the same manner. It  
 19 should be noted that for some pairs of weft and warp yarns, the vector pointing from the warp  
 to the weft yarn is given by  $\mathbf{n} = \mathbf{i}^+ \times \mathbf{j}^+$  instead, as is the case for yarns  $I$  and  $J + 1$ . The  
 21 vector  $\mathbf{n}$  is determined for each pair of weft and warp yarn based on the weave pattern of the  
 fabric prior to the start of the computation. Each yarn segment is considered to be a cylindrical

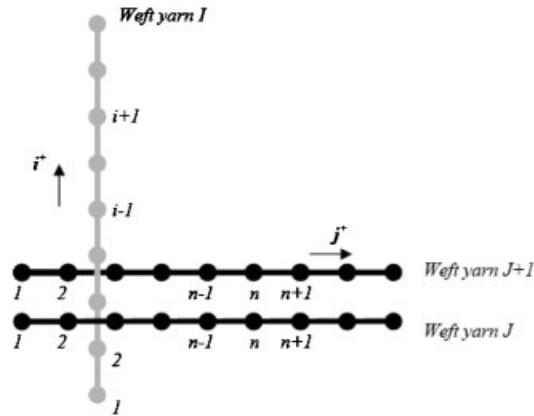


Figure 4. Node numbering for weft and warp yarns and definition of vectors  $\vec{i}^+$  and  $\vec{j}^+$ .

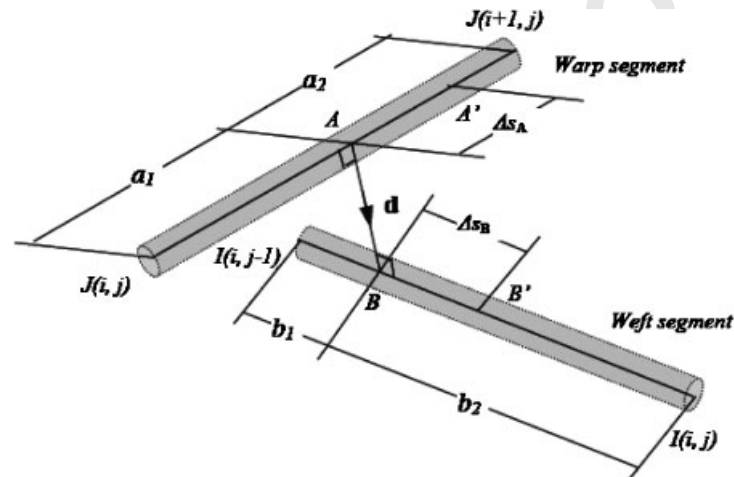


Figure 5. Illustration of interpenetrating yarn segments showing previous points of contact,  $A'$  and  $B'$ .

- 1 rod with a diameter  $D$  equivalent to half of the fabric thickness, hence, interpenetration also
- 2 occurs if  $|\mathbf{d}|$  is smaller than  $D$ .
- 3 The usual master-slave contact method to correct for interpenetration is not imposed here
- 4 because inter-yarn contact is symmetric. Instead, an adjustment is made to the nodal velocities
- 5 of interpenetrating yarn segments to move the yarn segments back toward the surface of each
- 6 other within a single time step. Figure 5 shows two interpenetrating segments before any
- 7 correction is made. Points  $A$  and  $B$  are the closest points on the segments and the magnitude
- 8 of the interpenetration is  $\delta = |\mathbf{d}| + D$ . Hence, an adjustment of  $\delta$  along vector  $\mathbf{d}$  is to be made
- 9 to the relative position of the yarn segments. Point  $A$  divides the warp yarn segment in the

1 ( $b_1 + b_2 = 1$ ). To move the segments by a total of  $\delta$ , the velocity of the nodes connected to  
 2 the contacting segments are adjusted by

$$\begin{aligned}
 & \Delta \mathbf{v}_{J(i,j)} = -a_2 \Delta \mathbf{v} \quad \text{and} \quad \Delta \mathbf{v}_{I(i,j-1)} = b_2 \Delta \mathbf{v} \\
 & \Delta \mathbf{v}_{J(i+1,j)} = -a_1 \Delta \mathbf{v} \quad \quad \quad \Delta \mathbf{v}_{I(i,j)} = b_1 \Delta \mathbf{v}
 \end{aligned} \tag{9}$$

3 where

$$|\Delta \mathbf{v}| = - \frac{\delta}{\Delta t (a_1^2 + a_2^2 + b_1^2 + b_2^2)} \tag{10}$$

4 and  $\Delta t$  is the time step increment for computation.

5 The velocity adjustments will result in a relative displacement of  $\delta$  within a time step and the  
 6 total linear and angular momentum of the yarn segments before and after the velocity correction  
 7 are conserved. The end result is similar to a normal contact force  $\mathbf{f}_n$  pushing points  $A$  and  $B$   
 8 back toward their yarn surfaces, where

$$\mathbf{f}_n = m \frac{\Delta \mathbf{v}}{\Delta t} \tag{11}$$

9 After adjustments to the nodal velocities to correct for yarn interpenetration as described, further  
 10 adjustments are made to account for frictional forces due to sliding of yarn elements over one  
 11 another during contact. The treatment of contact between rod elements formulated by Zavarise  
 12 and Wriggers [20] is adopted and modified for the current finite difference implementation.

13  $A'$  and  $B'$  in Figure 5 represent the points of contact between the two segments in the  
 14 previous time step. This means that the warp yarn  $J$  has slid along yarn  $I$  by a distance  
 15 of  $\Delta s_A$  within a time interval of  $\Delta t$ .

16 A trial friction force is first computed by a penalty formulation

$$\mathbf{f}_t^{\text{trial}} = K \Delta s_A \tag{12}$$

17  $K$  is a penalty parameter defined by  $K = \alpha(m/\Delta t^2)$ , where  $m$  is the nodal mass and  $\alpha$  is  
 18 a factor  $\alpha \leq 1.0$ . The trial friction force is applied to yarn  $J$  in the direction  $\overrightarrow{A'A}$ . When  
 19  $\alpha = 1$ ,  $\mathbf{f}_t^{\text{trial}}$  is equivalent to the force required to move the yarn segments back to the contact  
 20 configuration in the previous time step.

21 A frictional slippage function based on Coulomb friction is used to determine if sliding  
 22 occurs. The magnitude of the trial friction force is compared to the maximum Coulombic  
 23 friction force, i.e.

$$\varphi = |\mathbf{f}_t^{\text{trial}}| - \mu |\mathbf{f}_n| \tag{13}$$

24 where  $\mu$  is the friction coefficient and  $\mathbf{f}_n$  is obtained from Equation (11). If  $\varphi$  is negative, no  
 25 slippage occurs. The force applied to the  $J$  yarn element is  $\mathbf{f}_t = \mathbf{f}_t^{\text{trial}}$  in the direction  $\overrightarrow{A'A}$  and  
 26 the point  $A'$  is recorded as the point of contact on yarn  $J$ . A positive value of  $\varphi$  indicates  
 27 slippage. In this case, the force applied to the  $J$  yarn element is  $\mathbf{f}_t = \mu \mathbf{f}_n$  and the new contact  
 28 position on yarn  $J$  relative to  $A'$  is

$$\Delta s_A = \frac{|\mathbf{f}_t^{\text{trial}}| - \mu |\mathbf{f}_n|}{K} \tag{14}$$



1 Figure 5 also shows that the point of contact on yarn  $I$  moving from  $B'$  to  $B$  in the absence of  
 2 friction. When friction is present, the frictional force to be applied to yarn  $I$  is determined in  
 3 the same manner as for yarn  $J$ . It should be noted that the friction force  $\mathbf{f}_t$  applied to yarn  $J$   
 4 will also be applied to yarn  $I$  but in the opposite direction and *vice versa*, i.e. the actual  
 5 friction force  $\mathbf{f}_f$  acting at the contact point is the resultant of the two tangential forces along  
 6 the beam axes. The velocity changes for each node arising from friction can be calculated  
 7 from

$$\begin{aligned} \Delta \mathbf{v}_{J(i,j)} &= -a_2 \frac{\mathbf{f}_f \Delta t}{m} & \Delta \mathbf{v}_{I(i,j-1)} &= b_2 \frac{\mathbf{f}_f \Delta t}{m} \\ & \text{and} & & \\ \Delta \mathbf{v}_{J(i+1,j)} &= -a_1 \frac{\mathbf{f}_f \Delta t}{m} & \Delta \mathbf{v}_{I(i,j)} &= b_1 \frac{\mathbf{f}_f \Delta t}{m} \end{aligned} \quad (15)$$

9 Although the flexural rigidity of yarns is minimal, a small amount of bending stiffness between  
 10 interconnecting yarn segments is necessary to prevent unrestricted rotation of the pin-jointed  
 11 yarn elements. For this purpose, a compressive 'spring' equivalent to 5% the stiffness of the  
 12 yarn elements connects the middle of adjoining yarn elements.

13 It is possible to improve the accuracy of the yarn-to-yarn contact algorithm by performing  
 14 multiple iterations of the procedures described to Equations (5)–(15) in a single time step.  
 15 Johnson and Stryk [21] introduced this technique to obtain better compatibility between the  
 16 surfaces of master and slave elements for contact and sliding interface problems. Within one  
 17 time step,  $N$  iterations were carried out to gradually correct surface interpenetration and include  
 18 sliding frictional forces. A fraction,  $\alpha$ , of the required position and velocity changes is made  
 19 within one iteration, where  $\alpha$  is determined from

$$\alpha = \frac{1}{\sqrt{N - n - 1}} \quad (16)$$

21 where  $N$  is the specified total number of iterations and  $n$  is the current iteration number. Within  
 22 each iteration, the algorithm updates the contact information after making fine adjustments to  
 23 the nodal positions. For this iterative procedure Equation (10) is modified to

$$|\Delta \mathbf{v}| = -\alpha \frac{\delta}{\Delta t (a_1^2 + a_2^2 + b_1^2 + b_2^2)} \quad (17)$$

#### 25 4. ALGORITHM FOR SLIDING CONTACT WITH FRICTION BETWEEN 26 FABRIC AND PROJECTILE

27 In the computational simulation, the projectile is modelled as a rigid sphere striking a rectangular  
 28 fabric target normally at the centre. Hence, only a quarter of the fabric needs to be modelled  
 29 because of the double symmetry, and rotation of the projectile need not be considered.

30 The effect of friction between yarn elements and the projectile is accounted for in the  
 31 simulation. As with yarn-to-yarn interactions, the position of the projectile and the fabric  
 32 nodes are initially updated without considering contact between the projectile and fabric. After  
 33 the initial position update, some fabric nodes may have interpenetrated the projectile surface.

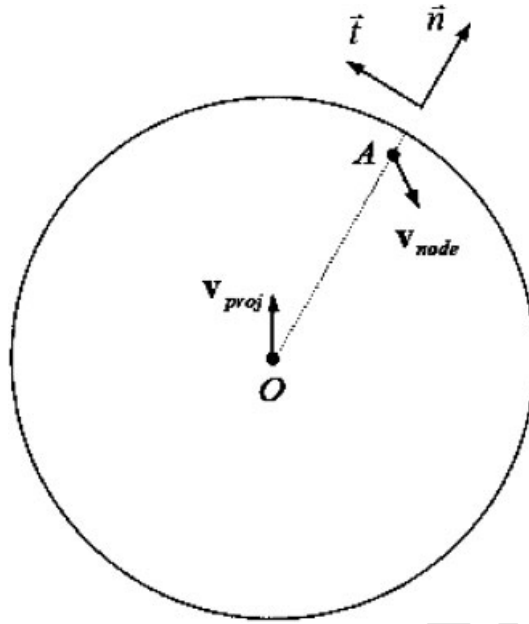


Figure 6. Interpenetration of yarn node A into spherical projectile.

1 Interpenetration occurs when

$$D_{\text{node}} = |\mathbf{x}_{\text{node}} - \mathbf{x}_{\text{proj}}| \leq R_{\text{proj}} \quad (18)$$

3 Here,  $R_{\text{proj}}$  is the radius of the projectile,  $\mathbf{x}_{\text{node}}$  is the position vector of the node and  $\mathbf{x}_{\text{proj}}$  is  
 5 the position vector of the projectile centre  $O$ . Figure 6 shows the situation where a node,  $A$ ,  
 7 has interpenetrated the projectile. Node  $A$  is treated as a slave node and a force,  $\mathbf{f}_n^p$ , normal  
 to the projectile surface (i.e. in the radial direction  $\mathbf{n}$ , of the spherical projectile) is applied to the  
 node to push it back onto the projectile surface within one time step. This force is determined  
 from

$$9 \quad \mathbf{f}_n^p = \frac{m(R_{\text{proj}} - D_{\text{node}})}{\Delta t^2} \mathbf{n} \quad (19)$$

In the presence of friction, the frictional force acting on the node,  $\mathbf{f}_t^p$ , also needs to be found.  
 11 The relative velocity of the node to the projectile is first determined, i.e.

$$\mathbf{v}_R = \mathbf{v}_{\text{node}} - \mathbf{v}_{\text{proj}} \quad (20)$$

13 The component of  $\mathbf{v}_R$  normal to  $\mathbf{n}$  determines the direction of the frictional force and is  
 given by

$$15 \quad \mathbf{v}_R^t = \mathbf{v}_R - (\mathbf{v}_R \cdot \mathbf{n})\mathbf{n} \quad (21)$$

1 The tangential force applied to the node is then

$$\mathbf{f}_t^p = -\min\left(m \frac{\mathbf{v}_R^t}{\Delta t}, \mu_p |\mathbf{f}_p^n| \mathbf{t}\right) \quad (22)$$

3 where  $\mu_p$  is the Coulomb friction coefficient between the fabric and projectile and  $\mathbf{t}$  is the unit vector parallel to  $\mathbf{v}_R^t$ .

## 5. COMPUTATIONAL SIMULATION OF WOVEN ARAMID FABRIC

Actual ballistic tests were performed by launching 12 mm (7 g) spherical steel projectiles to strike  $120 \times 120$  mm Twaron<sup>®</sup> CT 716 fabric specimens clamped along two opposite edges, with normal incidence. Twaron<sup>®</sup> is made from aramid (PPTA) fibrous material, similar to that of Kevlar<sup>®</sup>. The algorithms described were implemented to simulate the ballistic experiments and the relevant specifications of the fabric are listed in Table I. From the specifications, the parameters to construct the woven fabric network structure shown in Figure 1 were ascertained to be  $L_1 = 0.2$  mm and  $L_0 = 0.844$  mm, and the nodal mass was calculated to assume the value  $m = 9.4 \times 10^{-5}$  g.

The material constitutive relationship for yarn elements in the fabric model is represented by a linear spring–dashpot Zener viscoelastic model to describe the strain rate sensitivity of the polymer fibre [6]. The constitutive relationship is described by

$$\left(1 + \frac{K_2}{K_1}\right) \sigma + \frac{\mu_2}{K_1} \dot{\sigma} = K_2 \varepsilon + \mu_2 \dot{\varepsilon} \quad (23)$$

where  $\sigma$ ,  $\varepsilon$  and  $\dot{\varepsilon}$  are the stress, strain and strain rate, respectively. The constants defining the springs ( $K_1$ ,  $K_2$ ) and dashpot ( $\mu$ ) are obtained semi-empirically. At a constant strain rate, i.e.  $\dot{\varepsilon}(t) = d\varepsilon(t)/dt = \dot{\varepsilon}_0$ , the stress as a function of strain and strain rate, with initial conditions  $\varepsilon = 0$  and  $\sigma = 0$ , corresponding to Equation (23) are

$$\sigma = \frac{K_1 K_2}{K_1 + K_2} \varepsilon - \frac{K_1^2 \mu}{(K_1 + K_2)^2} \dot{\varepsilon}_0 (e^{-((K_1 + K_2)/\mu)\varepsilon/\dot{\varepsilon}_0} - 1) \quad (24)$$

23 and

$$\frac{d\sigma}{d\varepsilon} = \frac{K_1 K_2}{K_1 + K_2} + \frac{K_1^2}{K_1 + K_2} (e^{-((K_1 + K_2)/\mu)\varepsilon/\dot{\varepsilon}_0}) \quad (25)$$

Experimental data for Twaron yarns at different strain rates is scarce. Here, use of data reported by Gu [22] as shown in Table II, is employed. The values of the parameters for the viscoelastic material model are obtained from the best fit of Equation (25) to the reported Young's modulus, as shown in Figure 7. These values are  $K_1 = 7.28 \times 10^{10}$  Pa,  $K_2 = 4.17 \times 10^{11}$  Pa and  $\mu = 6.26 \times 10^8$  Pa s. The failure strain of yarn elements is approximated to be within the range of 5.2–5.7%.

The numerical model is implemented through an in-house code and simulations are performed on a 2.8 GHz Pentium IV personal computer. The time taken to complete a simulation depends on the assigned impact velocity of the projectile. For low impact velocities, the penetration

Table I. Twaron<sup>®</sup> fabric CT716 specification.

|                              |                        |
|------------------------------|------------------------|
| Specific density             | 1.44 g/cm <sup>3</sup> |
| Linear density warp and weft | 1100 dtex f 1000       |
| Areal density                | 280 g/m <sup>2</sup>   |
| Thickness                    | 0.40 mm                |
| Sett (per 10 cm)             | 122                    |

Table II. Twaron<sup>®</sup> fibre bundle properties [22].

| Strain rate (s <sup>-1</sup> ) | <i>E</i> (GPa) | $\epsilon_{\max}$ (%) |
|--------------------------------|----------------|-----------------------|
| 10 <sup>-2</sup>               | 62             | 5.19                  |
| 180                            | 69             | 5.22                  |
| 480                            | 70             | 5.47                  |
| 1000                           | 72             | 5.70                  |

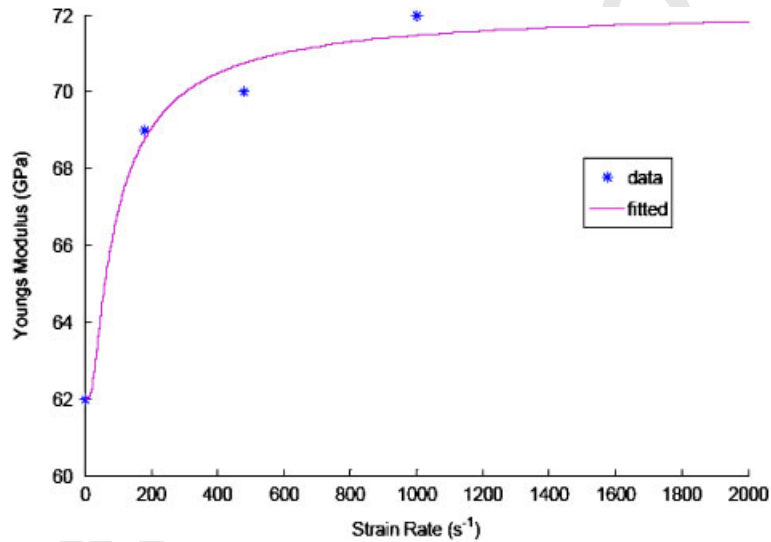


Figure 7. Variation of Young's modulus with strain rate based on Equation (24) and fit with experimental data from Gu [22].

1 process is slow and the simulations take about 1–2 h of computational time whereas high-  
 2 velocity impacts require only a few minutes.

### 3 6. VALIDATION OF FABRIC MODEL

4 The effects of yarn-to-yarn friction were investigated in this paper by performing simulations  
 5 over a range of friction coefficients. Briscoe and Motamedi [18] measured the friction coefficient

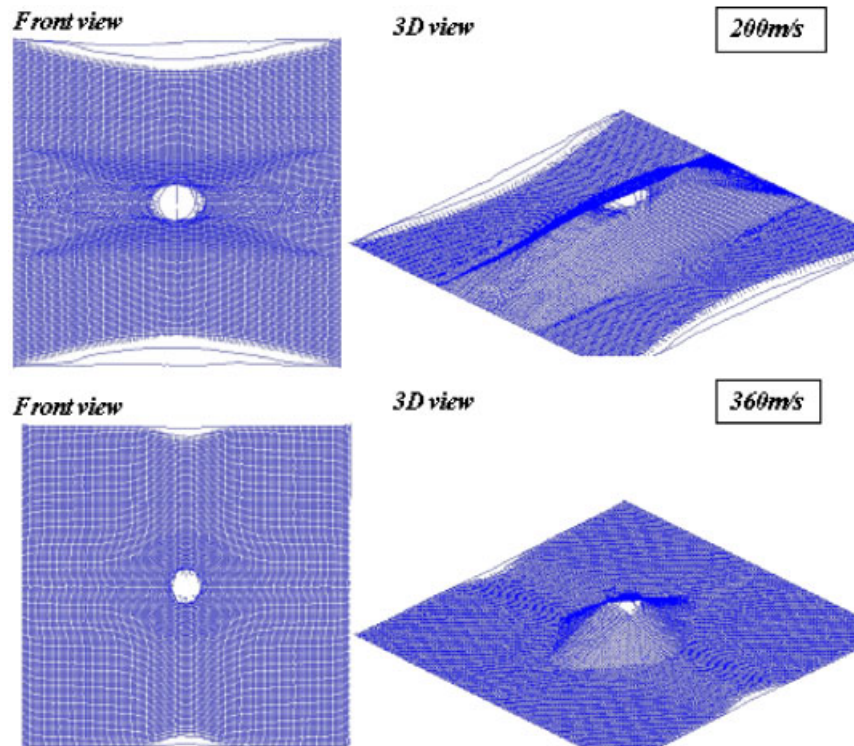


Figure 8. Predicted fabric perforation characteristics for impact at 200 and 360 m/s.

1 of the aramid yarn, Kevlar<sup>®</sup> 49, by using a hanging-yarn experimental method and reported  
 2 values of  $\mu = 0.22 \pm 0.03$ . Rebouillat [23] studied the effects of relative sliding speed on the  
 3 friction coefficient between Kevlar<sup>®</sup> yarns using a classical pin-on-disc tribometer. The yarn  
 4 coefficient of friction was reported to be  $\mu = 0.22 \pm 0.03$  at a sliding speed of 9.6 mm/min and  
 5  $\mu = 0.27 \pm 0.03$  at a speed of  $77 \times 10^3$  mm/min.

6 Figure 8 shows the deformation predicted by the fabric model with  $\mu = 0.2$ , after perforation  
 7 by a projectile striking at 200 and 360 m/s, respectively. The configurations predicted by  
 8 the computational model correlate closely with experimental observations. Deformation at low  
 9 impact velocity differs from that at high impact velocity. For an impact velocity of 200 m/s,  
 10 the fabric was fully stretched prior to perforation and a large area of fabric was deflected.  
 11 At 360 m/s, the transverse waves did not have enough time to propagate to the edges before  
 12 the perforation occurred, resulting in a small area of deflection. Similar observations have been  
 13 captured by high-speed photographs of ballistic tests, as shown in Figure 9.

14 The simulations also identified differences in the way the yarns break at the region of  
 15 impact, as shown in Figure 8. At a relatively low velocity of 200 m/s, the number of broken  
 16 warp and weft yarns differs significantly. Clamped yarns broke earlier than unclamped yarns  
 17 because unclamped yarns do not elongate significantly during impact. As the fabric deflects,  
 the unclamped yarns are simply pulled inwards whereas clamped yarns need to stretch to

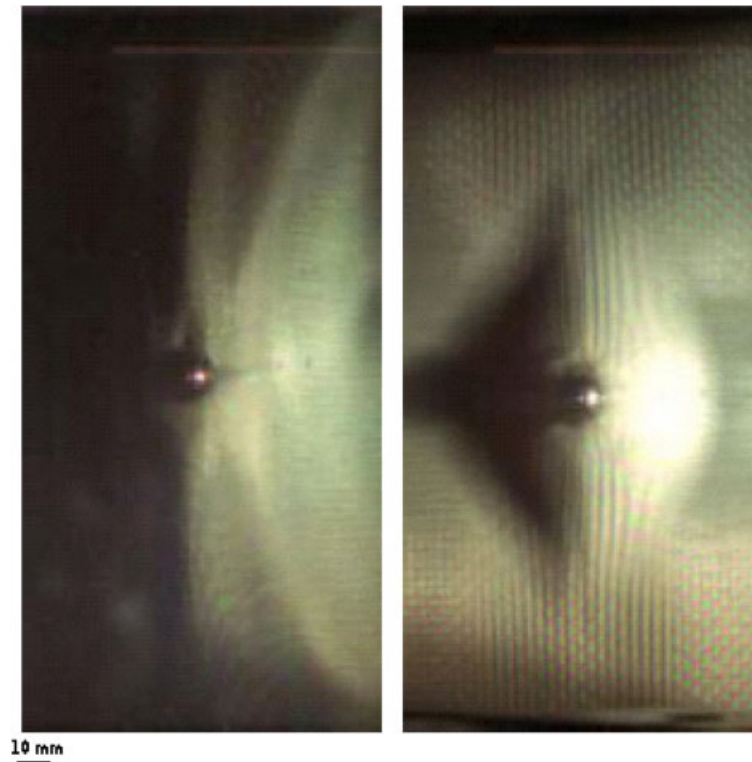


Figure 9. High-speed photographs of fabric perforation; left: low-velocity impact 200 m/s; right: high-velocity impact 360 m/s.

1 accommodate the fabric deflection. As a consequence, the primary clamped yarns rupture to  
 3 form a slit, while most of the unclamped yarns are able to slide against the projectile without  
 5 breaking. Only a few unclamped yarns remain in contact with the projectile and are pulled  
 7 to breakage. At a high impact velocity like 400 m/s, the fabric is perforated within a very  
 9 short time. Boundary conditions have less influence at high impact velocities because the fabric  
 is perforated before the primary yarns are uniformly stretched. Hence, the numbers of yarns  
 broken in the weft and warp directions are almost equal giving rise to square rather than  
 slit-like perforations. These features are also observed in actual fabric specimens, as illustrated  
 in Figure 10.

11 Another physical phenomenon that was replicated in the simulation is the ravelling of yarns  
 along the free edges. This ravelling gives an indication of the amount of energy that is  
 transmitted to the edges of the fabric.

13 The energy absorbed by a fabric during impact is one measure of its ballistic resistance  
 to small projectiles. As indicated in Table II, the failure strain varies with strain rate [22].  
 15 Computational simulations were performed using failure strains at the lower and upper bounds  
 of 5.2 and 5.7% for comparison. Simulation results of energy absorbed by the fabric over an  
 17 impact velocity range of 180–500 m/s are plotted together with experimental data in Figure 11.



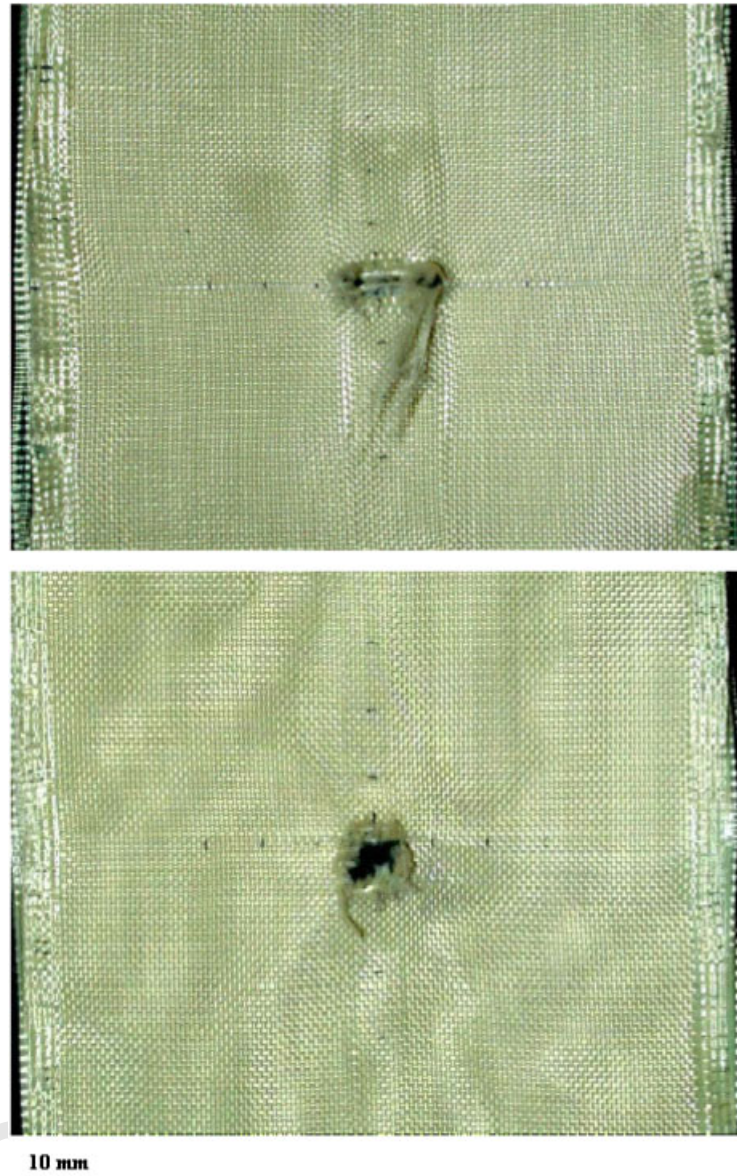


Figure 10. Single fabric ply after penetration, left: low-velocity impact 200 m/s, right: high-velocity impact 360 m/s.

- 1 The graphs show reasonable agreement between simulation and experiments. Both numerical and simulation results display a sudden decrease in energy absorption at about 250 m/s. It should
- 3 be noted that the material properties were based on experimental results from Gu [22] and not from tests on the Twaron<sup>®</sup> fabric used in the ballistic tests reported. Hence, only qualitative

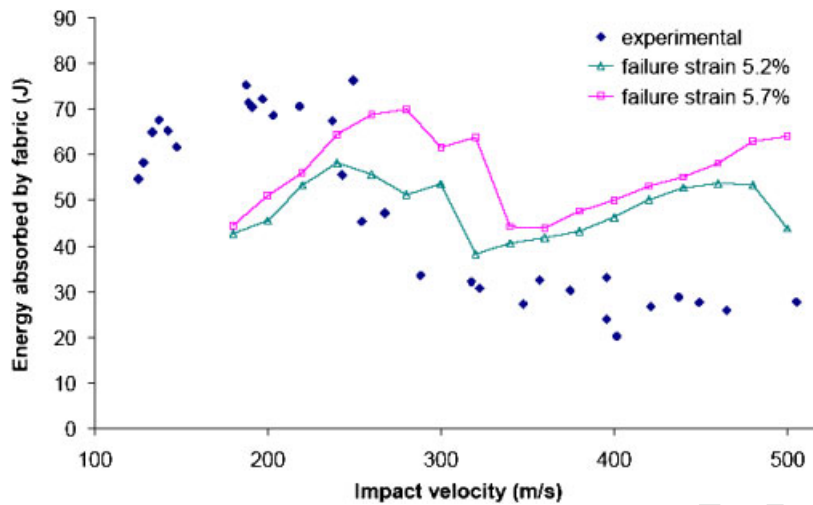


Figure 11. Variation of energy absorbed by fabric with impact velocity—simulation and ballistic test results.

1 comparisons are possible. Quantitative comparisons will require actual material properties of  
 2 the yarns from the fabric tested. A comparison of ballistic tests results for Twaron<sup>®</sup> fabric  
 3 with simulation using yarn properties from Gu [22] verifies that the mechanisms of fabric  
 4 deformation and perforation are reflected by the model and that predicted trends in energy  
 5 absorption are similar to experimental results.

## 7. EFFECTS OF FRICTION COEFFICIENT ON ENERGY ABSORBED BY FABRIC

7 A parametric study on the effects of friction coefficient was carried out by varying the coefficient  
 8 of friction between yarns from 0 to 1. The friction coefficient between the projectile and the yarn  
 9 is kept constant as 0.27 as measured in Reference [23] for all simulations. Different yarn–yarn  
 10 friction coefficients resulted in differences in the penetration processes, as shown in Figure 12.  
 11 A low friction ( $\mu=0.02$ ) corresponds to a fabric with very low structural coherence. It allows  
 12 a projectile to push yarns aside easily and adversely affects the ability of the fabric to retard  
 13 the projectile. For a more realistic friction level ( $\mu=0.2$ ), the numerical model predicts fabric  
 14 perforation that correlates closely with experimental observations. It also shows that unclamped  
 15 yarns are less likely to rupture than clamped yarns. A high contact friction ( $\mu=1.0$ ) reduces  
 16 yarn mobility and inter-yarn sliding. Hence, yarns in both unclamped and clamped directions  
 17 break at the same time.

18 The trends of energy absorption by fabric for different yarn friction coefficients are presented  
 19 in Figure 13. The range of coefficient is categorized into three groups—low ( $\mu=0-0.06$ ),  
 20 moderate ( $\mu=0.06-0.2$ ) and high ( $\mu=0.2-1$ )—based on their effects on energy absorption  
 21 by the fabric. For a low friction coefficient, the energy absorbed by the fabric increases with  
 22 impact velocity as well as with the friction coefficient. Energy absorption is strongly sensitive to  
 23 the magnitude of friction coefficient within this range of coefficients. For moderate coefficients



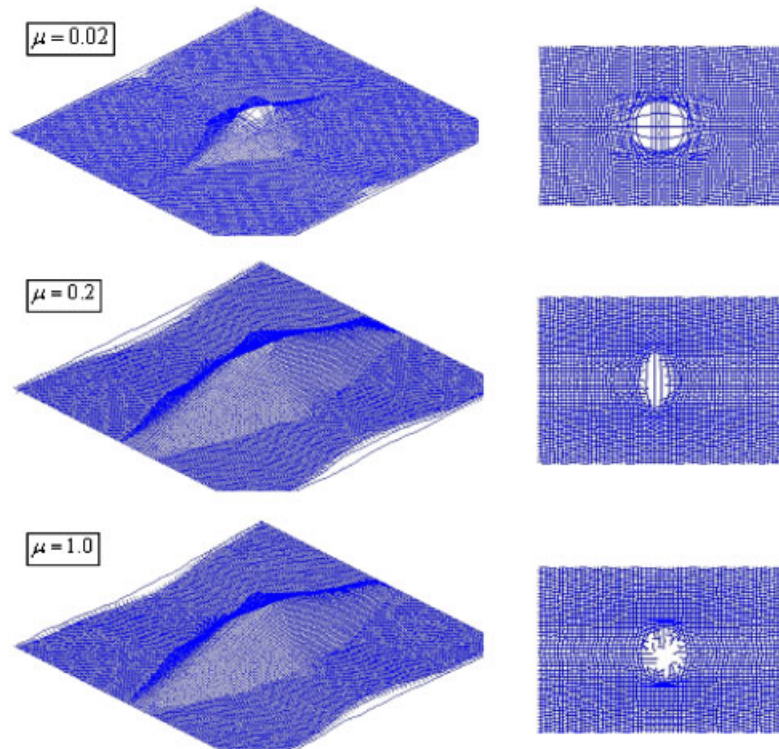


Figure 12. Simulation results for fabric perforation with different inter-yarn friction coefficients (impact velocity = 260 m/s).

1 of friction ( $\mu = 0.06\text{--}0.2$ ), energy absorption by the fabric increases with impact velocity until  
 2 about 250 m/s before it decreases with further increase in impact velocity. A friction coefficient  
 3 of  $\mu = 0.2$  yields the highest energy absorption for the whole range of impact velocities. At high  
 4 friction coefficients greater than 0.2, the energy absorbed start to decrease when the friction  
 5 coefficient is increased.

6 A better understanding of the effects of friction coefficient can be obtained by studying  
 7 the fabric response at specific impact velocities for different friction coefficients. Figure 14  
 8 presents the variation of projectile velocity with time for an impact velocity of 260 m/s. The  
 9 gradient is very gentle for low friction coefficients ( $\mu = 0\text{--}0.08$ ). Primary yarns which are in  
 10 direct contact with the projectile gradually slip aside during impact, leaving fewer and fewer  
 11 primary yarns to retard the projectile. Hence, the fabric offers a very low projectile resistance.  
 12 A friction coefficient of at least 0.1 is required to prevent yarns from slipping off the projectile  
 13 and to maintain integrity of the weave at the impact point. The rate of decrease in projectile  
 14 velocity is greatest for  $\mu = 0.2$ . In the case of  $\mu = 1.0$ , the projectile velocity history is similar  
 15 to that for  $\mu = 0.2$  except that perforation occurs earlier at 73  $\mu\text{s}$  after impact. A high friction  
 16 coefficient ( $\mu = 1.0$ ) restrains relative slippage between yarns and causes stress concentration  
 17 at the impact point, leading to a shorter time to failure.

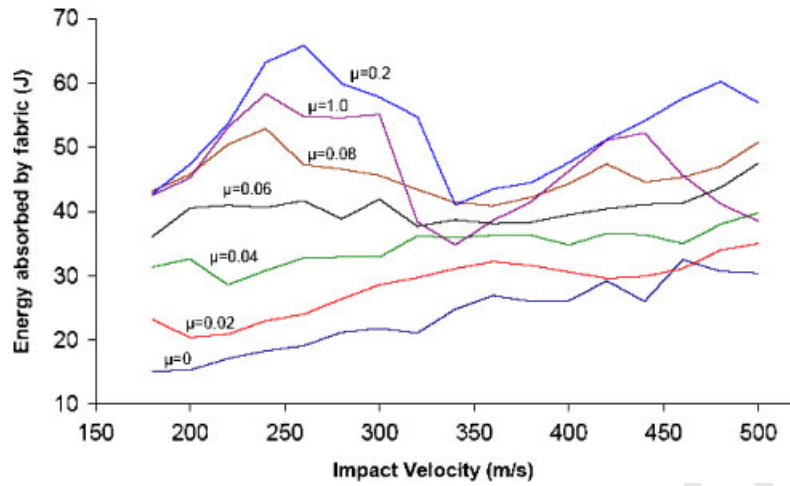


Figure 13. Variation of energy absorbed by fabric with impact velocity for different yarn friction coefficients.

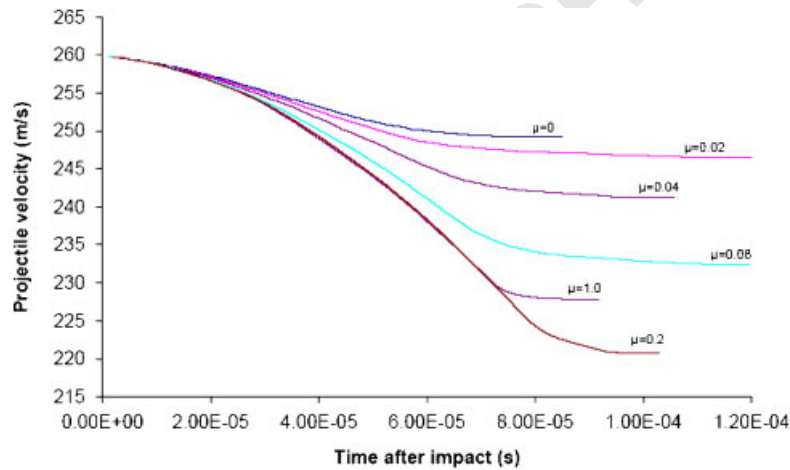


Figure 14. Projectile velocity histories for 260 m/s impact on fabric with different yarn friction coefficients.

1 The projectile impact energy is dissipated by the fabric via kinetic energy transferred to  
 2 the fabric and strain energy in deformed yarns. The variation of these energy components for  
 3 an impact velocity of 260 m/s is shown in Figure 15 for different friction coefficients. Both  
 4 the kinetic energy and strain energy increase with the friction coefficient up to a value of  
 5  $\mu=0.2$ . For higher values of inter-yarn friction, the energy trends appear similar up to the  
 6 point of perforation, which occurs earlier for a higher friction coefficient. This is exemplified  
 7 by the curves for  $\mu=0.2$  and  $\mu=1.0$ . Although yarn friction clearly affects the amount energy

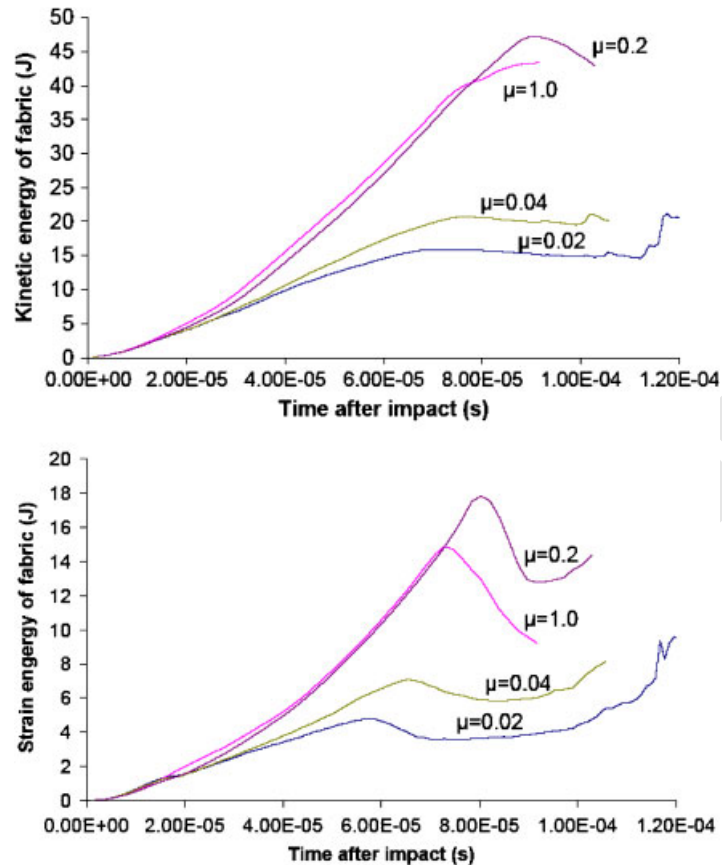


Figure 15. Kinetic energy and strain energy histories for impact at 260 m/s on fabric with different inter-yarn friction coefficients.

- 1 absorbed by the fabric, energy losses through friction are insignificant. For a friction coefficient of  $\mu=0.2$ , frictional losses are only about 1 J.

### 3 8. EFFECTS OF INTER-YARN FRICTION ON THE FABRIC BALLISTIC LIMIT

5 The ballistic limit (BL) is another measure of the effectiveness of fabric armour in resisting  
 6 ballistic impact. As shown in Figure 16, the BL increases with inter-yarn friction for values  
 7 of  $\mu$  from 0 to 0.1. The BL for  $\mu=0.1$  is almost twice that for perfectly smooth yarns ( $\mu=0$ ).  
 8 The BL for  $\mu=0.2$  (70 m/s) is the same as the BL for  $\mu=0.1$  (71 m/s). Further increases  
 9 in the friction coefficient, from a value of 0.2–0.6 yields little difference in the BL. As a  
 reference, the BL obtained from a fabric model with pin-jointed crossover points that does not  
 cater for sliding between yarns [8] is also plotted. This can be regarded as an extreme case

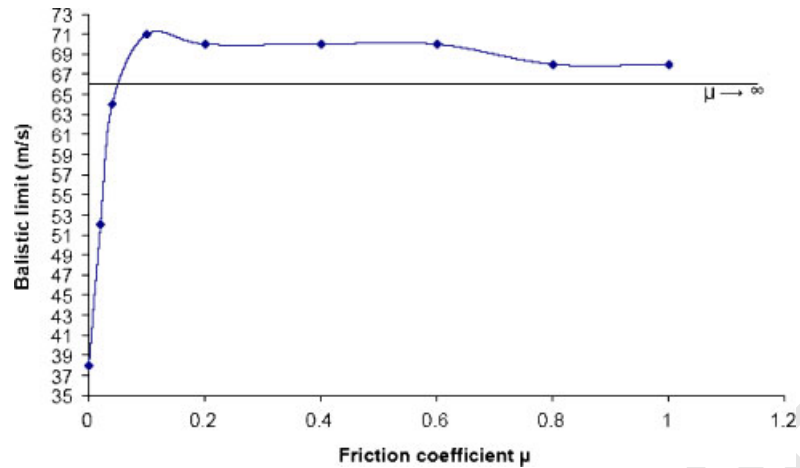


Figure 16. Variation of predicted ballistic limit with inter-yarn friction coefficient.

- 1 where the coefficient of friction is infinite ( $\mu = \infty$ ). The results show that an excessively high  
 2 friction coefficient has a negative effect on the BL.
- 3 Figure 17 shows predictions of fabric deformation at the BL for different inter-yarn friction  
 4 coefficients. Perforation of the fabric for the case of  $\mu = 0.02$  corresponds to extensive slippage  
 5 of yarns over the projectile near the point of impact. No yarns are broken and yarn ravelling  
 6 is obvious at the unclamped edges. The woven structure has poor integrity because of the  
 7 low contact friction between yarns. A friction coefficient of  $\mu = 0.1$  is required to preserve the  
 8 structural coherence of fabric and for the crossover points of yarns to transmit energy effectively  
 9 throughout the fabric. Hence, the BL is much higher than that for  $\mu = 0.02$ . Deformation and  
 10 perforation of the fabric at the BL for  $\mu = 0.2$  and  $\mu = 1.0$  do not show much difference. The  
 11 decrease in BL for higher friction coefficients of  $\mu = 0.8-1.0$  arises from stress concentrations  
 at the region of impact because of lack of mobility between yarns.

13

## 9. CONCLUSIONS

14 A numerical model for woven fabric armour, comprising a network of pin-jointed linear  
 15 viscoelastic elements is developed and implemented through an in-house code. Accommodation  
 16 of inter-yarn slippage is incorporated into the model to facilitate a more realistic representation  
 17 of the woven architecture of fabrics. This is effected with minimal degrees of freedom and yet  
 able to identify the effects of friction between yarns.

18 The viscoelastic parameters of the linear elements were obtained by fitting a Zener  
 19 viscoelastic model to experimental data reported by Gu [22]. Comparisons of computational  
 20 simulations with actual ballistic tests show that the proposed model yields good predictions  
 21 of the fabric deformation and perforation mechanisms, as well as the fabric energy absorption  
 22 characteristics.

23

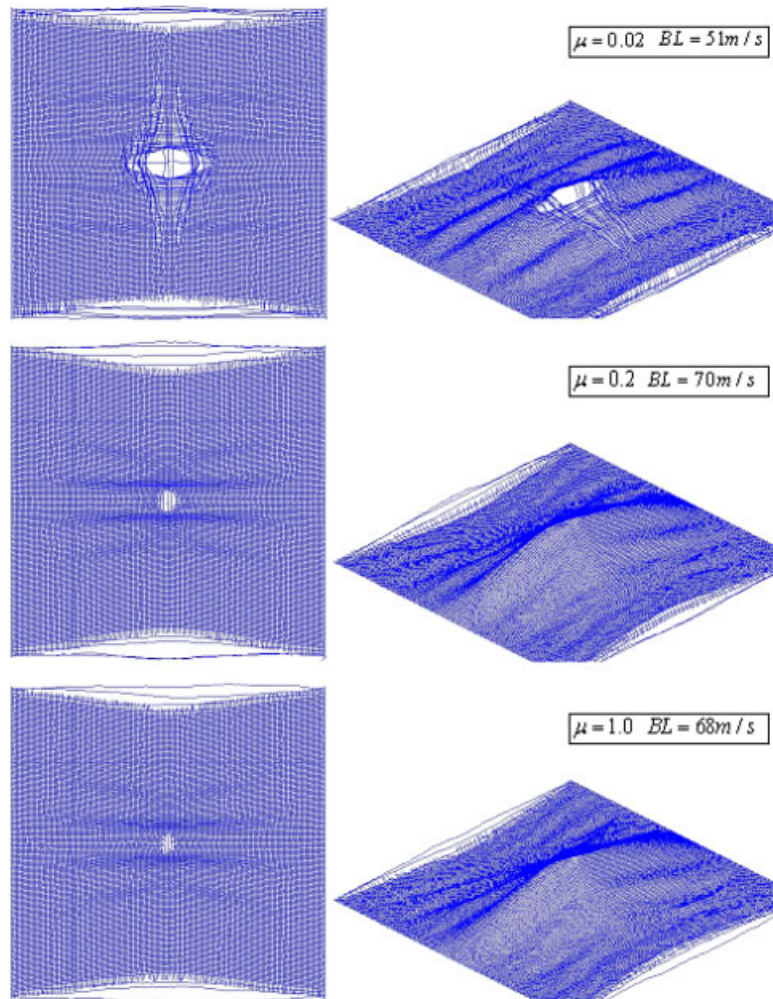


Figure 17. Simulation results for fabric with different inter-yarn friction coefficients perforated at the ballistic limit.

- 1 Simulations of ballistic impact on woven fabric with different values of inter-yarn friction
- 2 show that the influence of friction on ballistic performance can be categorized into three
- 3 levels of friction—low friction ( $\mu=0-0.06$ ), moderate friction ( $\mu=0.06-0.2$ ) and high friction
- 4 ( $\mu>0.2$ ). Low friction coefficients lead to poor weave integrity and undermine the ability of the
- 5 fabric to arrest projectiles. Excessive friction reduces relative sliding between yarns, resulting in
- 6 high stress concentrations at the point of impact and give rise to premature yarn rupture. The
- 7 ballistic resistance of a woven fabric appears to remain constant when the inter-yarn friction
- coefficient exceeds 0.2.



1

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