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Citation: Journal of Applied Physics **119**, 245103 (2016); doi: 10.1063/1.4954683 View online: http://dx.doi.org/10.1063/1.4954683 View Table of Contents: http://scitation.aip.org/content/aip/journal/jap/119/24?ver=pdfcov Published by the AIP Publishing

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Lithium concentration dependent structure and mechanics of amorphous silicon

H. S. Sitinamaluwa,¹ M. C. Wang,¹ G. Will,¹ W. Senadeera,¹ S. Zhang,² and C. Yan^{1,a)}

¹School of Chemistry, Physics and Mechanical Engineering, Queensland University of Technology (QUT), Brisbane QLD 4001, Australia

²Centre for Clean Environment and Energy, Environmental Futures Research Institute and Griffith School of Environment, Gold Coast Campus, Griffith University, QLD 4222, Australia

(Received 21 March 2016; accepted 11 June 2016; published online 23 June 2016)

A better understanding of lithium-silicon alloying mechanisms and associated mechanical behavior is essential for the design of Si-based electrodes for Li-ion batteries. Unfortunately, the relationship between the dynamic mechanical response and microstructure evolution during lithiation and delithiation has not been well understood. We use molecular dynamic simulations to investigate lithiated amorphous silicon with a focus to the evolution of its microstructure, phase composition, and stress generation. The results show that the formation of Li_x Si alloy phase is via different mechanisms, depending on Li concentration. In these alloy phases, the increase in Li concentration results in reduction of modulus of elasticity and fracture strength but increase in ductility in tension. For a Li_x Si system with uniform Li distribution, volume change induced stress is well below the fracture strength in tension. *Published by AIP Publishing*. [http://dx.doi.org/10.1063/1.4954683]

I. INTRODUCTION

Li-ion battery (LIB) is increasingly used in various portable electronic devices and electric vehicles (EV) due to high energy/power density and long lifespan.^{1,2} To reduce the structural weight in transportation systems, novel Li-ion batteries with higher energy density is required.³ Silicon, with its highest known theoretical specific capacity (i.e., 4200 mA h g^{-1}) has emerged as a promising candidate as Liion battery anode material.⁴ However, mechanical integrity of Si remains as a challenge due to its huge volume expansion (\sim 300%).⁵ When subjected to constraints imposed by electrode architecture, such enormous expansion can cause mechanical stress and fracture.^{5,6} Various strategies have been developed to improve the mechanical stability of silicon, including nanostructured electrodes,^{7,8} flexible binder materials,⁹ and carbon based hybrids.¹⁰ Recently, there has been increasing interest in amorphous silicon (a-Si), owing to its robust lithiation behavior compared with its crystalline counterpart.11,12

Lithiation of amorphous silicon occurs via formation of metastable Li_xSi alloy phases,¹³ identified as $Li_{15}Si_4$ for fully lithiated silicon.¹⁴ McDowell *et al.* showed that a-Si nanospheres undergo up to 332% volume expansion at full lithiation.¹¹ Several authors have examined mechanical behavior of electrochemically lithiated silicon. Berla *et al.*¹⁵ performed nanoindentation tests to evaluate the modulus and hardness of lithiated amorphous silicon as a function of Li concentration in active material. Also, *in-situ* atomic force microscopy (AFM) and transmission electron microscopy (TEM) have been used to measure Young's modulus and tensile strength of lithiated silicon.^{16,17} Furthermore, Berla *et al.* showed viscoelastic flow and creep behavior in lithiated silicon, indicating the possibility of plastic deformation

under high Li concentrations.¹⁵ In addition, recent experiments and theoretical studies suggest that lithiation induced plasticity can effectively prevent fracture.^{18,19} These studies demonstrated a strong dependence of mechanical properties on phase composition. From recent first-principle calculations, the microstructural evolution such as Si-Si bond breaking and regeneration has been explained as responsible for the composition-dependent mechanical properties of lithiated silicon.^{20–22} However, due to limitations in length-scale, the model cells in those first-principles calculations cannot fully represent the typical a-Si. Therefore, more work is required to understand the underlying structure-property relationship in lithiated silicon.

In this study, we use molecular dynamic (MD) simulations to investigate the volume expansion and stress development of amorphous silicon during lithiation and their dependence with phase composition and atomic structures. Through quantitative analysis of atomic structures, we find that the formation of Li_xSi phase is via different mechanisms, depending on the Li concentration. Furthermore, our results indicate that the changed atomic mobility in lithiated silicon structures play a major role in transient mechanical properties. The outcomes of this work are believed to be useful for design of Si based electrodes with improved mechanical stability.

II. METHODS

In the molecular dynamic (MD) simulation, the initial amorphous silicon (a-Si) atomic model is a cubic unit with a side length 2.7 nm, which contains 1000 silicon atoms as shown in Figure 1(a). The a-Si model was created by quenching molten silicon from $2400 \,^{\circ}$ C to room temperature at a cooling rate of $10^{12} \,\mathrm{K \, s^{-1}}$. A modified Stillinger-Weber potential for a-Si was used.²³ After structural relaxation, the density of the generated a-Si was calculated as 2.291 g cm⁻³. The a-Si

^{a)}E-mail: c2.yan@qut.edu.au



FIG. 1. (a) Schematic illustration (left panel) and atomic structure (right panel) of the lithiation process of the 1000-atom (a-Si) model constrained by a substrate. (b) A comparison between the radial distribution function (RDF) results determined by MD simulation and previous experiment.²⁴ (c) Bond angle distribution function of the a-Si atomic model.

atomic model was validated by comparing the radial distribution function (RDF) with the existing experimental data,²⁴ as shown in Figure 1(b). The broad first-neighbor peak and the absence of well-defined higher-neighbor peaks in the RDF confirm its amorphous structure. The average Si-Si coordination number (CN) was estimated to be about 4, confirming a tetrahedral structure as reported in literature.²⁴ Further, as shown in Figure 1(c), the average bond angle is 107.3° with standard deviation 18.9°, in good agreement with previous simulations.^{25,26}

For crystalline silicon, lithiation may occur via a twophase mechanism;^{27,28} however, detailed lithiation mechanism of a-Si still remains unclear.^{29–32} To understand the effect of Li concentration, it is reasonable to assume a uniform dispersion of Li ion in the a-Si. To simulate lithiation-delithiation, a stepwise Li insertion and extraction procedure was used.³³ For each insertion step, 25 Li atoms were randomly distributed in the a-Si cell, and the structural relaxation was carried out. Energy minimization was achieved when the forces on atoms were less than 10^{-8} eV/Å. The simulation cell was equilibrated at 300K in isothermal-isobaric (NPT) ensemble for further 10 ps. This procedure was repeated until the desired Li concentration was obtained. Similarly, delithiation was also simulated stepwise in which 25 Li atoms were removed randomly, and the structure was relaxed in the NPT ensemble. Atomic interactions during Li insertion and extraction were defined by using modified embedded atom method (MEAM) based interatomic potential for Li-Si alloys.³⁴

When the a-Si anode is attached to a rigid substrate as shown in Figure 1(a), non-zero intrinsic stress generates along the substrate (in-plane) direction due to volume expansion.^{19,33,35} To estimate such intrinsic stress in lithiated a-Si, the in-plane (x- and y-axis) lengths of the simulation cell

were fixed to mimic the constraint induced by the rigid substrate, and the volume expansion along the thickness (z axis) direction was set free. Based on the virial theorem,^{36,37} the atomic stresses along the in-plane direction were calculated after each lithiation step by running MD simulations under the NVT ensemble for further 10 ps. Li atoms distribution in the cell is approximately uniform, to avoid the concentration gradient dependent stress.³⁸ To understand the mechanical properties of lithiated a-Si, uniaxial tensile tests along x axis (as shown in Figure 1(a)) were also conducted to calculate the fracture strengths of several lithiated a-Li_xSi structures at a constant strain rate of 10^8 s^{-1} .^{39,40} All simulations were performed using LAMMPS software,⁴¹ which has been used for simulating mechanical properties of nanomaterials in our previous studies.⁴²⁻⁴⁴ Atomic structure analysis and visualization were carried out using OVITO⁴⁵ and VMD⁴⁶ postprocessing tools.

III. RESULTS AND DISCUSSION

A. Lithiation-induced microstructural transformations

Amorphous silicon has a disordered tetrahedral structure where Si atoms are covalently bonded forming continuous random network (CRN).²⁴ With lithiation process, Li atoms react with Si to form Li_xSi alloy phase, with Li to Si ratio (*x*) increasing up to 4.4.⁴ In this work, Li_xSi alloys (x = 0 to 3.75) were simulated as these Li to Si ratios have been observed in room temperature lithiation.¹⁴ Upon lithiation, gradual breakdown of a-Si CRN and formation of new Li-Si bonds are observed, leading to volume expansion. The maximum volume expansion observed in current model at full lithiation (Li_{3.75}Si) was 283%, in agreement with experiment.¹¹



FIG. 2. (a) Atomic structures of Li_xSi at different Li concentrations during one lithiation cycle. (b) Si-Si radial distribution function (RDF) at different stages of lithiation and delithiation.

Figure 2(a) shows the atomic structures at different stages of lithiation. For clarification, only Si-Si bonds are shown. The radial distribution function (RDF) of Si-Si and Li-Si pairs at different stages of lithiation was evaluated. RDF is an indicator of how atoms are packed in space and can provide short and medium range structure information, which can be expressed as

$$g(r) = \frac{N(r, r + \Delta r)}{4\pi r^2 \Delta r \rho},$$
(1)

where r is the distance between the reference particle and found particle; $N(r, r + \Delta r)$ represents the number of atoms in specified pairs within the region $(r, r + \Delta r)$, and ρ is the total bulk density of the considered particle.

Figure 2(b) shows the Si-Si RDFs for a-Si and a-Li_xSi at different composition. The initial bulk silicon structure (x=0) do not show sharp second neighbor peak due to the lack of long range ordered atomic structure. The largest Si-Si pair distance corresponding to the nearest neighbor peak is 2.6 Å, which is the cut-off distance for the Si-Si bond. Similarly, the Si-Li bond cut-off distance is calculated as 3.3 Å. The Si-Si RDFs as a function of Li concentration, during lithiation and delithiation, are shown in Figure 2(b). With the increase in Li, the Si-Si nearest neighbor peak shifts to larger pair distance, indicating weakened Si-Si bonds.⁴⁷ For example, at x=0, the largest Si-Si pair distance is 2.6 Å, and increases to 2.9 Å at full lithiation. Also, the reduced peak height during lithiation indicates the break of Si-Si bonds. During delithiation, the Si-Si bonds can be recovered, as indicated by the increased peak height. Further, the Si-Si nearest neighbor peak is visible at x = 3.75, associated with the unbroken Si-Si bonds. The bond break and recovery can be observed in the atomic structures shown in Figure 2(a).

The atomic coordination numbers (CN) were calculated by taking integral of the nearest neighbor peak of RDF for corresponding atomic pair. In Figure 3(a), Si-Si CN represents the average number of coordinated Si, while Si-Li CN represents the average number of Li coordinated by a Si atom. At x = 0, CN of a-Si structure is ~ 4 , as atoms are tetrahedrally bonded. With increasing Li concentration, Si-Si CN decreases due to break of Si-Si bonds. Simultaneously, Si-Li CN increases with Li invasion due to the formation of new Si-Li bonds. The CN curves show different trends when Li concentration (x) is above and below 2.5. Below x = 2.5, Si-Si CN gradually decreases and is saturated around ~ 1 . This is accompanied by the increase of Si-Li CN, indicating that a rapid breakdown of Si-Si bonds takes place below x = 2.5. Also, Si-Si CN is retained at ~ 4 until x reaches 0.25. First-principles calculation indicates that Li atoms initially take the positions at interstices of larger silicon rings,²⁰ with only minor distortions to the a-Si continuous random network (CRN). It is also reported that four Li atoms at the tetrahedral sites can result in a very small energy barrier (0.08 eV) against the break of Si-Si bond, and thus, it is believed that at least four Li atoms are needed to weaken and break a Si-Si bond.²²

We further analyze the breakdown of amorphous silicon structure via cluster analysis. A group of silicon atoms (two or more) connected with the Si-Si bonds is considered as a silicon cluster. At x = 0, the bulk silicon structure is one big cluster of 1000 atoms, and every silicon atom is a part of tetrahedral CRN. It is interesting to note that the silicon structure does not start to break until *x* exceeds ~0.25. At x = 3.75, silicon atoms are mainly arranged in three different



FIG. 3. (a) Si-Si and Si-Li coordination numbers during lithiation as a function of Li concentration. (b) Breakdown of silicon CRN into number of smaller structures during lithiation via cluster analysis.

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FIG. 4. Stress development in constrained a-Si thin film during lithiation/ delithiation process.

structures: (i) lithiated individual silicon atoms, (ii) lithiated Si-Si dumbbells, and (iii) lithiated silicon clusters of 3–9 atoms. Figure 3(b) gives the percentage of these structures as a function of Li concentration. In Figure 3(b), a rapid breakdown of large silicon clusters and formation of smaller structures observed up to $x \sim 2.5$. At x = 2.5, about 90% silicon atoms are in the form of dumbbells, small clusters or individual atoms and the rest 10% are unbroken larger clusters. This suggests that ~90% Si structure's break down takes place below x = 2.5. As indicated in Figure 3(a), a silicon atom has an average of 8.3 Li neighbors at x = 2.5, which increases up to 9.3 in full lithiation (x = 3.75). Therefore, lithiation beyond x = 2.5 creates more Li-rich phases around silicon clusters without significant structural change.

Based on the structure analysis, lithiation is featured by a three-step process. The first step occurs below $x \sim 0.25$, where Li atoms occupy interstitial sites in silicon structure with only minor distortions to a-Si CRN. The second step corresponds to the region between $x \sim 0.25$ and $x \sim 2.5$, where a-Si CRN structure undergoes considerable transformations via a rapid Si-Si bond breaking. In this step, the a-Si structure is mainly broken down to lithiated silicon atoms, dumbbells, and smaller clusters. Above $x \sim 2.5$, further lithiation creates more Li-rich phases with relatively less Si-Si bond breaking.

B. Effects of structural transformation on stress and deformation of a-Si electrode

Figure 4 shows lithiation-delithiation induced stress in a thin film electrode as a function of Li concentration.

During lithiation, a sharp increase in compressive stress up to ~ 2 GPa is observed as x = 0.3. It is followed by stress relaxation, where compressive stress decreases to ~ 1 GPa as x = 2.5. Finally, almost constant stress of ~ 1 GPa is observed for the rest of lithiation. It is interesting to note that lithiation induced stress shows a clear three-step response which corresponds to the structural transformations described in Section III A.

In the first step, Li atoms occupy interstices between silicon rings, without breaking the tetrahedral CRN. As the covalent Si-Si bonds dominating the structure at this stage, volume expansion creates significant compressive stress. As Li concentration increases up to x = 2.5, silicon structure breaks down to form Si-Li phases, with apparent stress relaxation. Zhao et al.²² demonstrated that Si-Li pair has a mixed ionic-covalent bond, weaker than covalent Si-Si bonds. Thus, stress relaxation observed at this stage may be attributed to the increased number of weak Li-Si pairs. This behavior is also described as plastic flow of Li_xSi structure.^{19,22} With more Li invasion (up to $x \sim 2.5$), compressive stress gradually decreases to ~ 1 GPa. Interestingly, when x > 2.5, no further increase in stress is observed. In general, plastic flow at a constant stress is usually observed when diffusing atoms are trapped in structural defects. For example, high temperature creep of some alloys is found to occur under constant stress, due to trapping of diffusing atoms in grain boundaries.⁴⁸ In this work, silicon clusters may act as diffusion barrier, contributing to constant-stress flow. During delithiation, the model volume starts to contract and elastic unloading occurs, where restriction from the substrate converts the stress from compression to tension, as shown in Figure 4. There is a sharp increase in tensile stress at lower Li concentrations, i.e., below ~ 0.5 . This is associated with the possible recovery of covalent Si bonds at low Li-ion concentrations, which results in significant increases in stress along with volume contraction.

To understand the intrinsic mechanical behavior of the lithiated Si, a uniaxial tensile test was simulated in several Li_xSi structures. The stress-strain curves of the uniaxial tensile tests are shown in Figure 5(a). As shown in Table I, the modulus of elasticity of the lithiated silicon decreases with increasing Li concentration. The stress-strain response of those Li_xSi with x = 0 and x = 0.25 is like brittle materials, and no notable plastic deformation is observed. In contrast, considerable plastic deformation can be observed in the Li_xSi of high Li concentrations, which can be attributed to



FIG. 5. (a) Stress-strain curves of a-Si and a-Li_xSi structure in uniaxial tension. (b) A comparison of tensile fracture stresses of Li_xSi structures with delithiation induced tensile stress.

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TABLE I. Mechanical characteristics of LixSi alloys and mean square displacement of Si atoms (MSD)Si during tensile test.

Structure	Modulus of elasticity (GPa)	Stress at fracture (GPa)	(MSD) _{Si} after 1000 ps (Å) ²
Bulk a-Si	125.6	10.0	12.6
Li _{0.25} Si	118.5	8.8	13.8
LiSi	80.7	6.0	27.8
Li ₂ Si	43.6	3.3	62.3
Li _{3.75} Si	39.8	2.7	198.8

the Li-rich phases with higher mobility. To further confirm this, the total mean square displacement (MSD) of silicon atoms during the tensile test was evaluated, as shown in Table I. It is clear that MSD increases with increasing Li, indicating higher atomic mobility with higher number of Li-Si pairs.

Figure 5(b) shows the fracture stress of the Li_xSi structures in the tensile test as a function of Li concentration. The fracture stress is found to be decreasing with the increase in Li concentration. In Figure 5(b), the simulated tensile stress during delithiation is also included for comparison. It is clear that the delithiation induced stress well below the fracture stresses in tension. This implies that the tensile stress induced by volume change in the constrained model is not sufficient to trigger fracture. Therefore, amorphous silicon electrodes may fail under repeated compressive-tensile stresses during charge-discharge cycles. Further work along this direction is required.

IV. CONCLUSIONS

In this work, molecular dynamic simulation was used to evaluate the evolution of microstructure and mechanical stresses in amorphous silicon under electrochemical lithiation-delithiation process. Through the short and medium range atomic structure analysis, it was found that Li-Si alloying reaction occurs via a three-stage mechanism. First, Li atoms occupy interstitial sites of amorphous silicon, followed by the second stage, i.e., rapid breaking of the original Si-Si structure to form $\sim Li_{2.5}Si$. Further lithiation leads to the formation of Li_{3.75}Si, accompanied by a relatively less bond breaking. At full lithiation (Li₁₅Si₄), the lithiated Si-Si dumbbells and small silicon clusters dominate the structure. With increasing Li concentration, reduced elastic modulus and fracture strength but increased ductility is observed in Li_xSi phases. Cyclic compressive-tensile stresses can be generated in geometrically laterally constrained amorphous silicon during lithiation-delithiation process. Interestingly, these stresses are found to be well below the fracture stress in tension, implying that the mechanical damage in the Si-based Li battery electrodes is likely to be triggered by the cyclic tensile and compressive stresses.

ACKNOWLEDGMENTS

The authors wish to acknowledge the High Performance Computing (HPC) and CARF at QUT for access to their facilities. The work was partially supported by the ARC DP project (DP150101717). H.S.S. would like to acknowledge Queensland University of Technology for APA, IPRS, and QUT excellence top up scholarships.

- ²K. E. Aifantis, S. A. Hackney, and R. V. Kumar, *High Energy Density* Lithium Batteries: Materials, Engineering, Applications (Wiley-VCH, Weinheim, 2010).
- ³M. S. Whittingham, Proc. IEEE 100, 1518 (2012).
- ⁴W.-J. Zhang, J. Power Sources **196**, 13 (2011).
- ⁵A. Mukhopadhyay and B. W. Sheldon, Prog. Mater. Sci. 63, 58 (2014).
- ⁶X. Gao, Z. Ma, W. Jiang, P. Zhang, Y. Wang, Y. Pan, and C. Lu, J. Power Sources 311, 21 (2016).
- ⁷J. R. Szczech and S. Jin, Energy Environ. Sci. 4, 56 (2011).
- ⁸H. Wu and Y. Cui, Nano Today 7, 414 (2012).
- ⁹M. Ling, Y. Xu, H. Zhao, X. Gu, J. Qiu, S. Li, M. Wu, X. Song, C. Yan, and G. Liu, Nano Energy 12, 178 (2015).
- ¹⁰H. Sitinamaluwa, S. Zhang, W. Senadeera, G. Will, and C. Yan, Mater. Technol. (published online).
- ¹¹M. T. McDowell, S. W. Lee, J. T. Harris, B. A. Korgel, C. Wang, W. D. Nix, and Y. Cui, Nano Lett. 13, 758 (2013).
- ¹²L. A. Berla, S. W. Lee, I. Ryu, Y. Cui, and W. D. Nix, J. Power Sources 258, 253 (2014).
- ¹³M. T. McDowell, S. W. Lee, W. D. Nix, and Y. Cui, Adv. Mater. 25, 4966 (2013).
- ¹⁴J. Y. Kwon, J. H. Ryu, and S. M. Oh, Electrochim. Acta 55, 8051 (2010).
- ¹⁵L. A. Berla, S. W. Lee, Y. Cui, and W. D. Nix, J. Power Sources 273, 41 (2015).
- ¹⁶Q. P. McAllister, K. E. Strawhecker, C. R. Becker, and C. A. Lundgren, J. Power Sources 257, 380 (2014).
- ¹⁷A. Kushima, J. Y. Huang, and J. Li, ACS Nano 6, 9425 (2012).
- ¹⁸K. Zhao, M. Pharr, S. Cai, J. J. Vlassak, and Z. Suo, J. Am. Ceram. Soc. 94, s226 (2011).
- ¹⁹V. A. Sethuraman, M. J. Chon, M. Shimshak, V. Srinivasan, and P. R. Guduru, J. Power Sources 195, 5062 (2010).
- ²⁰S. Huang and T. Zhu, J. Power Sources **196**, 3664 (2011).
- ²¹M. K. Chan, C. Wolverton, and J. P. Greeley, J. Am. Chem. Soc. 134, 14362 (2012).
- ²²K. Zhao, W. L. Wang, J. Gregoire, M. Pharr, Z. Suo, J. J. Vlassak, and E. Kaxiras, Nano Lett. 11, 2962 (2011).
- ²³R. L. C. Vink, G. T. Barkema, W. F. van der Weg, and N. Mousseau, J. Non-Cryst. Solids 282, 248 (2001).
- ²⁴K. Laaziri, S. Kycia, S. Roorda, M. Chicoine, J. Robertson, J. Wang, and S. Moss, Phys. Rev. Lett. 82, 3460 (1999).
- ²⁵A. France-Lanord, E. Blandre, T. Albaret, S. Merabia, D. Lacroix, and K. Termentzidis, J. Phys.: Condens. Matter 26, 055011 (2014).
- ²⁶E. Holmström, B. Haberl, O. Pakarinen, K. Nordlund, F. Djurabekova, R. Arenal, J. Williams, J. Bradby, T. Petersen, and A. Liu, J. Non-Cryst. Solids 438, 26 (2016).
- ²⁷X. H. Liu, A. V. Davydov, S. X. Mao, S. T. Picraux, S. L. Zhang, J. Li, T. Zhu, J. Y. Huang, J. W. Wang, S. Huang, F. F. Fan, X. Huang, Y. Liu, S. Krylyuk, J. Yoo, and S. A. Dayeh, Nat. Nanotechnol. 7, 749 (2012).
- ²⁸Z. Xie, Z. Ma, Y. Wang, Y. Zhou, and C. Lu, RSC Adv. 6, 22383 (2016).
- ²⁹J. W. Wang, Y. He, F. Fan, X. H. Liu, S. Xia, Y. Liu, C. T. Harris, H. Li, J. Y. Huang, S. X. Mao, and T. Zhu, Nano Lett. 13, 709 (2013).
- ³⁰H. Liao, K. Karki, Y. Zhang, J. Cumings, and Y. Wang, Adv. Mater. 23, 4318 (2011).
- ³¹H. Ghassemi, M. Au, N. Chen, P. A. Heiden, and R. S. Yassar, ACS Nano 5, 7805 (2011).
- ³²G. A. Tritsaris, K. Zhao, O. U. Okeke, and E. Kaxiras, J. Phys. Chem. C 116, 22212 (2012).
- ³³F. F. Fan, S. Huang, H. Yang, M. Raju, D. Datta, V. B. Shenoy, A. C. T. van Duin, S. L. Zhang, and T. Zhu, Modell. Simul. Mater. Sci. Eng. 21, 074002 (2013).
- ³⁴Z. Cui, Z. Cui, F. Gao, and J. Qu, J. Power Sources **207**, 150 (2012).
- ³⁵V. A. Sethuraman, M. J. Chon, M. Shimshak, N. Van Winkle, and P. R. Guduru, Electrochem. Commun. 12, 1614 (2010).
- ³⁶D. H. Tsai, J. Chem. Phys. 70, 1375 (1979).
- ³⁷M. Zhou, Proc. R. Soc. London, Ser. A **459**, 2347 (2003).
- ³⁸I. Ryu, J. W. Choi, Y. Cui, and W. D. Nix, J. Mech. Phys. Solids **59**, 1717 (2011).
- ³⁹M. Wang, C. Yan, L. Ma, and N. Hu, Comput. Mater. Sci. 68, 138 (2013).

¹J.-M. Tarascon and M. Armand, Nature **414**, 359 (2001).

- ⁴⁰M. Wang, C. Yan, L. Ma, N. Hu, and M. Chen, Comput. Mater. Sci. 54, 236 (2012).
- ⁴¹S. Plimpton, J. Comput. Phys. **117**, 1 (1995).
- ⁴²M. C. Wang, C. Yan, L. Ma, N. Hu, and G. P. Zhang, Comput. Mater. Sci. **75**, 69 (2013).
- ⁴³M. C. Wang, C. Yan, D. Galpaya, Z. B. Lai, L. Ma, N. Hu, Q. Yuan, R. X. Bai, and L. M. Zhou, J. Nano Res. 23, 43 (2013).
- ⁴⁴Y. Xu, M. Wang, N. Hu, J. Bell, and C. Yan, RSC Adv. 6, 28121 (2016).
 ⁴⁵A. Stukowski, Modell. Simul. Mater. Sci. Eng. 18, 015012 (2010).
- ⁴⁶W. Humphrey, A. Dalke, and K. Schulten, J. Mol. Graphics **14**, 33 (1996).
- ⁴⁷H. Kim, C.-Y. Chou, J. G. Ekerdt, and G. S. Hwang, J. Phys. Chem. C **115**, 2514 (2011).
- ⁴⁸Y.-H. Yeh, H. Nakashima, H. Kurishita, S. Goto, and H. Yoshinaga, Mater. Trans., JIM 31, 284 (1990).