A giant planet undergoing extreme ultraviolet irradiation by its hot massive-star host

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The amount of ultraviolet irradiation and ablation experienced by a planet depends strongly on the temperature of its host star. Of the thousands of extra-solar planets now known, only four giant planets have been found that transit hot, A-type stars (temperatures of 7300– 10,000 K), and none are known to transit even hotter B-type stars. WASP-33 is an A-type star with a temperature of \sim 7430 K, which hosts the hottest known transiting planet¹: the planet is itself as hot as a red dwarf star of type M². The planet displays a large heat differential between its day-side and night-side², and is highly inflated, traits that have been linked to high insolation^{3,4}. However, even at the temperature of WASP-33b's day-side, its atmosphere likely resembles the molecule-dominated atmospheres of other planets, and at the level of ultraviolet irradiation it experiences, its atmosphere is unlikely to be significantly ablated over the lifetime of its star. Here we report observations of the bright star HD 195689, which reveal a close-in (orbital period \sim 1.48 days) transiting giant planet, KELT-9b. At \sim 10,170 K, the host star is at the dividing line between stars of type A and B, and we measure the KELT-9b's day-side temperature to be \sim 4600 K. This is as hot as stars of stellar type K4⁵. The molecules in K stars are entirely dissociated, and thus the primary sources of opacity in the day-side atmosphere of KELT-9b are likely atomic metals. Furthermore, KELT-9b receives \sim 700 times more extreme ultraviolet radiation (wavelengths shorter than 91.2 nanometers) than WASP-33b, leading to a predicted range of mass-loss rates that could leave the planet largely stripped of its envelope during the main-sequence lifetime of the host star⁶.

The first transiting planets were discovered around cool, solar-type stars^{7,8}, primarily because hot stars have few spectral lines and rotate rapidly, making Doppler confirmation of planets more difficult. Only in the past few years have transiting planets been confirmed around hot stars of

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types early-F and A^{9,10}, inspired by the discovery of WASP-33b¹. That discovery demonstrated that it is possible to confirm transiting planets around rapidly rotating hot stars via a combination of relatively low-precision radial-velocity measurements and Doppler tomography. However, even the hottest of these few A-type transiting-planet host stars only reach temperatures of \sim 7500 K. Thus, while transit surveys, in particular *Kepler*¹¹, have extended the census of planets around low-mass stars, our understanding of planets around massive, hot stars remains poor.

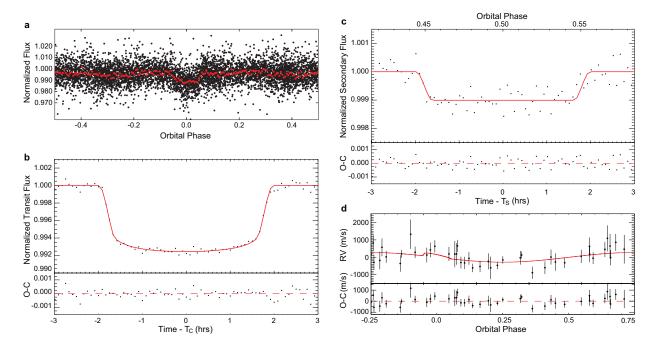
Massive stars cool and spin down as they evolve, enabling precise Doppler measurements. Thus the primary strategy to search for planets around high-mass stars has been surveys of "retired A-stars"¹², high-mass stars that have already evolved into subgiant and giant stars. These stars have revealed a paucity of short-period giant planets relative to sun-like main-sequence stars¹². One interpretation is that the initial planet population of high-mass stars is similar to that seen in unevolved sun-like stars, but that the short-period planets are subsequently engulfed during the evolution of their parent stars or ablated by the intense irradiation of their host stars while they are still hot⁶. Another interpretation is that these stars actually have masses similar to the Sun¹³, implying that the paucity of short-period planets among the retired A-stars is indeed a signature of planet engulfment¹⁴, as sun-like stars do not emit strong ultraviolet radiation with which to ablate their planets⁶.

It is therefore critical to assay the population of short-period planets around *bona fide* highmass stars while they are still on the main sequence, and then to map the evolution of these planets through to their later evolutionary phases. Although there have been radial-velocity surveys targeting unevolved high-mass stars 15,16 , there are still no known transiting planets around unevolved stars more massive than 2 M_{\odot} that produce high levels of extreme ultraviolet irradiation.

The Kilodegree Extremely Little Telescope (KELT) is an all-sky survey for planets transiting bright (visual magnitude 8–11) stars^{17,18}. HD 195689 (hereafter KELT-9), exhibited repeating transit-like events of \sim 0.6% depth with a period of $P \sim$ 1.48 d (Figure 1), and was selected as a candidate transiting planet (see the Methods). KELT-9's basic properties (Table 1) include a high effective temperature and rapid rotation. Following the approach that led to the discovery of WASP-33b, we obtained follow-up observations (see Figure 1, Figure 2, and the Methods) that ultimately confirmed KELT-9b as a transiting planet.

KELT-9 is a hot ($T_{\rm eff} \simeq 10,170$ K), massive ($M_* \simeq 2.5~M_{\odot}$) star of spectral type B9.5–A0, with a relatively young age of 300 million years (Methods), comfortably in its 500-million-year main-sequence phase of evolution; it has evidently not yet begun its evolution toward becoming a "retired A star". Indeed, from comparison with other known planet-hosting stars (Fig. 3), it is clear that KELT-9 is a likely progenitor of at least a subset of the putative "retired A star" hosts of planets detected by radial-velocity surveys. KELT-9, and other A-star transiting planet hosts, thereby provides an important missing link between these samples of planets, and planets detected in more traditional radial-velocity surveys of sun-like stars.

KELT-9 is only the fifth A-type star known to host a transiting giant companion, and is by a significant margin the hottest, most massive, and most luminous known transiting giant planet host. The host star also has the brightest V-band magnitude of any transiting hot Jupiter host, being



KELT-9 discovery and confirmation data. Shown are the discovery light curve, follow-up primary transit and secondary eclipse light curves, and reflex Doppler signal due to the transiting companion. (a) Field 11 of the northern KELT survey was observed 6001 times from UT 2007 May 30 to UT 2013 June 14, following the same procedures as described for the first KELT planet²⁹ The black points show the unbinned data folded at the period of the planet (P = 1.4810932 days), while the red points show these points binned in phase using a bin size of 0.05 in phase. (Methods). (b) Combined, binned follow-up light curves of the primary transit, where TC is the time of the center of the transit. The transit shape is consistent with an opaque, circular planet occulting a star with the limb darkening expected for KELT-9's spectral type of B9.5-A0 The black points show the average of all follow-up light curves, combined in 5-minute bins. The combined best-fit models binned the same way are shown as a solid red line. (c) Combined, binned follow-up light curves of the secondary eclipse. The magnitude of the signal gives a temperature for the day-side of the planet of \sim 4600 K. This is consistent with the range of expected equilibrium temperatures of the planet and suggests poor heat redistribution to the night side of the planet. The data points and red curve are binned in the same manner as in b. TS is the time of the center of the secondary eclipse. (d) Doppler reflex curve of the primary star. Only the data that were not used in the Doppler Tomography analysis (see Fig. 2) are shown and were used to measure the Doppler reflex curve. The best-fit model of the Doppler reflex curve is shown as the red line. Note the weak Rossiter-McLaughlin signal near phase zero, which is a prediction based on the Doppler Tomographic signal exhibited in Figure 2. A periodogram of this data alone yields a significant peak with an emphemeris that matches that of the photometric transit data to within $\lesssim 10^{-4}$ days, thereby confirming the reflex signal is due to the transiting planet and verifying the planet mass estimate.

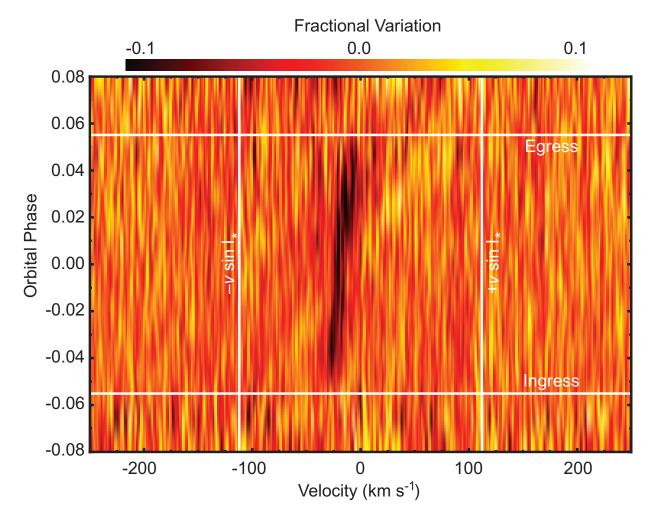


Figure 2: Combined Doppler tomographic measurements from three separate transits. All three transits (see Methods) clearly show a "Doppler shadow" at the time of transit, with the velocity width as expected given the spectroscopically-measured $v\sin I_*$ and the photometrically-measured transit depth and impact parameter, thus confirming the reality of the Doppler tomographic signal and that the planet orbits KELT-9. The color coding shows the fractional variation in the spectroscopic signal from the null hypothesis of no shadow due to a transiting planet. Darker regions indicate the Doppler shadow as the planet crosses the face of the host star. The horizontal line indicate the expected beginning (bottom) and end (top) of the transit, and the vertical lines indicate the negative and positive width of the spectral lines due to broadening caused by the project rotation speed of the star (see Methods). The projected path of the planet across the face of the star is nearly coincident with the stellar rotation axis, with an impact parameter in units of the stellar radius of ~ 0.2 and a projected spin-orbit misalignment of ~ -85 degrees, indicating that the planet is on a nearly polar orbit. The host star's rapid rotation and resulting oblateness implies that the orbit of the planet is likely to exhibit precession, which should be detectable by the year 2022^{30} .

slightly brighter in V than HD 209458b^{7,8}.

Given the high stellar luminosity and close orbit, the planet receives a large stellar insolation flux (Table 1). As a result, it has an extremely high equilibrium temperature, assuming zero albedo and perfect heat redistribution, of \sim 4050 K. This is as hot as a late K-type star⁵, and thus we expected a large thermal emission signal, which we easily confirmed with our z'-band detection of the secondary eclipse with a depth of $\sim 0.1\%$ (Figure 1). This measurement implies an even hotter day-side temperature of $\sim 4600 \pm 150$ K, likely indicating poor redistribution of energy to the night side of the planet and a temperature closer to that of a mid-K star. The planet is also extremely inflated relative to theoretical models, with a radius of $\sim 1.9~R_{\rm J}$. Poor redistribution of heat and radius inflation (both of which are also observed in WASP-33b^{1,2}), have been linked to high stellar insolation^{3,4}, although the exact physical mechanisms remain uncertain.

Thus, although other transiting planets have been found around A-type stars, and indeed some (such as WASP-33b) have been discovered with the temperatures of low-mass stars, the KELT-9 planet and host star are hotter by ~ 1000 K and ~ 2500 K, respectively, than any other known transiting gas-giant system. Consequently, one expects all of the opacity sources on the day-side to be atomic, as in a K-type star. In contrast, all other known transiting planets, which have day-side temperatures of < 3850 K⁵, are cool enough to contain molecular species. Furthermore, the very high flux of extreme ultraviolet radiation (wavelengths shorter than 91.2 nanometers) from KELT-9, ~ 700 times higher than WASP-33, may lead to unique photochemistry in the planet atmosphere¹⁹. This places KELT-9b in a qualitatively new regime of planetary atmospheres, and

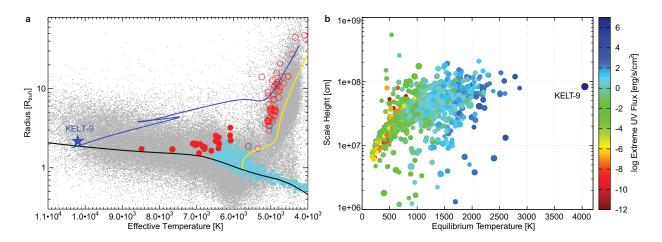


Figure 3: The extreme properties of the KELT-9 system in context and the prospects for follow-up. (a) Hertzsprung-Russell (R_* vs. $T_{\rm eff}$) diagram of hosts of known planets detected by the radial-velocity (open circles) and transit methods (filled circles), as well as nearby stars in the Hipparcos catalog (grey points). Only planet-hosting systems with $V \leq 10$ are shown for clarity. Cyan symbols are low-mass planet hosts with $M < 1.5~M_{\odot}$, red symbols indicate massive planet hosts with $M_* \geq 1.5~M_{\odot}$. The yellow line shows the evolutionary trajectory for a solar analog ($M_* = M_{\odot}$ and solar metallicity), whereas the blue track shows the evolutionary trajectory for an analog of KELT-9 ($M_* = 2.5~M_{\odot}$ and solar metallicity). We also show the zero-age main sequence for solar-metallicity stars from the YY isochrones (black curve). KELT-9 is hotter than any other known transiting planet host by $\sim 1500~{\rm K}$. (b) Atmospheric scale height versus equilibrium temperature for known transiting planets with measured masses. Color represents the amount of incident extreme ultraviolet ($\lambda \leq 91.2$ nanometers) flux the planet receives from its parent star, and the symbol size is inversely proportional to the V magnitude of the host. KELT-9b is hotter than any other known transiting gas giant planet by $\sim 1000~{\rm K}$ and receives $\sim 700~{\rm times}$ more extreme ultraviolet flux.

makes characterization of the atmosphere of KELT-9b particularly compelling.

Fortunately, the brightness of KELT-9 and the extreme properties of its transiting planet make the prospects for detailed characterization of this system promising. Observations using ground-based facilities, Spitzer, the $Hubble\ Space\ Telescope\ (HST)$, and ultimately the $James\ Webb\ Space\ Telescope$, will allow for the measurement of the phase-resolved spectrum of its thermal emission from the far-optical through the infrared ($\sim 30\ \mu m$). The low surface gravity of KELT-9b combined with the high temperature lead to one of the largest atmospheric scale heights of

any known transiting planet (Figure 3). The expected few-percent variations in the transmission spectrum during the primary transit should be easily detectable.

The future evolution of the KELT-9 system is uncertain but certainly interesting. The high ultraviolet flux impinging on KELT-9b likely means that its atmosphere is being significantly ablated, with an estimated mass loss rate⁶ of $\sim 10^{10}$ – 10^{13} g s⁻¹. At the upper end of these rates, the planet may be completely stripped of its outer envelope in < 600 Myr, roughly the time scale for the host to evolve from the main-sequence to the base of the red giant branch (see Methods). These estimates are so uncertain because the physics of planet evaporation is extraordinarily complex, particularly due to the unknown magnitude of the host star activity, and thus very high-energy (extreme ultraviolet and X-ray) non-thermal radiation and stellar wind. In any event, even the lower end of mass-loss rates we estimate should be easily detectable and measurable with HST observations.

On the other hand, as KELT-9 eventually exhausts its core hydrogen supply in \sim 200 million years, it will grow from its current radius of \sim 2.4 R_{\odot} to a radius of \sim 5 R_{\odot} , while simultaneously cooling to $T_{\rm eff}$ \sim 8000 K. Soon after, it will rapidly traverse the 'subgiant branch' whereby it will cool to \sim 5000 K and expand to \sim 8 R_{\odot} . As it reaches the base of the red giant branch, the stellar surface of KELT-9 will encroach upon the orbit of KELT-9b. Exactly what will happen to the star and planet at this point is far from clear. If the mass loss from ablation is lower than estimated above, the planet may remain intact as a gas giant, and it will be engulfed by its host star, perhaps leading to a bright transient event²⁰, and an unusually rapidly rotating red giant star

with enriched lithium provided by the dissolved planetary companion²¹. On the other hand, if the planet possesses a rocky core and is fully ablated before this point, this could imply the existence a population of close-in, super-Earth remnant cores orbiting subgiant stars, a prediction that could be testable with the upcoming Transiting Exoplanet Survey Satellite (TESS) mission.

More detailed theoretical studies are needed to provide a clearer picture of the future evolution of the KELT-9 system and its analogs, and provide testable predictions. Further follow-up observations using ground and space-based telescopes will test models of heat redistribution, radius inflation, unusual photochemistry, and rapid ablation of planetary atmospheres. The leap from WASP-33b to KELT-9b should invigorate further exploration of the planet population of even higher-mass host stars, complementing efforts to discover planets orbiting ever lower-mass host stars²². The KELT-9 system provides an important benchmark system for understanding the nature of planetary systems around massive stars, from birth to death.

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Author Contributions B.S.G. led the process from initial candidate selection to final planet confirmation. B.S.G. and K.G.S. wrote the majority of the main manuscript and contributed to the analysis of the results. K.A.C. principally coordinated the assemblage of the final results and production of the Methods. T.G.B. and G.Z. lead the Doppler Tomographic analysis, and also contributed to the final modeling and interpretation. J.D.E., R.J.S., D.J.S., L.A.B., J.P., J.E.R., K.D.C., M.C.J., and M.P. all provided critical insight, analysis, or interpretation of the system. G.D., V, B., S.C.N., M.T.D., T.E., C.G., H.J.-C., D.H.K, A.F., J.G., A.I., J.F.K., M.M., K.M., N.N., T.E.O., P.A.R., G.S., D.C.S., R.R.Y., R.Z., B.J.F., and A.H., all provided photometric or radial velocity data that was important for the interpretation of the system. J.C. and E.J.G. provided the observations and analysis of the Keck AO data. E.L.N.J., D.J.J., D.B., I.A.C., D.L.D., G.A.E., A.G., M.D.J., R.B.K, J.L.-B., M.B.L, J.M. K.K.M., R.W.P, H.R. D.C.S., C.S., T.G.T, M.T, and P.T. have all been essential for the initiation and successful operation of the KELT-North and KELT-South surveys. All of the authors have read both the main manuscript and supplemental material and concur with the conclusions therein.

Author Information The authors declare that they have no competing financial interests. Correspondence and requests for materials should be addressed to B.S.G. (email: gaudi.1@osu.edu).

Stellar Parameters:		
	M (35)	2 = 2+0 25
$M_* \dots \dots$	Mass (M_{\odot})	$2.52^{+0.25}_{-0.20}$
R_*	Radius (R_{\odot})	$2.362^{+0.075}_{-0.063}$
$L_* \dots \dots$	Luminosity (L_{\odot})	$2.362_{-0.063}^{+0.075}\atop{-0.063}\atop{53_{-10}^{+13}}$
$ ho_*\dots\dots$	Density (cgs)	0.2702 ± 0.0029
$\log g_* \dots \dots$	Surface gravity (cgs)	4.093 ± 0.014
$T_{\mathrm{eff}} \dots \dots$	Effective temperature (K)	10170 ± 450
[Fe/H]	Metallicity	-0.03 ± 0.20
$v\sin I_* \dots \dots$	Rotational velocity (km/s)	111.4 ± 1.3
λ	Spin-orbit alignment (degrees).	-84.8 ± 1.4
Planetary Parameters:		
P	Period (days)	1.4811235 ± 0.0000011
a	Semi-major axis (AU)	$0.03462^{+0.00110}_{-0.00093}$
$M_P \dots \dots$	Mass ($M_{ m J})$	2.88 ± 0.84
$R_P \dots \dots$	Radius ($R_{ m J})$	$1.891^{+0.061}_{-0.053}$
$ ho_P \ldots \ldots$	Density (cgs)	0.53 ± 0.15
$\log g_P \dots \dots$	Surface gravity	$3.30^{+0.11}_{-0.15}$
$T_{eq} \dots \dots$	Equilibrium temperature (K)	4050 ± 180
$\langle F \rangle \dots \langle F \rangle$	Incident flux (10 9 erg s $^{-1}$ cm $^{-2}$)	$61.1_{-9.8}^{+11.0}$

Table 1: Median values and 68% confidence intervals for the physical properties of the KELT-9 system from a global fit to the light curves, radial velocities, and Doppler tomography data, using a custom version of the EXOFAST transit fitting code²³ (Methods). We use the Yonsei-Yale (YY) isochrones²⁴ as well as empirically-calibrated stellar relations²⁵ as constraints. In addition, KELT-9 has a measured parallax from $Hipparcos^{26}$ and $Gaia^{27}$, thereby allowing us to estimate the stellar radius empirically²⁸ from its distance, effective temperature, bolometric (wavelength-integrated) flux, and interstellar extinction (Methods). We adopt the values derived using constraints from the YY isochrones and Hipparcos-derived radius as our fiducial system parameters.

Methods

KELT-North Observations and Photometry

KELT-North field 11 is centered on $\alpha=19^h\ 27^m\ 00^s$, $\delta=31^\circ\ 39'\ 56''.16$ (J2000) and was observed 6001 times from UT 2007 May 30 to UT 2013 June 14. Following the standard KELT candidate selection strategy²⁹, we reduced the data and extracted $\sim 116,000$ light curves from the east orientation and $\sim 138,000$ light curves from the west. The combined east and west light curves that passed the reduced proper motion cut³¹ were searched for transiting exoplanet candidates. One bright ($V\sim 7.6$) candidate, KC11C043952 (HD 195689, TYC 3157-638-1, 2MASS J20312634+3956196) located at $\alpha=20^h\ 31^m\ 26^\circ\!35401$, $\delta=39^\circ\ 56'\ 19''.7744$ (J2000), robustly passed our selection criteria²⁹ making it a top candidate. Hereafter, we refer to the candidate host star as "KELT-9" and the candidate planet as "KELT-9b". KELT-9 is also located in KELT-North field 12, which is centered on $\alpha=21^h\ 22^m\ 52^\circ\!\!8$, $\delta=31^\circ\ 39'\ 56''.16$ (J2000). The field was observed 5,700 times from UT 2007 June 08 until UT 2013 June 14. Although not originally used to select KELT-9 as a candidate, the signal in the light curve from field 12 bolstered KELT-9 as a strong candidate.

The KELT-North field 11 phased KELT-9 discovery light curve is shown in the top panel of Fig. 1 in the main manuscript. The field 11 light curve shows apparent out-of-transit variations (OOTVs), while these are absent in the field 12 light curve. We find that the source of the field 11 OOTVs is due to variability, possibly caused by saturation, of a star a few arc-minutes southeast of KELT-9 that is blended in the field 11 KELT-9 aperture. Given the magnitude of the variability, we suspect the contaminating source is the bright object HD 195728, rather than one of the fainter

neighbors. The variation is long-term and not periodic, so it tends to inject power at and near aliases of our sampling rate. The variability from the neighbor does not affect the field 12 light curve because the point spread function of KELT-9 in field 12 is smaller and is elongated in a different direction than in field 11. As a result, KELT-9 is minimally blended with the bright neighbor HD 195728, as well as the other faint neighbors near HD 195728.

In short, the apparent out-of-transit variations (OOTVs) seen in the field 11 light curve are not due to intrinsic variability of KELT-9 itself, and thus do not lead us to question its reality as a bona fide planet-host candidate.

Following the discovery of the primary transit signal, we carried out an intensive photometric followup campaign through which we obtained a total of 17 primary and 7 secondary transit light curves. These will be described in a forthcoming paper (Collins et al. 2017, in preparation), and are used in the global fit of the system presented below. The combined, binned follow-up light curve is shown in Figure 1 of the main manuscript in order to highlight the statistical power of this follow-up dataset. We only fit to the individual data; the binned data are shown only for illustration.

Spectroscopic Follow-up

To constrain the mass and enable eventual Doppler tomographic (DT) detection of KELT-9b, we obtained a total of 115 spectroscopic observations of the host star with the Tillinghast Reflector Echelle Spectrograph (TRES) on the 1.5 m telescope at the Fred Lawrence Whipple Observatory, Arizona, USA. Each spectrum delivered by TRES has a spectroscopic resolution of $\lambda/\Delta\lambda=41000$ over the wavelength range of $3900-9100\,\text{Å}$ over 51 echelle orders. This includes 40 observations

covering the entire orbital phase to constrain the mass of the planet, and 75 observations made intransit over three epochs to perform the tomographic line profile analysis. We measured the relative radial velocity from 104 of the observations (see Extended Data Table) and used a total of 43 out-of-transit RVs (40 plus one out-of-transit RV from each of the spectroscopic transit observations) to constrain the planet's orbit and mass. The phased radial velocities are displayed in Figure 1 of the main article.

Three spectroscopic transits of KELT-9 were observed with TRES on 2014-11-15, 2015-11-06, and 2016-06-12. The line broadening kernel is derived from each spectrum via a least-squares deconvolution^{1,32,33}. The planetary shadow is seen crossing the stellar surface on all three nights, as shown in Extended Data Figure 3. In addition, the rotational profiles allowed us to accurately determine a rotational velocity of $v \sin I_* = 111.4 \pm 1.3 \, \mathrm{km \, s^{-1}}$.

Host Star Properties

Extended Data Table lists various properties and measurements of KELT-9 collected from the literature and derived in this work. The data from the literature include the four monochromatic near-UV fluxes from the Catalog of Stellar UV Fluxes³⁴, UBV photometry³⁵, optical fluxes in the $B_{\rm T}$ and $V_{\rm T}$ passbands from the Tycho-2 catalog³⁶, $I_{\rm C}$ from the TASS catalog³⁷, near-infrared (IR) fluxes in the J, H and $K_{\rm S}$ passbands from the 2MASS Point Source Catalog^{38,39}, near- and mid-IR fluxes in four WISE passbands^{40,41}, distances from Hipparcos⁴² and Gaia²⁷, and proper motions from the NOMAD catalog^{43,44}.

SED Analysis

We construct an empirical, broad-band spectral energy distribution (SED) of KELT-9, shown in Extended Data Figure 1. We use the 17 photometric measurements from the literature discussed in Section and shown in Extended Data Table . In total, the observed SED spans the wavelength range 0.16– $22 \,\mu\text{m}$. We fit this observed SED to Kurucz stellar atmosphere models⁴⁵. For simplicity we adopted a fixed $\log g_{\star} = 4.0$, based on the light curve transit analysis. The fit parameters were thus the effective temperature (T_{eff}) , the metallicity ([Fe/H]), the extinction (A_V) , and the overall flux normalization. The maximum permitted extinction was set to $A_V = 0.30$ based on the total line-of-sight extinction in the direction of KELT-9 from Galactic dust maps⁴⁶.

Stellar Parameters from SED

The best fit model has a reduced χ^2 of 2.56 for 13 degrees of freedom (17 flux measurements, 4 fit parameters). We find $A_{\rm V}=0.09\pm0.05$, $[{\rm Fe/H}]=0.0\pm0.2$, and $T_{\rm eff}=9560\pm550$ K. We note that the quoted statistical uncertainties on $A_{\rm V}$ and $T_{\rm eff}$ are likely to be slightly underestimated because we have not accounted for the uncertainty in the value of $\log g_{\star}$ used to derive the model SED, although this parameter generally does not strongly affect the overall shape of the SED, and moreover $\log g_{\star}$ is strongly constrained from the light curve transit analysis.

We can integrate the best-fit SED to obtain the (unextincted) bolometric flux at Earth, $F_{\rm bol}=3.18(\pm0.09)\times10^{-8}~{\rm erg~s^{-1}~cm^{-2}}$. Together with the best-fit $T_{\rm eff}$ and the distance newly provided by the *Gaia* parallax, we obtain a direct constraint on the stellar radius of $R_{\star}=2.37\pm0.35~R_{\odot}$. From the distance provided by the Hipparcos parallax, we obtain a direct constraint²⁸ on the stellar

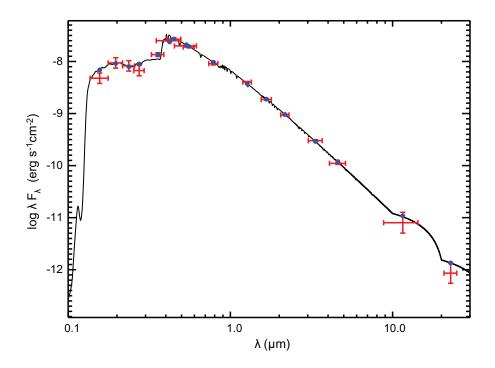


Figure 1: KELT-9 spectral energy distribution (SED). Crosses represent the measured fluxes, with vertical bars representing the measurement uncertainties and the horizontal bars representing the width of the bandpass. The blue dots are the predicted passband-integrated fluxes of the best-fit theoretical SED corresponding to our observed photometric bands. The black curve represents the best-fit Kurucz stellar atmosphere⁴⁵.

radius of $R_{\star}=2.17\pm0.33~R_{\odot}$. As noted in the footnote of Extended Data Table , we generally believe the Hipparcos-derived distance and stellar radius to currently be more reliable, and use this for most of the analysis in this paper.

Stellar Models and Age

With $T_{\rm eff}$ and $\log g_{\star}$, and an estimated stellar mass from the global analysis (see below), we can place the KELT-9 system in the Hertzsprung-Russell diagram for comparison with theoretical stellar evolutionary models (Extended Data Figure 2). From the values obtained with the initial SED fit, we infer a system age of ≈ 0.4 Gyr; the final age estimate using the final global fit parameters is ≈ 0.3 Gyr. Thus, it is clear that the KELT-9 system is nearly unevolved from the zero-age main-sequence, and in any event is at an early stage of evolution well before the "blue hook" transition to the subgiant and eventual red giant evolutionary phases.

Global System Fit

We determined the physical and orbital parameters of the KELT-9 system by jointly fitting 17 primary and 7 secondary light curves, 43 TRES out-of-transit RVs, and 3 Doppler tomographic data sets (see the section Doppler tomographic model below). To perform the global fit, we used MULTI-EXOFAST (MULTIFAST hereafter), which is a custom version of the public software package EXOFAST²³. MULTIFAST first performs an AMOEBA⁴⁷ best fit to each of the RV and light curve data sets individually to determine uncertainty scaling factors. The uncertainties are scaled such that the probability that the χ^2 for a data set is larger than the value we achieved, $P(>\chi^2)$, is 0.5, to ensure the resulting parameter uncertainties are roughly accurate. The re-

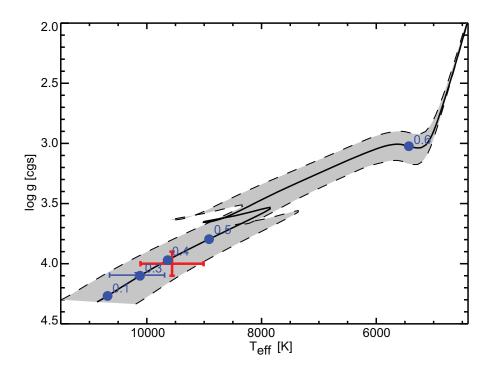


Figure 2: Evolution of the KELT-9 system in the Hertzsprung-Russell diagram. The red cross represents the system parameters from the initial SED fit, the blue cross represents the final system parameters from the global fit. The black curve represents the theoretical evolutionary track for a star with the mass and metallicity of KELT-9, and the grey swath represents the uncertainty on that track based on the uncertainties in mass and metallicity. Nominal ages in Gyr are shown as blue dots. When KELT-9 evolves to the base of the giant branch in 200-300 million years, it will encroach upon the orbit of KELT-9b. The fate of the system at that point is highly uncertain $^{20,21,56-58}$.

sulting RV uncertainty scaling factor is 1.23. The DT uncertainties were scaled based on the χ^2 of the out-of-transit data relative to the median value. The uncertainties of the UT 2014-11-15, UT 2015-11-06, and UT 2016-06-12 DT observations were scaled by 0.79, 0.78, and 0.79, respectively. Finally, MULTIFAST performs a joint AMOEBA model fit to all of the datasets and executes Markov Chain Monte Carlo (MCMC), starting at the global best fit values, to determine the median and 68% confidence intervals for each of the physical and orbital parameters. Siverd et al.²⁹ provide a more detailed description of MULTIFAST, except the Doppler tomographic model implementation, which was newly implemented as part of this work.

Doppler Tomographic Model

To model the Doppler tomographic signal, we construct and integrate our own models into the MULTIFAST fitting process⁴⁸. We treat the Doppler shadow of the planet as a combination of three different broadening mechanisms which we account for consecutively. Our shadow model begins as a Gaussian profile, with a standard deviation equal to the mean inherent spectral line width (in velocity space) of an equivalent non-rotating star. Second, we convolve this base Gaussian profile with a second Gaussian of width $\sigma = c/R$ to account for the finite spectral resolution of the TRES spectrograph ($R \approx 44,000$). Third, and finally, we convolve the resulting shadow model with a normalized rotational broadening kernel.

For our rotational broadening kernel we use the kernel given by Equation 18.14 in Gray⁴⁹. For simplicity we assume that the linear limb-darkening coefficient (ϵ) for the kernel is equal to zero. We set the velocity width of the kernel equal to (R_p/R_*) $v \sin I_*$, to represent the fraction of

the stellar rotational surface actually obscured by the planet.

We normalize the resulting shadow model using the depth of a V-band transit light curve for the system at each observation time, so that the integrated light obscured by the shadow matches the actual depth of the transit. To calculate the trajectory of the shadow in time-velocity space, we follow the method given by Equations 7, 8, and 10 in Collier Cameron et al. 50

Besides the physical parameters already included in the transit model, our Doppler tomography model therefore has three free parameters: the rotation velocity of the host star ($v \sin I_*$), the inherent line width of the non-rotating stellar spectrum (σ_0), and the spin-orbit angle (λ). Since KELT-9 has $v \sin I_* = 111.4 \pm 1.3 \, \mathrm{km \, s^{-1}}$, σ_0 is essentially negligible, so we use a prior on $\sigma_0 = 10.0 \pm 1.0 \, \mathrm{m \, s^{-1}}$ (i.e. a value near zero that avoids significant numbers of negative σ_0 MCMC trials) to improve convergence of the MCMC chains. To evaluate the goodness-of-fit for a particular set of model parameters, we use the $\Delta \chi^2$ between the predicted shadow model and the observed Doppler tomographic observations. Since the observations themselves are super-sampled below the instrumental resolution of the spectrograph during the data reduction process, we divide the resulting $\Delta \chi^2$ value by a factor of $(c/R)/\Delta v$, where c is the speed of light, R is the instrumental spectral resolution, and Δv is the velocity interval between the individual points in the super-sampled spectra. This reduction of the $\Delta \chi^2$ accounts for the fact that due to the instrumental spectral resolution, adjacent points in the super-sampled spectra were not truly independent.

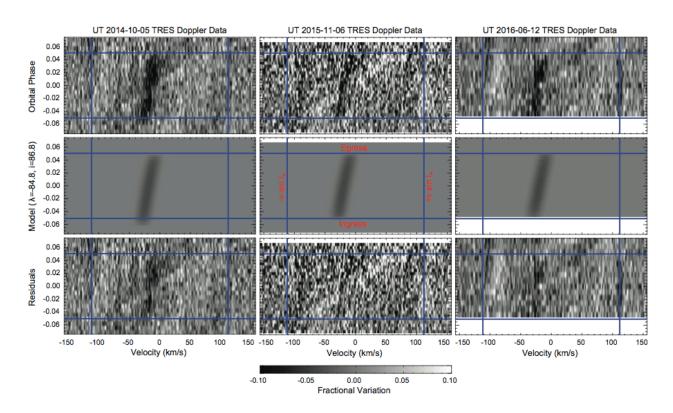


Figure 3: Doppler tomographic line profile plots. The data, models, and residuals from all three nights where KELT-9 was observed during transit are shown in three columns. In each case, the top plot shows the spectroscopic date, the middle panels show the derived model, and the bottom panels show the residuals. In each panel, the vertical blue lines denote the width of the convolution kernel (i.e., $v \sin I_*$), and the horizontal blue lines show the duration of the transit. Time increases vertically for each panel. The apparent extension of the Doppler shadow before ingress on UT 20014-10-05 is an artifact of uneven and widely-spaced sampling in time. The greyscale shows the fractional variation in the spectroscopic signal from the null hypothesis of no shadow due to a transiting planet. Darker regions indicate the Doppler shadow as the planet crosses the face of the host star. Note the transit is nearly coincident with the projected stellar spin axis, implying a nearly polar orbit for the planet.

Global Model Results

We adopt a fiducial model (Model 1) with YY constraints, a fixed circular orbit, a fixed RV slope $\dot{\gamma}=0$, and a Hipparcos-based R_{\star} prior, and compare the results to those from ten other global models that systematically explore the results of differing constraints. The posterior median parameter values and 68% confidence intervals are shown in Extended Data Table for the initial TTV and final ten global model fits. The KELT-9 fiducial model indicates the system has a host star with mass $M_{\star}=2.52\,M_{\odot}$, radius $R_{\star}=2.362\,R_{\odot}$, and effective temperature $T_{\rm eff}=10,170$ K, and an extremely hot planet with $T_{\rm eq}=4050$ K, mass $M_{\rm P}=2.88M_{\rm J}$, and radius $R_{\rm P}=1.891\,M_{\rm J}$.

The initial TTV model (Model 0) parameter median values are well within 1σ of the fiducial model, and the uncertainties are essentially identical to those of the fiducial model results, except the ephemeris parameter uncertainties (T_C and P) are $\sim 25\%$ larger due to the additional seven secondary transits in the fiducial model and the linear ephemeris constraint.

We also explored several other models involving different choices of observational constraints and free-fit parameters, which will be described in a forthcoming paper (Collins et al. 2017, in preparation) but are summarized in Extended Data Table.

In summary, we find that all combinations of stellar constraints result in system parameter values that are within 1σ , and in almost all cases, well within 1σ of the fiducial model.

False Positive Analysis

KELT-9 was a sufficiently unusual system that we were especially vigilant about ruling out false positive scenarios. We have many lines of evidence that rule out essentially all viable false positive scenarios. First, false positives around rapidly rotating stars can be ruled out via relatively imprecise RV measurements, high precision light curves, and a positive Rossiter-McLaughlin^{51,52} (RM) or DT detection^{1,53}. An upper limit on the Doppler signal with relatively imprecise RV precisions of a few hundred m s⁻¹ can rule out low-mass stars and brown dwarfs as the occulter of the primary star. Precise follow-up light curves, along with the $v \sin I_*$ measured from the spectra, can be used to predict the magnitude, impact parameter, and duration of the DT signal (although not the direction λ). Thus a measurement of the DT signal that is consistent with the light curve effectively confirms the planet interpretation.

In the case of KELT-9, our first spectroscopic measurement during transit yielded a very weak RM signal, which was sufficiently noisy that we considered it could simply be due to the relatively large uncertainties of the RVs. This lack of a RM signal would be surprising for a transiting planet given the transit depth and large $v \sin I_*$ of the star. We thus originally concluded that the system was likely a false positive. However, we decided to proceed with a full DT analysis, which showed tentative evidence for a planet shadow that was nearly coincident with the projected stellar rotation axis, thus possibly explaining the small RM signal. Improved reduction methods and two additional DT measurements definitively showed that the planet does indeed transit nearly along the projected stellar rotation axis, as shown in Extended Data Figure 3.

Second, we were also concerned about the lack of a definitive measurement of the Doppler reflex signal. Our initial fits to a subset of the data presented here provided only upper limits to K, and thus we were originally not able to measure the planet mass. However, after acquiring additional data, we were ultimately able to measure K to roughly $\sim 30\%$ precision. Indeed, a Lomb-Scargle^{54,55} periodogram of the out-of-transit RV data yielded several significant signals. The two signals with the highest power are at long periods of ~ 100 and ~ 200 days (depending on the exact choice of datasets that are used), which are likely aliases of each other, and may be due to another planet, intrinsic long-term stellar activity, or may simply be due to systematic errors in the RVs. However, the signal with the next highest power has a ephemeris of $T_0 = 2456969.7901$ and P = 1.4811228. This period is consistent with the final period we derive from our follow-up light curves to within $\sim 10^{-5}$ d, and projecting the ephemeris from our follow-up light curves forward yields a value of T_0 that is consistent with that from the RV periodogram to within $\sim 10^{-4}$ d. We are therefore confident that the reflex RV signal is real and due to the transiting planet.

Finally, there are several other pieces of evidence that support the planet hypothesis. (1) The stellar density we infer from the global fit to the light curve (which is essentially a direct observable) is consistent with what one would expect for an unevolved A0 star with the spectroscopically-measured temperature. (2) The limb darkening in the redder bandpasses are noticeably smaller by eye than planet transits of a cooler star, as one would expect for a hot star and bandpasses near the Rayleigh-Jeans part of the SED. Although we do not fit for the limb darkening, the fact that the light curves spanning from U to z are well-fit by the smaller limb darkening coefficients that are predicted for a star of this temperature and surface gravity (as compared to, e.g., those expected

for a solar-type star) provides quantitative support for this qualitative conclusion. (3) Adaptive optics observations (Collins et al. 2017, in preparation) do not reveal a blended stellar companion that could cause a false positive, down to the mass of $\sim 0.14~M_{\odot}$ with a projected separation of $\gtrsim 200~{\rm AU}$. (4) We detect a secondary eclipse in the z band with a depth that is consistent with what one would expect given the amount of stellar irradiation that the planet receives. (5) The basic consistency between the model fits using various constraints (YY, Torres, and Hipparcos and Gaia-inferred radii) also provides support for our interpretation.

We therefore conclude that, despite the very unusual and extreme nature of the system, all available data are consistent with the interpretation that KELT-9 is being transited by an extremely irradiated, highly inflated planet on a near-polar orbit of only ~ 1.5 d.

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Data Availability The data that support the findings of this study are available from the corresponding author upon reasonable request.

BJD _{TDB}	RV*	σ_{RV}^{\dagger}	Orbit [‡]	DT§	BJD _{TDB}	RV*	σ_{RV}^{\dagger}	Orbit [‡]	DT§
			Y						Y
2456914.630558	1284.62	370.91		_	2457332.565948	919.48	472.04	-	
2456931.621504	0.00	168.62	Y	_	2457332.573447	1089.69	326.53	-	Y
2456932.830099	656.46	328.19	Y	_	2457332.583568	1101.27	248.54	_	Y
2456933.768611	1322.24	394.09	Y	_	2457332.591091	20.88	492.65	-	Y
2456934.720143	595.47	168.62	Υ	Ÿ	2457332.598573	369.91	520.93	_	Y
2456935.607094	216.76	312.86	-	Y	2457332.606056	-35.51	458.92	_	Y
2456935.642110	264.80	328.13	_		2457332.613468	646.39	483.26		
2456935.651924	352.14	416.06	-	Y	2457332.620922	1.55	327.82	_	Y
2456935.661420	398.38	310.71	_	Y	2457332.628363	762.52	542.15	-	Y
2456935.670957	550.61	302.90	_	Y	2457332.635932	130.36	566.49	-	Y
2456935.680453	152.78	341.89	_	Y	2457332.643363	1236.65	393.37	-	Y
2456935.690024	508.93	256.70	_	Y	2457332.651204	592.91	712.93	-	
2456935.699561	254.60	417.88	_	Y	2457332.658749	667.61	366.74	_	Y
2456935.709213	441.60	249.04	_	Y	2457332.666156	683.63	500.68	-	Y
2456935.718750	599.89	354.15	_	Y	2457332.673598	817.72	505.93	-	Y
2456935.728345	532.89	354.56	_	Y	2457332.682452	786.98	586.05	-	Y
2456935.738095	486.15	358.01	_	Y	2457332.689940	533.08	627.12	-	Y
2456935.747713	620.59	319.60	-	Y	2457332.697746	1098.64	560.02	-	Y
2456935.757337	466.26	400.36	_	Y	2457332.705547	104.44	591.02	-	Y
2456935.766989	524.75	388.66	_	Y	2457332.713127	373.94	664.04	-	Y
2456935.777376	856.08	344.75	_	Y	2457332.720592	865.83	521.93	_	Y
2456935.786924	958.58	284.88	_	Y	2457332.728277	714.07	520.84	-	Y
2456935.796444	683.37	414.85	-	Y	2457332.735950	673.97	475.68	-	Y
2456935.806102	727.88	528.20	_	Y	2457332.743554	1491.43	714.56	_	Y
2456935.815841	225.39	400.57	_	Y	2457332.751343	237.89	754.52	-	Y
2456935.825494	583.12	460.85	_	Y	2457332.758825	31.25	735.54	_	Y
2456935.835134	660.37	505.61	Y	Υ	2457332.766400	648.09	609.21	Y	Υ
2456936.603281	414.40 595.45	273.39	Y	_	2457345.641623 2457346.603403	176.68	352.38 292.41	Ý	_
2456941.592333 2456942.760407	650.62	344.62	Ý	_		-130.35 402.68	298.04	Ý	_
2456942.760407	-419.72	253.95	Ý	_	2457347.657077 2457348.636906		407.04	Ý	_
		280.81		_		284.17 196.01		Ý	_
2456944.766967 2456945.588041	141.71	285.10	Y	_	2457349.624490 2457350.574309	164.04	407.52 275.50	Ý	_
2456958.583258	441.26	353.95	Ϋ́	_					Ÿ
2456959.711817	478.53 233.56	349.79 372.13	Ý		2457551.802134 2457551.815416	1074.34 1223.97	386.10 407.75	_	Ý
	109.61	382.89	Ý	_	2457551.828883	928.88	365.34	_	Ϋ́
2456960.697436			Ý	_				_	Ϋ́
2456961.622370	150.95	241.25	-	_	2457551.848623	574.42	330.36	_	
2456962.586162 2456963.580571	-144.65 488.55	196.71 276.30	Y	_	2457551.861396 2457551.874632	905.46 1240.58	449.48 402.07	_	Y
2456964.677590	520.05		Ý	_	2457551.888099	1164.19	326.72		Ý
2456965.589717	-60.44	308.90 261.99		_	2457551.866099	736.73	305.39	_	Ý
			Y	_				_	Ϋ́
2456968.611833	-141.05 659.45	526.54		_	2457551.914298	1107.55	266.56 480.33	_	
2456969.591490		187.51	Y	_	2457551.927713	839.50		_	Y
2456970.648671 2456971.642401	1093.42 314.33	414.61 173.67	Y	_	2457551.941145 2457551.959219	1218.15 712.03	444.26 375.84	_	Y
	887.63	342.10	Ý	_	2457551.959219	621.58	178.27		Ϋ́
2456972.715791				_				Y	Ϋ́
2456973.611765	637.78	370.60	Y	_	2457551.987496	1119.25	280.59	Ý	
2456975.592297	751.58	266.34		_	2457561.930813	1039.28	413.36	Ý	_
2456976.659911	917.54 434.85	352.82	Y	_	2457585.790394 2457589.873837	1789.64	685.84 562.82	Ý	_
2456977.664878		317.53		_		1112.21		Ý	_
2456978.665650	1095.99	286.15	Υ	Y	2457589.953527	929.04 1553.14	669.05	Ý	_
2457332.558483	330.89	571.73		T	2457595.783396	1000.14	605.75	r	

Extended Data Table 1: Radial Velocity Measurements of KELT-9 * relative RVs (m s $^{-1}$)

[†]unrescaled relative RV errors (m s⁻¹)

[‡]Y indicates that the RV was included in the RV orbit analysis

 $[\]S Y$ indicates that the spectrum was included in the Doppler tomography analysis

Parameter	Description	Value	Source§
Names		KELT-9	
		HD 195689	
		TYC 3157-638-1	
		2MASS J20312634+3956196	
α _{J2000}	Right Ascension (RA)	20h 31m 26s35401	Hipparcos
$\delta_{\rm J2000}$	Declination (Dec)	39° 56′ 19″7744	Hipparcos
156.5 nm	USST* passband flux (cW m-2 nm-1)	3.05 ± 0.22 × 10 ⁻¹²	USST
196.5 nm	USST passband flux (cW m ⁻² nm ⁻¹)	$4.78 \pm 0.29 \times 10^{-12}$	USST
236.5 nm	USST passband flux (cW m ⁻² nm ⁻¹)	$3.48 \pm 0.16 \times 10^{-12}$	USST
274.0 nm	USST passband flux (cW m ⁻² nm ⁻¹)	$2.44 \pm 0.07 \times 10^{-12}$	USST
Johnson U	magnitude `	7.617 ± (not available)	Mermilliod
Johnson B	magnitude	7.587 ± 0.014	Mermilliod
Johnson V	magnitude	7.550 ± 0.010	Mermilliod
BT	Tycho B _T magnitude	7.634 ± 0.015	Tycho-2
VT	Tycho V _T magnitude	7.585 ± 0.011	Tycho-2
lc J	magnitude	7.570 ± 0.061	TASS
J H	2MASS magnitude	7.458 ± 0.018 7.492 ± 0.021	2MASS 2MASS
п К _S	2MASS magnitude 2MASS magnitude	7.482 ± 0.021	2MASS
WISE1	WISE passband magnitude	7.460 ± 0.026	WISE
WISE2	WISE passband magnitude	7.497 ± 0.02	WISE
WISE3	WISE passband magnitude	7.456 ± 0.045	WISE
WISE4	WISE passband magnitude	7.412 ± 0.218	WISE
Distance†	Hipparcos Distance (pc)	188.7 ± 18.2	Hipparcos
Distance	Gaia Distance (pc)	205.6 ± 18.7	Gaia
μα	Proper Motion in RA (mas yr-1)	17.3 ± 0.4	NOMAD
μδ	Proper Motion in DEC (mas yr-1)	21.8 ± 0.5	NOMAD
RV	Absolute RV (km s ⁻¹)	-20.567 ± 0.1	TRES Spectroscopy
v sin l*	Stellar Rotational Velocity (km s ⁻¹)	111.40 ± 1.27	TRES Spectroscopy
U [‡]	Space motion (km s ⁻¹)	- 19.9 ± 2.5	This work
V	Space motion (km s ⁻¹)	-2.0 ± 0.5	This work
W	Space motion (km s ⁻¹)	5.4 ± 0.4	This work

Extended Data Table 2: Stellar Properties of KELT-9 Obtained from the Literature and This Work *Ultraviolet Sky Survey Telescope

[†]We generally use the Hipparcos parallax as the default value for most the analysis in this paper, because the Hipparcos and Gaia parallaxes have similar fractional precisions, they agree to within their uncertainties, and the Gaia parallaxes are known to have a small systematic error, although this error is small compared to the absolute value of the parallax of KELT-9. Nevertheless, when warranted, we compare results obtained using both parallax determinations.

[‡]U is positive in the direction of the Galactic Center.

[§]See text for references.

Parameter	Description/Units	Model 0 (Initial TTV Run)	Model 1 (adopted)	Model 2	Model 3	Model 4
Global Fit Configuration:						
M _⋆ & R _⋆ Constraints	Torres, Yonsei-Yale, None	Yonsei-Yale	Yonsei-Yale	Yonsei-Yale	Yonsei-Yale	Yonsei-Yale
Orbital Constraint	Circular, Eccentric	Circular	Circular	Circular	Circular	Circular
RV Slope	No Slope, Slope	No Slope	No Slope	No Slope	No Slope	Slope
Distance Prior	Hipparcos, Gaia, None	Hipparcos	Hipparcos	Gaia	None	Hipparcos
TTV Fit	Yes, No	Yes	No	No	No	No
Stellar Parameters:						
M_*	Mass (M_{\odot})	$2.51^{+(25)}_{-(20)}$	$2.52^{+(25)}_{-(20)}$	$2.56^{+(25)}_{-(21)}$	$2.57^{+(26)}_{-(21)}$	$2.53^{+(25)}_{-(20)}$
R_*	Radius (R_{\odot})	$2.361^{+(75)}_{-(63)}$	$2.362^{+(75)}_{-(63)}$	$2.373^{+(76)}_{-(65)}$	$2.376^{+(78)}_{-(67)}$	$2.365^{+(74)}_{-(64)}$
$L_* \dots \dots$	Luminosity (L_{\odot})	$53^{+(13)}_{-(10)}$	53+(13)	55 ⁺⁽¹³⁾ ₋₍₁₁₎	56+(13)	$54^{+(13)}_{-(11)}$
ρ_*	Density (cgs)	0.2690(30)	0.2702(29)	0.2704(28)	0.2704(28)	0.2703(28)
$\log g_{\star}$	Surface gravity (cgs)	4.091(14)	4.093(14)	4.096(14)	4.096(15)	4.094(14)
T_{eff}	Effective temperature (K)	10140(450)	10170(450)	10240(450)	10260(450)	10220(450)
[Fe/H]	Metallicity	-0.03(20)	-0.03(20)	-0.02(20)	-0.01(20)	-0.04(20)
Planetary Parameters:						
$M_P \dots \dots$	Mass (M _J)	2.86(84)	2.88(84)	2.91(84)	2.92(85)	2.75(86)
$R_P \dots \dots$	Radius (R _J)	1.888+(62)	$1.891^{+(61)}_{-(53)}$	$1.896^{+(62)}_{-(54)}$	$1.898^{+(63)}_{-(56)}$	$1.889^{+(61)}_{-(53)}$
ρ_P	Density (cgs)	0.53(15)	0.53(15)	0.53(15)	0.53(15)	0.51(16)
$\log g_{\mathrm{P}}$	Surface gravity	$3.30^{+(11)}_{-(15)}$	$3.30^{+(11)}_{-(15)}$	$3.30^{+(11)}_{-(15)}$	$3.30^{+(11)}_{-(14)}$	$3.28^{+(11)}_{-(16)}$
T_{eg}	Equilibrium temperature (K)	4040(180)	4050(180)	4080(180)	4080(180)	4070(180)
(F)	Incident flux (10 ⁹ erg s ⁻¹ cm ⁻²)	$60.4^{+(11.0)}_{-(9.8)}$	$61.1^{+(11.0)}_{-(9.8)}$	$62.0^{+(12.0)}_{-(10.0)}$	$63.0^{+(12.0)}_{-(10.0)}$	$62.1^{+(12.0)}_{-(10.0)}$
Orbital Parameters:	,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,	-(9.8)	-(9.8)	-(10.0)	-(10.0)	-(10.0)
T_C (BJD _{TDB})	Primary Center Time -2457095	0.68573(19)	0.68572(14)	0.68573(14)	0.68573(14)	0.68573(13)
T_S (BJD _{TDB})	Secondary Center Time -2457094	0.00010(10)	0.94516(14)	0.94517(14)	0.94517(14)	0.94517(13)
P	Period (days)	1.4811233(16)	1.4811235(11)	1.4811235(11)	1.4811235(11)	1.4811235(11)
a	Semi-major axis (AU)	$0.03456^{+(110)}_{-(92)}$	$0.03462^{+(\hat{1}10)}_{-(93)}$	$0.03478^{+(110)}_{-(96)}$	$0.03483^{+(110)}_{-(100)}$	$0.03466^{+(110)}_{-(93)}$
λ	Spin-orbit alignment (degrees)	-84.8(1.4)	-84.8(1.4)	-84.8(1.4)	-84.8(1.4)	-84.7(1.4)
e	Eccentricity	0.0 (fixed)	0.0 (fixed)	0.0 (fixed)	0.0 (fixed)	0.0 (fixed)
ω*	Arg. of periastron (degrees)	90 (fixed)	90 (fixed)	90 (fixed)	90 (fixed)	90 (fixed)
RV Parameters:	rug. or portablion (degrees)	oo (moa)	00 (1100)	oo (iixoo)	oo (moa)	oo (nxoo)
K	RV semi-amplitude (m/s)	275(78)	276(79)	276(79)	276(78)	263(81)
γ _{TRES}	RV offset (m/s)	463(57)	462(57)	462(57)	463(56)	481(59)
γ	RV slope (m/s/day)	0.0 (fixed)	0.0 (fixed)	0.0 (fixed)	0.0 (fixed)	0.37(30)
Transit Parameters:	,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,	(,	(,	()	()	(00)
$R_P/R_\star \dots$	Radius of the planet (R_*)	0.08219(44)	0.08228(43)	0.08212(43)	0.08211(43)	0.08211(43)
a/R*	Semi-major axis (R_*)	3.148(12)	3.153(11)	3.153(11)	3.153(11)	3.153(11)
i	Inclination (degrees)	86.72(27)	86.79(25)	86.79(25)	86.79(25)	86.79(26)
b	Impact parameter	0.180(14)	0.177(14)	0.177(13)	0.177(13)	0.177(14)
$\delta_p \dots \delta_p$	Primary transit depth (10^{-3})	6.755(72)	6.770(72)	6.744(71)	6.743(70)	6.742(71)
δ	Secondary transit depth (10^{-3})		1.006(97)	1.005(96)	1.003(96)	1.004(97)
T_{14}	Total duration (days)	0.16330(48)	0.16316(48)	0.16309(47)	0.16309(47)	0.16310(46)
au	Ingress/egress duration (days)	0.01320(11)	0.01319(11)	0.01316(11)	0.01316(11)	0.01316(11)

Extended Data Table 3: Median Values and 68% Confidence Intervals for Selected Physical and Orbital Parameters of the KELT-9 System

The values within parentheses represent the uncertainty in the same number of least significant digits. For readability, we report asymmetric uncertainties with a difference of less than 5% in the absolute value of the positive and negative component as a symmetric uncertainty with a value equal to the maximum of the absolute values of the asymmetric components. Green text highlights the adopted fiducial model values. Yellow text highlights nominal parameter values that differ slightly (but are well within 1σ) from the fiducial model values. Red text highlights parameter nominal values and uncertainties that differ significantly (on the order of 1σ) from the fiducial model values. Red text in the Global Fit Configuration section highlights configuration settings that differ from the fiducial model settings.

Parameter	Description/Units	Model 5	Model 6	Model 7	Model 8	Model 9	Model 10
Global Fit Configuration	:						
M. & R. Constraints	Torres, Yonsei-Yale, None	Yonsei-Yale	Torres	Torres	Torres	None	None
Orbital Constraint	Circular, Eccentric	Eccentric	Circular	Circular	Eccentric	Circular	Circular
RV Slope	No Slope, Slope	No Slope	No Slope	No Slope	No Slope	No Slope	No Slope
Distance Prior	- Pp	Hipparcos	Hipparcos	None	Hipparcos	Hipparcos	Gaia
TTV Fit	Yes, No	No	No	No	No	No	No
Stellar Parameters:							
M_*	Mass (M_{\odot})	$2.53^{+(25)}_{-(20)}$	$2.55^{+(20)}_{-(19)}$	$2.57^{+(21)}_{-(19)}$	$2.54^{+(20)}_{-(19)}$	$1.74^{+(98)}_{-(71)}$	$2.28^{+(1.20)}_{-(0.89)}$
R*	Radius (R_{\odot})	$2.355^{+(76)}_{-(66)}$	2.368(61)	2.376(62)	2.357(64)	2.080(330)	2.280(350)
L_*	Luminosity (L_{\odot})	$54^{+(13)}_{-(11)}$	$55^{+(13)}_{-(11)}$	$56^{+(13)}_{-(11)}$	$55^{+(13)}_{-(11)}$	$43^{+(18)}_{-(14)}$	$51^{+(21)}_{-(16)}$
	Density (cgs)	$0.2727^{+(79)}_{-(49)}$	0.2705(28)	0.2705(28)	$0.2730^{+(86)}_{-(50)}$	0.2704(28)	0.2704(28)
$\log g_*$	Surface gravity (cgs)	4.098(16)	4.095(12)	4.096(12)	4.099(14)	$4.040^{+(65)}_{-(76)}$	$4.079^{+(62)}_{-(72)}$
T_{eff}	Effective temperature (K)	10220(440)	10250(440)	10280(450)	10250(440)	10290(450)	10280(450)
[Fe/H]	Metallicity	-0.03(20)	-0.02(20)	-0.00(20)	-0.02(20)	-0.01(20)	-0.01(20)
Planetary Parameters:							+(1.90)
M_P		2.87(85)	2.90(83)	2.91(84)	2.88(87)	2.18(1.10)	$2.61^{+(1.30)}_{-(0.98)}$
R_P	Radius (R _J)	$1.881^{+(62)}_{-(54)}$	$1.890^{+(51)}_{-(50)}$	$1.896^{+(52)}_{-(51)}$	1.881(53)	1.660(270)	1.820(280)
ρ _P	Density (cgs)	0.53(16)	0.53(15)	0.53(15)	0.54(16)	$0.61^{+(21)}_{-(18)}$	$0.55^{+(19)}_{-(17)}$ $3.30^{+(11)}_{-(14)}$
$\log g_{\mathbb{P}}$,	$3.30^{+(11)}_{-(15)}$	$3.30^{+(11)}_{-(14)}$	$3.30^{+(11)}_{-(14)}$	$3.30^{+(11)}_{-(15)}$	$3.30^{+(11)}_{-(14)}$	$3.30^{+(11)}_{-(14)}$
T_{eq}		4060(180)	4080(170)	4090(180)	4070(180)	4100(180)	4090(180)
$\langle F \rangle$	Incident flux $(10^9 \text{ erg s}^{-1} \text{ cm}^{-2})$	$61.6^{+(11.0)}_{-(9.8)}$	$62.0^{+(11.0)}_{-(10.0)}$	$63.0^{+(12.0)}_{-(10.0)}$	$62.0^{+(11.0)}_{-(10.0)}$	$63.0^{+(12.0)}_{-(10.0)}$	$63.0^{+(12.0)}_{-(10.0)}$
Orbital Parameters:		()	()	()	()	(,	()
T_C (BJD _{TDB})	Primary Center Time -2457095	0.68574(13)	0.68573(13)	0.68573(13)	0.68574(14)	0.68572(14)	0.68573(14)
T_S (BJD _{TDB})	Secondary Center Time -2457094	0.94358(170)	0.94517(13)	0.94517(13)	0.94358(170)	0.94516(14)	0.94516(14)
P	Period (days)	1.4811235(11)	1.4811235(11)		1.4811235(11)	1.4811235(11)	1.4811235(11)
a	Semi-major axis (AU)	$0.03466^{+(110)}_{-(93)}$	$0.03472^{+(90)}_{-(88)}$	$0.03483^{+(92)}_{-(90)}$	$0.03472^{+(90)}_{-(89)}$	0.03060(490)	0.03350(510)
λ	Spin-orbit alignment (degrees)	-84.8(1.4)	-84.8(1.4)	-84.8(1.4)	-84.8(1.4)	-84.8(1.4)	-84.8(1.4)
e	Eccentricity	$0.0045^{+(94)}_{-(31)}$	0.0 (fixed)	0.0 (fixed)	$0.0046^{+(100)}_{-(33)}$	0.0 (fixed)	0.0 (fixed)
ω	Arg. of periastron (degrees)	$-117^{+(23)}_{-(110)}$	90 (fixed)	90 (fixed)	$-115^{+(21)}_{-(110)}$	90 (fixed)	90 (fixed)
RV Parameters:		(220)			(110)		
K	RV semi-amplitude (m/s)	274(81)	277(78)	276(78)	274(82)	277(78)	277(78)
7TRES	RV offset (m/s)	462(58)	462(56)	462(56)	463(58)	462(56)	462(57)
ή	RV slope (m/s/day)	0.0 (fixed)					
Transit Parameters:							
R_P/R_*	Radius of the planet (R_*)	0.08209(43)	0.08203(42)	0.08204(43)	0.08202(43)	0.08203(42)	0.08204(42)
a/R_*	Semi-major axis (R_*)	$3.162^{+(30)}_{-(19)}$	3.154(11)	3.154(11)	$3.163^{+(33)}_{-(19)}$	3.154(11)	3.153(11)
í	Inclination (degrees)	86.81(25)	86.79(25)	86.79(25)	86.81(25)	86.79(25)	86.79(25)
b	Impact parameter	0.177(14)	0.176(14)	0.177(13)	0.177(14)	0.177(13)	0.177(13)
δ_p	Primary transit depth (10 ⁻³)	6.739(71)	6.728(69)	6.731(70)	6.728(70)	6.729(69)	6.730(70)
δ_s	Secondary transit depth (10^{-3})	1.003(98)	1.004(96)	1.004(96)	1.004(97)	1.004(97)	1.005(96)
T_{14}		0.16309(47)	0.16306(46)	0.16305(46)	0.16305(47)	0.16305(46)	0.16307(46)
τ	Ingress/egress duration (days)	0.01315(11)	0.01314(11)	0.01314(11)	0.01314(11)	0.01314(11)	0.01314(11)

Extended Data Table 3: Continued.