# The TESS-Keck Survey. XX. 15 New TESS Planets and a Uniform RV Analysis of All Survey Targets

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#### Abstract

The Transiting Exoplanet Survey Satellite (TESS) has discovered hundreds of new worlds, with TESS planet candidates now outnumbering the total number of confirmed planets from Kepler. Owing to differences in survey design, TESS continues to provide planets that are better suited for subsequent follow-up studies, including mass measurement through radial velocity (RV) observations, compared to Kepler targets. In this work, we present the TESS-Keck Survey's (TKS) Mass Catalog: a uniform analysis of all TKS RV survey data that has resulted in mass constraints for 126 planets and candidate signals. This includes 58 mass measurements that have reached  $\geqslant 5\sigma$ precision. We confirm or validate 32 new planets from the TESS mission either by significant mass measurement (15) or statistical validation (17), and we find no evidence of likely false positives among our entire sample. This work also serves as a data release for all previously unpublished TKS survey data, including 9,204 RV measurements and associated activity indicators over our three-year survey. We took the opportunity to assess the performance of our survey and found that we achieved many of our goals, including measuring the mass of 38 small ( $\langle 4 R_{\oplus}$ ) planets, nearly achieving the TESS mission's basic science requirement. In addition, we evaluated the performance of the Automated Planet Finder as survey support and observed meaningful constraints on system parameters, due to its more uniform phase coverage. Finally, we compared our measured masses to those predicted by commonly used mass–radius relations and investigated evidence of systematic bias.

Unified Astronomy Thesaurus concepts: [Exoplanet astronomy](http://astrothesaurus.org/uat/486) (486); [Radial velocity](http://astrothesaurus.org/uat/1332) (1332); [Hot Jupiters](http://astrothesaurus.org/uat/753) (753); [Super Earths](http://astrothesaurus.org/uat/1655) (1655); [High resolution spectroscopy](http://astrothesaurus.org/uat/2096) (2096); [Catalogs](http://astrothesaurus.org/uat/205) (205)

# 1. Introduction

The Kepler (Borucki et al. [2010](#page-47-0)) mission revealed a rich exoplanet landscape, showing that planets with short orbital periods ( $P < 100$  days) between the size of Earth and Neptune  $(1-4 R_{\oplus})$  are ubiquitous in our Galaxy (Dressing & Charbonneau [2013](#page-47-0); Fulton & Petigura [2018](#page-47-0)). However, the stars Kepler targeted are mostly too distant, and hence too faint, for precise RV follow-up. In contrast, the Transiting Exoplanet Survey Satellite (TESS, Ricker et al. [2015](#page-48-0)) mission was designed to preferentially discover planets around brighter, closer host stars that are highly amenable to follow-up efforts, including known planetary systems (Kane et al. [2009,](#page-48-0) [2021](#page-48-0)). The current era of precise RV mass measurements is beginning to reveal structure within the mass–radius distribution of small planets. For cool M dwarfs, Luque  $\&$  Pallé ([2022](#page-48-0)) found that the dichotomy of sub-Neptunes and super-Earths may indeed be a trichotomy composed of rocky, icy, and gaseous planets. While the connection between these groups and specific formation scenarios is debated (Rogers et al. [2023](#page-48-0)), resolving nuances in this region of mass–radius space is a significant motivator to expand the sample of planets around FGK stars with precisely determined masses.

Understanding the processes that result in the observed planet population hinges not only on well-characterized planets but also on the mitigation of bias and non-astrophysical scatter in data sets. This is especially important for the derivation of



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mass–radius relations. For example, publication bias (Burt et al. [2018](#page-47-0)) can not only bias a mass–radius relation to higher masses but may also artificially decrease intrinsic scatter. Additionally, uniform analyses have the power to eliminate scatter that may be inherent in heterogeneous data sets such as those compiled from various literature sources. Previous derivations of mass– radius relations have found evidence for more scatter in planet masses than can be described by measurement error alone (Weiss & Marcy [2014;](#page-48-0) Wolfgang et al. [2016](#page-48-0)). Some of this intrinsic scatter arises from the diversity of exoplanet compositions, but some may also be due to heterogeneous data sets. Mass–radius relations derived from homogeneous samples have the potential to more clearly distinguish between scatter caused by real compositional diversity as opposed to differences in analysis and target selection (Teske et al. [2021](#page-48-0)).

The plethora of planets now accessible with ground-based spectrographs has given rise to large collaborations that are able to pool many nights of telescope time together to follow up these targets and amass RV time series data. The TESS Follow-up Observing Program effort $53$  is a global collaboration encompassing many different efforts to confirm planet candidates including transit recovery, high-contrast imaging, and RV follow-up. Within RV follow-up, large collaborations include the GAPS program (Covino et al. [2013](#page-47-0)), the HARPS-N Rocky Planet Search (Motalebi et al. [2015](#page-48-0)), the CARMENES search for exoplanets around M dwarfs (Reiners et al. [2018](#page-48-0)), the Magellan-TESS survey (Teske et al. [2021](#page-48-0)), and many others. Another major RV follow-up effort, which is the focus of this work, is the TESS-Keck Survey (TKS).

TKS is an RV follow-up program that targeted TESSdiscovered systems with Keck/HIRES and APF/Levy. The main goals of TKS, as well as our target selection function, are detailed in full by Chontos et al. ([2022](#page-47-0), hereafter TKS0), with the primary goal being to measure 100 planet masses over a 3 yr time span. Our target selection process<sup>54</sup> identified TESS Objects of Interest (TOIs) that were valuable across both

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<sup>53</sup> https://[tess.mit.edu](https://tess.mit.edu/followup/)/followup/ <sup>54</sup> https://github.com/[ashleychontos](https://github.com/ashleychontos/sort-a-survey)/sort-a-survey

<span id="page-2-0"></span>complementary and competing science cases. This resulted in a final target list of 108 transiting planet candidates in 86 systems.<sup>55</sup> While many TKS publications have focused on the detailed characterization of individual systems, in this work we combine the homogeneous stellar properties of MacDougall et al. ([2023](#page-48-0)) with a uniform analysis of all RV survey data to compile a catalog of planet masses, eccentricities, and densities. This includes a reanalysis of RV data for previously confirmed planets in addition to 32 unpublished TESS systems.

The paper is organized as follows. Section 2 describes our data collection for high-resolution imaging and RV. Section [3](#page-4-0) details the algorithm we developed for homogeneously derived masses for all planets in our sample, and we expand upon caveats to this algorithm in Section [4.](#page-6-0) In Section [5](#page-6-0), we compare our measured masses with those previously published in TKS and in the wider literature. In Section  $6$ , we discuss how well our survey met its original goals, and we assess the contribution of APF to our survey. In Section [7,](#page-14-0) we review the ability of commonly used mass–radius relations to predict a planet's mass.

#### 2. Observations

# 2.1. Radial Velocities

All of our spectroscopic data used in this analysis were taken entirely by the High Resolution Echelle Spectrometer (HIRES, Vogt et al. [1994](#page-48-0)) on the Keck I telescope at W. M. Keck Observatory, and the Levy Spectrometer located on the Automated Planet Finder (APF, Vogt et al. [2014a](#page-48-0)) telescope. HIRES is an echelle spectrometer with a resolving power of ∼67,000, depending on the slit used, and a wavelength range spanning from 374 to 970 nm. We extracted RVs from the spectroscopic data (reduced using the standard procedure described in Howard et al. [2010](#page-48-0)), using the iodine cell method outlined in Butler et al. ([1996](#page-47-0)). The APF is a robotic telescope operated at Lick Observatory (Vogt et al. [2014b](#page-48-0)). The Levy spectrograph has a resolving power of 100,000, and also utilizes the same iodine cell method to extract precise RVs as for HIRES. The APF has the same wavelength span as HIRES; 374–970 nm.

The total survey includes 7894 HIRES RVs taken over the course of 974 Keck observing nights when including archival RVs. Narrowing to only HIRES RVs taken during TKS observing programs, we collected 4943 HIRES RVs over 301 observing nights. We also use 5,989 APF RVs taken over the course of 1205 APF observing nights. Again, narrowing to only observations taken during TKS-affiliated observing semesters, we acquired 4261 APF RVs over 710 observing nights.

Observational baseline is a key component of RV surveys, as a long baseline allows for improved sensitivity to long-period, non-transiting companions. When utilizing archival HIRES RVs, our maximum survey baseline was 6502 days and our minimum was 59 days. The average RV baseline of the survey was 1427 days. A more detailed breakdown by instrument is given in Table [1](#page-3-0). In addition, many of the stars in our survey

have RVs from other instruments, but for this work we do not include non-TKS data.

In general, the TKS survey was broken down into a number of different science cases with different observation and cadence goals, and such science cases were the primary drivers in our observing strategy. Rather than detail all of the TKS science cases here, we refer the interested reader to Chontos et al. ([2022](#page-47-0)), which describes our strategies in detail.

# 2.2. High-resolution Imaging

High-resolution imaging (HRI) is an important step in validating planet candidates, due to its ability to resolve additional stars not resolved by Gaia. Visual binary companions, regardless of whether they are gravitationally bound or not, can dilute the TESS lightcurve, leading to underestimated planet radii (Ciardi et al. [2015](#page-47-0)). In this section, we describe the observations taken of the systems that are newly validated in this work and discuss the presence of a companion to TOI-1443. Since many of our TOIs were observed by multiple instruments, in Table [5](#page-19-0) we provide a summary of the imaging observations made.

#### 2.2.1. Keck-NIRC2

NIRC2 is an adaptive optics enabled imaging camera on the on the 10 m Keck II telescope. Narrow-angle mode was used, providing a pixel scale of around  $0$ ." $01 \text{ px}^{-1}$  and a full field of view of about 10". Observations were taken in either the Br- $\gamma$ ,  $K$ , or  $K'$  filters. Observations were taken using the natural guide star AO system in the standard three-point dither pattern to avoid the lower-left quadrant of the detector, which is typically noisier than the other three quadrants. The dither pattern has a step size of 3″. Each dither position was observed three times, with  $0\rlap.{''}5$  positional offsets between each observation, for a total of nine frames.

#### 2.2.2. Palomar-PHARO

PHARO is an adaptive optics enabled imaging camera on the 5.1 m Hale telescope with a pixel scale of  $0''025 \text{ px}^{-1}$  and a full field of view of around 25″. Observations were taken in the Br- $\gamma$  and acquired using the natural guide star AO system P3K (Dekany et al. [2013](#page-47-0)) in the standard five-point quincunx dither pattern with steps of 5″. Each dither position was observed three times, with  $0$ ." 5 positional offsets between each observation, for a total of 15 frames.

#### 2.2.3. Shane-ShARCS

ShARCS is an adaptive optics enabled camera on the Shane 3 m telescope at Lick Observatory (Kupke et al. [2012;](#page-48-0) Gavel et al. [2014;](#page-47-0) McGurk et al. [2014](#page-48-0)). Observations were taken with the Shane adaptive optics system in natural guide star mode in order to search for nearby, unresolved stellar companions. We used a variety of filters including  $K_s$  ( $\lambda_0 = 2.150 \mu m$ ,  $\Delta\lambda = 0.320 \mu m$ ) and J ( $\lambda_0 = 1.238 \mu m$ ,  $\Delta\lambda = 0.271 \mu m$ ). Information about our observing strategy can be found in Savel et al. ([2020](#page-48-0)). We reduced the data using the publicly available SIMMER pipeline (Savel et al. [2022](#page-48-0)).<sup>56</sup> Further information about these data can be found in C. D. Dressing et al. (2024, in

<sup>&</sup>lt;sup>55</sup> Targets in this work are sometimes referred to by their TKS ID, which is the preparation) and on ExoFOP. name used in the California Planet Search queue system. For reference, stars in the Henry Draper catalog are referred to by their numerical designation (e.g., HD 95072  $\rightarrow$  95072), while TOIs are given a shorthand reference (e.g., TOI–1776  $\rightarrow$  T001776)

<sup>56</sup> https://[github.com](https://github.com/arjunsavel/SImMER)/arjunsavel/SImMER

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<b>Observation Statistics</b>								
Statistic	HIRES $(w / \text{archival})$	APF $(w / \text{archival})$	HIRES (TKS only)	APF (TKS only)				
$N_{\rm RV}$	7894	5989	4943	4261				
Nights Observed	974	1205	301	710				
Longest Baseline (days)	6502	3391	1076	1249				
Shortest Baseline (days)	59		59					
Average Baseline (days)	1427	965	828	745				

Table 1

#### 2.2.4. Gemini-NIRI

NIRI (Hodapp et al. [2003](#page-47-0)) is an adaptive optics enabled imaging camera on the Gemini North 8 m telescope. In each case, we observed using the Br $\gamma$  filter centered on 2.166  $\mu$ m, and collected nine individual frames, dithering the telescope in a grid pattern between each frame. Integration times for individual frames were calculated based on the stellar magnitude, and were 4.4 s, 6.6 s, and 4 s, respectively, for TOI–1184, TOI–1443, and TOI–1451. To process the data, we first constructed a sky background by median combining the dithered science frames. We then removed bad pixels, flat fielded, and subtracted the sky background. Cleaned frames were aligned by fitting the stellar PSF in each frame, and then coadded. Sensitivity of each observation was calculated as a function of radius by injecting fake companions and scaling their background so they could be redetected at  $5\sigma$ . For visual companions identified in the Gemini-NIRI data, we calculated astrometry relative to the host star by fitting a Gaussian to each of the star and companion, and then applying the pixel scale of 0"021 pixel<sup>-1</sup>. We performed aperture photometry on both the star and the visual companion in our reduced images, using a small aperture around the position determined in the previous section.

For each target, we reached contrast sensitivities of at least 5 mag beyond ∼300 mas, and at least 7mag in the backgroundlimited regime beyond ∼1″; images have a field of view of  $26'' \times 26''$ , though sensitivity is reduced in the outermost  $\sim$ 3″ where fewer dithered images are being combined.

#### 2.2.5. Gemini-'Alopkeke and Zorro

Alopeke and Zorro are high-resolution optical speckle instruments mounted on the Gemini North and South 8 m telescopes, respectively. Alopeke and Zorro are identical instruments that provide simultaneous speckle imaging in two bands, usually operated using narrowband filters centered at 562 and 832 nm.Speckle imaging consists of obtaining many thousands of very short exposures (typically 60 msec/frame at Gemini), which are then combined and subjected to Fourier analysis in our standard reduction pipeline (see Howell et al. [2011;](#page-48-0) Horch et al. [2012](#page-47-0)). Output data products include a reconstructed image in each filter and determination of robust  $5\sigma$  contrast limits on companion detections in each bandpass. If a close  $(0''02-1''2)$  companion source is discovered, its position angle, separation, and flux ratio are reported. The instruments are described in detail in Scott et al. ([2021](#page-48-0)).

# 2.2.6. WIYN-NESSI

The NESSI speckle imager (Scott [2019](#page-48-0)) is mounted on the 3.5 m WIYN telescope at Kitt Peak. NESSI simultaneously acquires data in two bands centered at 562 and 832 nm using high-speed electron-multiplying CCDs (EMCCDs). We collected and reduced the data following the procedures described in Howell et al. ([2011](#page-48-0)). The resulting reconstructed images achieved a contrast of  $\Delta$ mag ~ 6 at a separation of 1″ in the 832 nm band.

#### 2.2.7. Speckle Polarimeter

The Speckle Polarimeter (Safonov et al. [2017](#page-48-0)) is mounted on the the 2.5 m telescope at the Caucasian Observatory of Sternberg Astronomical Institute (SAI) of Lomonosov Moscow State University. Electron Multiplying CCD Andor iXon 897 was used as a detector. For each target, we accumulated 4000 frames with 30 ms exposure times. The detector has a pixel scale of 20.6 mas  $pixel^{-1}$ , and the angular resolution was 89 mas. For all targets except TOI–1443, we did not detect any stellar companions.

#### 2.2.8. AstraLux

The AstraLux instrument (Hormuth et al. [2008](#page-47-0)) is mounted on the 2.2 m telescope of the Calar Alto Observatory (Almería, Spain). AstraLux uses the lucky-imaging technique, obtaining a large number of very short exposures (shorter than the coherence time of the atmospheric turbulence) that are subsequently selected and combined. The datacube is then processed by the automatic observatory pipeline (Hormuth et al. [2008](#page-47-0)), which, besides doing the basic reduction of the individual frames, also selects the 10% with the highest Strehl ratio (Strehl [1902](#page-48-0)) and combines them into a final high-spatialresolution image. All images were obtained in the SDSSz bandpass (peak of the instrument resolution) with an individual exposure time of 10 or 20 ms (depending on the source and weather conditions) and a field of view windowed to  $6'' \times 6''$ . This final image is used to compute the sensitivity limits (also known as the contrast curve). Such limits are computed by using our own developed astrasens package<sup>58</sup> with the procedure described in Lillo-Box et al. ([2012](#page-48-0), [2014](#page-48-0)).

# 2.2.9. A Companion to TOI–1443

The only target in our sample where a diluting star is seen is TOI-1443, which was imaged with Gemini-NIRI and the SAI speckle polarimeter (Figure [1](#page-4-0)). TOI–1443 was imaged with the speckle polarimeter in four epochs: 2020 December 26, 2022 June 3, 2023 September 23, and 2023 September 29, in the  $I_c$ band. Each time the object was resolved, the binarity parameters were determined by an approximation of the average power spectrum with the model of the binary star (Strakhov et al. [2023](#page-48-0)). The separation and position angle of secondary companion for these epochs was found to be

<sup>57</sup> https://[www.gemini.edu](https://www.gemini.edu/sciops/instruments/alopeke-zorro/)/sciops/instruments/alopeke-zorro/ <sup>58</sup> https://[github.com](https://github.com/jlillo/astrasens)/jlillo/astrasens

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Figure 1. Contrast curves for TOI-1443 from Gemini-NIRI (blue) and SAI Speckle Polarimeter (black). The  $1\sigma$  error for the NIRI contrast curve is given as the blue envelope. A companion is observed for TOI-1443 at ∼0.<sup>"</sup> 5 that is not detected in Gaia DR3.

 $518 \pm 6$  mas and  $68^\circ$   $2 \pm 0^\circ$   $4$ ,  $531 \pm 6$  mas and  $69^\circ$   $3 \pm 0^\circ$   $3$ , 537  $\pm$  5 mas and 68°.6  $\pm$  0°.4, and 507  $\pm$  4 mas and  $68^\circ$   $2 \pm 0^\circ$   $4$ . The total proper motion reported by Gaia DR3 for TOI–1443 is  $90.511$  mas yr<sup>-1</sup>. This would yield a total relative shift of 249.54 mas between the first and last epochs, if the secondary companion was a background star. We conclude that this binary is a gravitationally bound system. We additionally note that, while orbital motions are statistically detected through these observations, they are not sufficient to provide meaningful orbital constraints.

From Monet et al. ([2003](#page-48-0)), TOI-1443 has  $I_c = 9.75$ , giving the primary an  $V-I_c$  color of 1.08 mag. Interpolating Table 5 of Pecaut & Mamajek ([2013](#page-48-0)), this gives an effective temperature around 5300 K (spectral type of K0V), consistent with its spectroscopically derived parameters. The Speckle Polarimeter gives an average magnitude difference between the components of  $\Delta I_c = 5.0$  mag, while NIRI gives  $\Delta K = 3.64 \pm 1.0$ 0.01 mag. From these values, we obtain a  $K-I_c$  color of 2.47 mag for the secondary component. Again interpolating Table 5 of Pecaut & Mamajek ([2013](#page-48-0)) provides an approximate effective temperature of 3100 K, consistent with an M4.5V spectral type. A  $\Delta I_c$  of 5.0 mag results in ~0.5% contamination, assuming the planet transits the primary.

We also observed a significant RV trend of 0.029  $\pm$ 0.0044 m s−<sup>1</sup> d−<sup>1</sup> . Using ethraid (Van Zandt & Peti-gura [2024](#page-48-0)),<sup>59</sup> we use the RV trend and the NIRI sensitivity curve to constrain the combinations of mass and semimajor axis of an object that could be responsible for the trend. We find that an M4.5V dwarf ( $M_* \sim 0.18$   $M_{\odot}$ ) at the observed angular separation is a plausible explanation for the observed trend in our RVs, lending additional support to a bound scenario.

## 3. Methods

In this section, we describe the analysis of our survey data. In an effort to achieve homogeneity, we outline the standardized decision tree that informed our RV fits. We also

describe the uniform set of priors used. In the few instances where RVs from Keck/HIRES and/or APF exist prior to the inception of our survey, we do include those in this analysis. This decision is motivated by a desire to release these data, and by the fact that the existence of RVs on a given target was a favorable metric in our target selection efforts (Chontos et al. [2022](#page-47-0)).

#### 3.1. Radial Velocity Fitting

A schematic of our methodology is shown in Figure [2](#page-5-0). We searched for signals in each RV data set using RVSearch (Rosenthal et al. [2021](#page-48-0)), which uses an iterative approach to fit candidate signals, as determined by the false-alarm probability,  $FAP > 1\%$ , and then compare against an N versus N-1 planet model. We seeded the search with parameters of known, transiting planets in each system taken from the TKS Systems Properties Catalog (MacDougall et al. [2023](#page-48-0)).

Once all candidate signals within the RV data set are identified by RVSearch, we begin our RV fit. Using RadVel (Fulton & Petigura [2018](#page-47-0)), we adopt a top-down approach whereby we assume a maximally complex RV model where all parameters (see below) are initially included. Free parameters included Keplerian signal parameters  $(\rho_i, T_{0,i}, K_i, \sqrt{\overline{e_i}} \sin(\omega_i))$ ,  $\sqrt{e_i}$  cos( $\omega_i$ )) for all known planets (transiting and nontransiting), and candidate signals, linear  $(\dot{\gamma})$  and quadratic  $(\ddot{\gamma})$ trend terms, and offset/jitter terms  $(\gamma, \sigma)$  specific to each instrument.

We ran RadVel in the fitting basis described above and explored the posterior parameter space using Markov Chain Monte Carlo (MCMC) sampling (emcee, Foreman-Mackey et al. [2013](#page-47-0)). This basis reparameterizes e (eccentricity) and  $\omega$ (argument of periastron) into  $\sqrt{e}$  cos( $\omega$ ) and  $\sqrt{e}$  sin( $\omega$ ) for better MCMC performance. The MCMC was run until the chains were well mixed, i.e., they exhibited a G-R statistic less than  $1.001$  (Gelman  $&$  Rubin  $1992$ ; Ford  $2006$ ) and the number of independent samples was greater than 1000 for all free parameters for at least five consecutive checks. We then determined which combination of model parameters was justified using the Akaike Information Criterion (AIC), with the preferred model being the one with the smallest AIC. In cases where the lowest model AIC values differed by  $\langle 2$ , the model with the fewest free parameters was selected. One exception is when model selection rejects a signal of a known, transiting planet. In these cases, we instead accept the model with the lowest AIC score that still includes all transiting planets.

We aimed to uniformly derive mass estimates for as many planets as possible. For systems with high planet multiplicity and/or especially active host stars, our methodology may not return the best planet parameters. Nonetheless, the primary goal of this analysis is the homogeneous derivation of masses, not necessarily the best derivation of masses. Similarly, in cases where additional data from other instruments exists, a more precise mass could be determined. We discuss caveats to our uniform fits in Section [4](#page-6-0) and compare our masses to literature values in Section [5.1.](#page-7-0)

# 3.2. Selection of Priors

We implemented a homogeneous set of priors for every planet in the sample detailed in Table [2](#page-7-0). Parameters such as trend, curvature, and jitter were given broad, uniform priors.

<sup>59</sup> https://[github.com](https://github.com/jvanzand/ethraid)/jvanzand/ethraid

<span id="page-5-0"></span>

Figure 2. A flowchart of our algorithm for homogeneous determination of planet masses.

For transiting planets, the priors applied to their period and time of conjunction were Gaussian, with centers and widths equal to the posteriors of the transit fits performed in MacDougall et al. ([2023](#page-48-0)). Non-transiting planets were given broad Jeffreys priors (Jeffreys [1946](#page-48-0)) on their orbital period, with the initial guess sourced from the RVsearch results. These planets were also given a broad, uniform prior on time of conjunction, with lower and upper bounds extending half of the orbital period. All other free parameters, including K-amplitude, K, were unrestricted and allow for negative values.

## 3.3. Gaussian Processes

A Gaussian process (GP) model can often be very effective at accounting for stellar activity from RV data sets (e.g., Haywood et al. [2014](#page-47-0); López-Morales et al. [2016](#page-48-0)). GPs are flexible, however, and can overfit on sparse data sets (e.g., Blunt et al. [2023](#page-47-0)). A GP is used on a TKS systems if (1)  $\log R'_{HK}$  -4.9, and (2) if  $N_{obs} > 30$ . The first condition is to ensure that stellar activity modeling is indeed warranted, as  $\log R'_{HK}$  values are found to positively correlate with increased stellar activity (Wright [2005](#page-48-0); Isaacson & Fischer [2010](#page-48-0); Hillenbrand et al. [2015](#page-47-0)). The second condition is intended to prevent overfitting while using a GP.

We used the QuasiPeriodic GP kernel to model activity (Haywood et al. [2014;](#page-47-0) López-Morales et al. [2016](#page-48-0)). The covariance between observations taken at times  $t_i$  and  $t_j$  is given as

$$
C_{i,j} = \eta_1^{2*} exp \left( \frac{-|t_i - t_j|^2}{\eta_2^2} - \frac{\sin^2 \left( \frac{\pi |t_i - t_j|}{\eta_3} \right)}{2\eta_4^2} \right).
$$
 (1)

An advantage of the QuasiPeriodic GP kernel is the interpretability of its hyperparameters, as most correspond, at

least approximately, with a physical parameter.  $\eta_1$  is the GP amplitude in m  $s^{-1}$ , and it corresponds to the amplitude of the stellar variability.  $\eta_2$  is the exponential decay length, and it approximately represents the lifetime of a stellar spot.  $\eta_3$  is the recurrence timescale, which is usually well represented by the rotation period of a star.  $\eta_4$  is the periodic scale length, which varies from 0 to 1 and approximately controls how periodic the structure of the GP signal is.

Similarly to the planet parameters, we strove for homogeneous priors on GP hyperparameters (Table [2](#page-7-0)). The GP amplitude  $\eta_1$  is not known a priori, and we adopt a broad Jeffreys prior. We also use a broad Jeffreys prior for  $\eta_2$ , though we prevent the value from falling below the system rotation period,  $\eta_3$ , as recommended by Kosiarek & Crossfield ([2020](#page-48-0)).  $\eta_3$  is informed by either a periodogram of the photometry or Svalues. Finally, we restrict  $\eta_4$  using a Gaussian prior recommended by López-Morales et al. ([2016](#page-48-0)). This is to prevent overfitting with the GP, which can happen for small values of  $\eta_4$ .

Even when a GP is included in our maximally complex initial fit, if the GP is not favored in model comparison, it is not included in our final fit. The fitted values for our GP hyperparameters are given in Table [7](#page-32-0).

# 3.4. Derivation of Physical Parameters

The stellar and transit properties for this work are taken from MacDougall et al. ([2023](#page-48-0)). In summary, stellar spectroscopic properties ( $T_{\text{eff}}$ , [Fe/H],  $\log(g)$ ) were derived from HIRES spectra using SpecMatch (Petigura [2015](#page-48-0); Yee et al. [2017](#page-48-0)), while  $M_*$  and  $R_*$  were obtained using isoclassify (Huber et al. [2017](#page-48-0); Berger et al. [2020](#page-47-0), [2023](#page-47-0)) with the spectroscopic parameters used to inform input priors. Similarly, the transiting planet parameters are from the same study's analysis of TESS lightcurves. In MacDougall et al. ([2023](#page-48-0)), stellar parameter uncertainties do not account for systematic errors in model

<span id="page-6-0"></span>isochrones (Tayar et al. [2022](#page-48-0)) but do provide a recommendation to inflate stellar parameter uncertainties. In the following derivation of physical properties, we make the recommended corrections.

We derive the parameters listed in Table [4](#page-16-0) by multiplying the MCMC chains from each fit. For the planet mass, we use the equation:

$$
\frac{m_p \sin(i)}{M_J} = \frac{K}{K_0} \sqrt{1 - e^2} \left(\frac{M_*}{M_\odot}\right)^{2/3} \left(\frac{P}{1 \text{ yr}}\right)^{1/3},\tag{2}
$$

where  $K_0 = 28.4329 \text{m s}^{-1}$ . The semimajor axis, *a*, is calculated using Kepler's third law in the form:

$$
a = \left(\frac{P^2 \ G \ M_{\text{Total}}}{4\pi^2}\right)^{1/3}.
$$
 (3)

The majority of the planets in our sample are transiting, allowing us to derive the true mass with knowledge of the planet's inclination. We calculate the inclination of the transiting planets using: $60$ 

$$
b = \frac{a}{R_*} \cos(i),\tag{4}
$$

where *a* is taken from Equation (3) and both  $R_{*}$  and *b* are from MacDougall et al. ([2023](#page-48-0)). For  $R_{*}$ , we assume Gaussiandistributed uncertainties; however, the posterior distribution for the impact parameter in transit analyses is typically non-Gaussian. To reflect this, we use the MCMC chains from MacDougall et al.  $(2023)$  $(2023)$  $(2023)$  <sup>61</sup> and use sklearn's Kernel Density Estimation function to construct the impact parameter's distribution. We sample the resulting distribution to obtain a chain of samples for  $b$  of equal length to those used in our RadVel analyses. The chains for  $b$ ,  $R_*$ , and  $a$  are used to derive i. For each derived parameter in Table [4](#page-16-0), we list the median value with the 15.6% and 84.1% percentiles as uncertainties.

# 4. Caveats to a Homogeneous RV Fit

We developed our fitting methodology to be applicable to most systems. However, there were some instances where we departed from the steps stated in Section [3](#page-4-0). In the interest of reproducibility, here we clarify decisions made for three systems.

# 4.1. TOI–1444

TOI–1444 hosts a confirmed ultra-short-period planet (b) at 0.47 days along with a candidate non-transiting companion (c) at ∼16 days (Dai et al. [2021a](#page-47-0)). Since the publication of this system, we collected 20 RVs in addition to the existing 56; however, these additional observations did not significantly increase in the detection significance for this signal. Curiously, when both the "b" and "c" components are removed, RVSearch reveals a strong peak at 89 days. When only removing the "b" component, neither the 16 day or 89 day signals are confidently detected. According to our

methodology, the 89 day signal would be included as a candidate signal, but since its detection is dependent on the removal of a signal that has yet to be definitively confirmed, we decide to exclude it from our fits. This period of 89 days is also close to the 90 day alias seen in HIRES data (Rosenthal et al. [2021](#page-48-0)).

# 4.2. TOI–1726

TOI–1726 is a young Sun-like star that hosts two super-Earths of comparable size (Damasso et al. [2023](#page-47-0)) with a rotation period of ∼6.4 days that is clearly seen in periodograms of both our RVs and our activity indicators. This rotation period happens to be close to the orbital period of TOI–1726 b at 7.1 days, which, combined with the low expected semiamplitude for both planets, motivated us to stop observing this target well before it hit its 100 observations quota. According to our methodology, we placed an informative prior on the periodic component of the GP informed by the rotation period of the star, resulting in a mass constraint for the c planet of  $17.3 \pm 9.9M_{\oplus}$  that is in agreement with the previous publication of this system. However, we obtain a mass for the "b" planet of  $37 \pm 9.6M_{\odot}$ , which is significantly higher than both the estimate from Damasso et al. ([2023](#page-47-0)) and the masses of all other sub-Neptunes of similar size.<sup>62</sup> Given the similarity between the rotation and orbital period, we suspect the GP might be overfitting and possibly inflating the RV semi-amplitude for the "b" planet to some degree. We keep this result in the catalog for completeness but caution accepting it at face value.

#### 4.3. TOI–1759

TOI–1759 is an M0V dwarf with moderate levels of activity  $(\log R'_{HK} = -4.69)$  that hosts a  $\sim 3 R_{\oplus}$  planet (Martioli et al. [2022a](#page-48-0); Espinoza et al. [2022](#page-47-0)). A stellar rotation period of 35.65 days was identified through SPIRou polarized spectra. Though this signal was not clearly seen in either the CARMENES or SPIRou RVs, it is prominent in the HIRES RVs. This could be due to the different band passes of the instruments and the expected spot temperatures. Additionally, this signal is absent in both our  $S_{HK}$  time series and the TESS photometry. Given our methodology for building RV fits, this stellar signal would have been treated as Keplerian with a blind GP applied. In this case, we chose to place a Gaussian prior on the rotation term in the GP. While we are only able obtain an upper limit for the mass, a reanalysis of this system using additional new data sets is underway (A. S. Polanski et. al 2024, in preparation).

# 5. Results

We report mass estimates or upper limits for a total of 126 planets and candidate signals across 86 planetary systems. Of these planets, 112 transit their host star. In Figure [3](#page-8-0), we show our full sample in comparison to all known exoplanets with a significant mass measurement,  $63$  while in Figure [4](#page-9-0) we show a detailed view of the mass–radius space for small planets. In Table [6](#page-22-0), we report values for all fitted parameters, while Table [4](#page-16-0) lists the derived mass, density and semimajor axis for each candidate along with the previously published mass estimate (where available).

<sup>60</sup> The precise definition of impact parameter includes a factor of  $(1 - e^2)/(1 + e \sin \omega)$ , but given that transiting planets have near-zero cos*i*, this factor usually has a negligible contribution to the derived mass.

 $61$  The chains for each transit fit are available at https://[github.com](https://github.com/mason-macdougall/tks_system_properties)/masonmacdougall/[tks\\_system\\_properties](https://github.com/mason-macdougall/tks_system_properties).

 $\frac{62}{2}$  Hadden & Lithwick ([2014](#page-47-0))'s transit timing analysis of Kepler planets found sub-Neptunes comparable in mass to our estimate for TOI–1726c. However, they are still consistent with the highest RV-determined sub-Neptune mass within their uncertainties.<br> $63 \text{ F}$  and  $84 \text{ F}$ 

 $3$  From the NASA Exoplanet Archive as of Sept. 2023

<span id="page-7-0"></span>

#### Table 2 Survey Priors

#### Notes.

<sup>a</sup>  $\mathcal N$  is a normal prior with  $\mathcal N$  (mean, standard deviation).<br>
<sup>b</sup>  $\mathcal U$  is a uniform prior with  $\mathcal U$  (lower,upper).<br>
<sup>c</sup>  $\mathcal J$  is a Jeffreys prior with  $\mathcal J$  (lower,upper).<br>
<sup>c</sup>  $\mathcal J$  is a Jeffreys prior wit handed coordinate system.

 $\mu_{M23}$  and sd<sub>M23</sub> refers to the mean and standard deviation of the posterior taken from MacDougall et al. ([2023](#page-48-0)).

In this section, we compare our results to previously published measurements, present newly confirmed planetary systems, and statistically validate planets that did not reach a mass detection significance threshold of  $3\sigma$  (see Section [5.3](#page-10-0)).

## 5.1. Comparison with Confirmed Planets

Figure [5](#page-9-0) compares our derived masses with literature masses for 62 planets in 33 previously confirmed TESS systems. For the majority of planets, we find that our approach produces consistent results, with most measurements within  $1\sigma$ .

Of the 11 planets where we were only able to determine an upper limit, 5 of them were in the TOI–1136 system where both multiplicity and stellar activity conspired to require a much more comprehensive analysis (Dai et al. [2023](#page-47-0), Beard et al. in review). Although expected, this highlights the difficulty in homogeneous mass measurements for multi-planet systems—a difficulty that is exacerbated by high levels of activity. Many of our mass measurements with the largest measurement error  $(M/\sigma_M < 3\sigma)$  in Figure [5](#page-9-0) come from such systems (e.g., TOI–561, HD 191939, and TOI–1136). In two cases, HD 25463c and HD 12572c, only an upper limit was obtained in both this work and in their dedicated TKS publication, despite a more in-depth joint RV-transit fit (Akana Murphy et al. [2023](#page-47-0)), suggesting these planets require additional data. Conversely, for 24 previously published planets, we are

able to obtain a mass measurement more precise than previous literature values, with an average improvement in the mean mass uncertainty of 20%.

#### 5.2. New RV Confirmation of TOIs

We achieve a mass significance of  $\geq 3\sigma$  for 15 previously unconfirmed candidate planets. Given the vetting performed on these targets in Chontos et al. ([2022](#page-47-0)), we confirm these as bona fide planets. While periodogram analyses of the RV data for each of these systems all reveal some amount of power at the candidate's period, eight had a false-alarm probability less than 1%. For the remaining seven systems, we take the additional step of statistically validating them (Section [5.3](#page-10-0)). Individual descriptions of these newly confirmed planetary systems are given in this section, and our periodogram analyses and RV fits can be found in the [Appendix.](#page-21-0)

# 5.2.1. TOI–260 b

TOI–260 b (HIP 1532) is a super-Earth  $(1.47^{+0.12}_{-0.08} R_{\oplus})$ orbiting a K-type star (4050  $\pm$  109 K) with an orbital period of ∼13.5 days. Our S-index time series contains a stellar activity signal at ∼37 days that is also seen in our RVs; however, we note that this is significantly different from the rotation period of ∼15 days obtained through TESS photometry by Howard

<span id="page-8-0"></span>

Figure 3. Mass–radius diagram for the full TKS sample (blue) in comparison to the all known exoplanets with  $a \ge 3\sigma$  mass. Markers are scaled according to the mass precision (size  $\alpha M/\sigma_M$ ), while triangles represent 3 $\sigma$  upper limits. Points outlined in gold represent newly confirmed and validated planets in our sample. The region outlined in black represents the region of M-R space shown in more detail in Figure [4.](#page-9-0)

et al. ([2021](#page-48-0)). An RV fit (Figure [11](#page-33-0)) with a GP included allows us to modestly constrain the mass of the transiting planet to  $3.3 \pm 1.1$   $M_{\oplus}$ , placing it between the Earth-like and 50/50 water world composition curves, although it is consistent with an Earth-like composition within the mass error. With an RV semi-amplitude of 1.3 m s<sup>-1</sup>, a mass measurement for this planet lies at the extreme end of HIRES' capabilities, and the system would benefit from additional observations with higherprecision spectrographs.

# 5.2.2. TOI–1173 b

TOI–1173 b is a Saturn-sized (9.02 $^{+0.16}_{-0.15}R_{\oplus}$ ) planet orbiting a G-type star (5414 ± 125 K) with an orbital period of  $\sim$ 7.1 days. We were able to obtain a  $7\sigma$  mass constraint of 28.3  $\pm$  4.0  $M_{\oplus}$  (Figure [12](#page-34-0)) with 28 RVs, giving TOI–1173 b a bulk density of  $0.21 \pm 0.04$  g cm<sup>-3</sup>, which is a third of that of Saturn. TOI-1173 has an estimated TSM of 182, making it a valuable target for transmission spectroscopy.

# 5.2.3. TOI–1184 b

TOI–1184 b is a sub-Neptune  $(2.41^{+0.10}_{-0.14} R_{\oplus})$  orbiting a K-type star (4616 ± 104 K) with an orbital period of  $\sim$ 5.7 days. Our RV fits (Figure [13](#page-35-0)) were able to constrain the planet mass to  $6.8 \pm 2.3$   $M_{\oplus}$ , making it one of the least dense sub-Neptunes in the TKS sample.

## 5.2.4. TOI–1194 b

TOI–1194 b is a Saturn-sized  $(8.72^{+0.16}_{-0.15} R_{\oplus})$  planet orbiting a G-type star  $(5393 \pm 124 \text{ K})$  with an orbital period of  $\sim$ 2.3 days. We were able to obtain a 30 $\sigma$  mass constraint of  $120 \pm 6$   $M_{\oplus}$  (Figure [14](#page-36-0)) with 32 RVs, corresponding to a density slightly higher than that of Saturn. The shorter period of TOI–1194 b places it at the upper edge of the Neptune desert, and its high density could be the result of past or ongoing atmospheric erosion.

#### 5.2.5. TOI–1248 b

TOI–1248 b is a sub-Saturn-sized  $(6.81^{+0.12}_{-0.12} R_{\oplus})$  planet orbiting a K-type star (5205  $\pm$  120 K) with an orbital period of ∼4.4 days. Similarly to TOI–1173 b, our RVs (Figure [15](#page-37-0)) constrain the mass to 27.4  $\pm$  3.6  $M_{\oplus}$ , giving it a density of  $0.47 \pm 0.1$  g cm<sup>-3</sup>.

#### 5.2.6. TOI–1279 b

TOI–1279 b is a sub-Neptune  $(2.66^{+0.14}_{-0.11} R_{\oplus})$  orbiting a G-type star (5457 ± 127 K) with an orbital period of  $\sim$ 9.6 days. We obtain a  $4\sigma$  mass constraint of  $10.6 \pm 2.5$   $M_\oplus$  (Fig. [16](#page-38-0)), placing this planet in a region of the mass–radius diagram where many of the sub-Neptunes in our sample are concentrated. We also measure a significant downward RV trend of -0.0079  $\pm$  0.0020 m s<sup>-1</sup> d<sup>-1</sup> that does not correlate

<span id="page-9-0"></span>

Figure 4. The mass–radius space for planets below  $4 R_{\oplus}$  in the TKS sample. Points are color-coded based on the planets' equilibrium temperature, while the size is scaled according their mass significance: larger points have a more precise mass measurement. Points outlined in gold are planets whose mass constraints are newly published in this work. Blue contours represent a Gaussian Kernel Density Estimate (KDE) of all known exoplanets around FGK stars. We show the composition tracks for pure iron, an Earth-like composition, and a 50/50 water-silicate mix (Zeng et al. [2019](#page-48-0)). In addition, we include the photoevaporation mass–radius distributions from Rogers et al. ([2023](#page-48-0)).



Figure 5. A comparison of mass measurements for the 62 planets in our sample for which masses had been published previously. Black stars indicate the literature value comes from a TKS publication, while blue circles indicate a mass publication outside of our survey. The insets show the comparison for objects exceeding 500  $M_{\oplus}$ (upper left) and a zoomed-in view of the dense region between 5 and 14  $M_{\oplus}$ (lower right).

with activity indicators, suggesting the possible presence of a long-period companion. Additionally, TOI–1279 b is the only single-TOI system where MacDougall et al. ([2023](#page-48-0)) detect significant transit timing variations.

#### 5.2.7. TOI–1451 b

TOI–1451 b is a sub-Neptune  $(2.61^{+0.13}_{-0.10} R_{\oplus})$  orbiting a G-type star (5800  $\pm$  139 K) with an orbital period of ~16.5 days. We measure a  $5\sigma$  mass of  $15.2 \pm 2.8$   $M_{\oplus}$  (Figure [17](#page-39-0)), placing TOI–1451 b in a region of the mass–radius diagram where many of the sub-Neptunes in our sample are concentrated (see Figure 4). Our periodogram analysis of the RVs after subtraction of the transiting planet also reveals a strong peak at ∼90 days that is not seen in the activity indicators. We treat this signal with caution because multiple ∼90 day aliases were seen in the California Legacy Survey data (Rosenthal et al. [2021](#page-48-0)). A periodogram of the window function for all RV data reveals no strong peak near 90 days; however, when considering the window function of *only* HIRES data, multiple peaks at integer fractions of a year are seen, including at 90 days. Because of this, we consider the 90 day signal to be an alias in the HIRES data.

## 5.2.8. TOI–1472 b

TOI–1472 b is a Neptune analog  $(4.16_{-0.13}^{+0.16} R_{\oplus})$  orbiting a K-type star with a period of ∼6.4 days. Our RV fits constrain the mass to  $16.5 \pm 5.1$   $M_{\oplus}$  (Figure [18](#page-40-0)), giving it a mass also comparable to that of Neptune.

# 5.2.9. HD 156141 b (TOI–1742 b)

HD 156141 b (TOI–1742 b) is a sub-Neptune  $(2.37^{+0.06}_{-0.05} R_{\oplus})$ orbiting a G-type star (5814  $\pm$  137 K) with an orbital period of  $\sim$ 21.3 days. We obtain a nearly 5.5 $\sigma$  mass constraint of  $9.7 \pm 1.9$   $M_{\oplus}$  (Figure [19](#page-41-0)), placing this planet in a region of the mass–radius diagram where many of the sub-Neptunes in our

<span id="page-10-0"></span>sample are concentrated. Our RVs also constrain the orbit of the planet to a modest eccentricity of  $0.3 \pm 0.1$ , making this the second-highest-eccentricity small ( $R < 4R_{\oplus}$ ) planet in the TKS sample after TOI–1751 b. We also detect a significant amount of curvature in our RV time series that correlates with our activity indicators.

#### 5.2.10. TOI–1753 b

TOI-1753 b is a sub-Neptune  $(2.34^{+0.08}_{-0.06} R_{\oplus})$  orbiting a G-type star (5620 ± 134 K) with an orbital period of  $\sim$ 5.3 days. Our RV fits (Figure [20](#page-42-0)) constrain the mass of the transiting planet to  $16.6 \pm 5.2$   $M_{\oplus}$ . It falls in the relatively common population of sub-Neptunes, though its mass and radius, consistent with either a large H/He envelope or some water enrichment (see Figure [4](#page-9-0)). A more precise mass would be required to disambiguate different planetary compositions.

# 5.2.11. TOI–1775 b

TOI–1775 b is a Saturn-sized (8.05 $^{+0.14}_{-0.13}$   $R_{\oplus}$ ) planet orbiting a late G/early K-type star (5283  $\pm$  116 K) with an orbital period of ∼10.2 days. We are obtain over a 6 $\sigma$  (Figure [21](#page-43-0)) mass constraint of  $96 \pm 15$   $M_{\oplus}$ , giving TOI–1775 b a mass consistent with that of Saturn.

#### 5.2.12. HD 77946 b (TOI–1778 b)

HD 77946 b (TOI-1778 b) is a sub-Neptune  $(2.9^{+0.19}_{-0.14}R_{\oplus})$ planet orbiting an F-type star  $(6000 \pm 147 \text{ K})$  with an orbital period of ∼10.2 days. We obtain a nearly 6.5σ (Figure [22](#page-44-0)) mass constraint of  $10.9 \pm 1.7$   $M_{\oplus}$ , placing this planet in a region of the mass–radius diagram where many of the sub-Neptunes in our sample are concentrated (see Figure [4](#page-9-0)). Our RVs obtain a  $3\sigma$  constraint on the orbital eccentricity of  $0.21 \pm 0.07$ , also placing this among the highest-eccentricity small  $(R < 4R_{\oplus})$  planets in the TKS sample. We also find evidence of curvature in our RV time series that does not correlate with our activity indices, suggesting the potential presence of a long-period companion.

# 5.2.13. TOI–1798 b and c

TOI–1798 is a K-type star  $(5106 \pm 115 \text{ K})$  that hosts two transiting planets, the first being TOI–1798 b, a  $2.76_{-0.1}^{+0.13}R_{\oplus}$ sub-Neptune with a period of ∼8.0 days. RV fits (Figure [23](#page-45-0)) constrain the mass to  $6.5 \pm 2 M_{\oplus}$ .

The second planet is TOI–1798c, a  $1.4^{+0.06}_{-0.07}R_{\oplus}$  super-Earth with a period of ∼0.44 days, making this another ultra-shortperiod planet in the TKS sample. We were able to obtain an  $8\sigma$ mass of  $5.6 \pm 0.7$   $M_{\oplus}$ , giving this planet a high density of  $11.5 \pm 2.0$  g cm<sup>-3</sup>. Dai et al. ([2021a](#page-47-0)) compiled all known USPs with a mass at the time and observed a correlation between planet mass and disk surface density of solid material  $(10^{[Fe/H]}\times M_*)$ , suggesting that planet mass increases with the availability of planetary building blocks within 1 au. Estimating the surface density of solids available in the protoplanetary<br>disk as  $10^{\text{[Fe/H]}} \times M_*$ , we find TOI–1798c is placed comfortably within the population of USPs studied in Dai et al. ([2021a](#page-47-0)). TOI–1798 is the third system in the TKS sample to host an inner USP accompanied by a longer-period, sub-Neptune-mass companion.

#### 5.2.14. TOI–1823 b

TOI-1823 b is a sub-Saturn-sized  $(7.54^{+0.33}_{-0.25} R_{\oplus})$  planet orbiting a K-type star (4926  $\pm$  113 K) with an orbital period of  $\sim$ 38.8 days. We obtain an 8 $\sigma$  mass constraint of 67.4  $\pm$  8.2  $M_{\oplus}$ (Figure [24](#page-46-0)) with 32 RVs, giving it a density slightly higher than that of Saturn.

#### 5.3. Statistical Validation of TOIs

For the 30 systems that either did not reach  $\geq 3\sigma$  mass significance or did not have a clear and independent periodogram detection, we used TRICERATOPS (Giacalone et al. [2021](#page-47-0)) to calculate the false-positive and nearby false-positive probabilities (FPP/NFPP) and statistically validate these targets. Statistical validation attempts to rule out astrophysical false-positive scenarios such as eclipsing binaries. While statistical validation has been employed since the Kepler era (Borucki et al. [2012](#page-47-0); Crossfield et al. [2016](#page-47-0)) with packages such as BLENDER (Torres et al. [2004](#page-48-0)) and VESPA (Morton [2015](#page-48-0)), we opt to use TRICERATOPS since it is actively maintained and designed specifically for use on TESS candidates (Morton et al. [2023](#page-48-0)).

We began by pulling the TOI's short-cadence lightcurve from the MAST database using Lightkurve (Lightkurve Collaboration et al. [2018](#page-48-0)) and then use transit information from the TKS Systems Properties catalog (MacDougall et al. [2023](#page-48-0)) to phase fold the data. We bin the time series data down to 100 bins, in the interest of computational speed.<sup>64</sup> The apertures used in the SPOC pipeline are pulled and used as input to TRICERATOPS. High-resolution imaging data is also incorporated; however, TRICERATOPS can only run with a single HCI data set at a time, so we chose the imaging set with the best available contrast (Section [2.2](#page-2-0)).

When calculating the NFPP, we remove any NFPP candidates flagged by the TESS Input Catalog (TIC, Stassun et al. [2018](#page-48-0)) as "artifacts," to avoid biasing the NFPP value. If the TIC is missing stellar parameters for NFPP candidates, we use the candidate's Bp–Rp color to estimate  $T_{\text{eff}}$ ,  $M_*$ , and  $R_*$  from Pecaut & Mamajek ([2013](#page-48-0)).

To obtain an accurate estimate of the FPP and NFPP, we run TRICERATOPS 20 times per candidate planet and report the average and standard deviation in Table [3](#page-11-0). Typically, for a planet to be validated, it must satisfy  $FPP < 0.015$  and  $NFPP < 0.001$ , In this work, we compute the average and standard deviation of the FPP and NFPP from the 20 runs and require a candidate to clear those thresholds to  $3\sigma$  to be considered validated. In total, we graduate 17 candidates to validated planets, while six targets, TOI–1436.01, 1438.01, 1438.02, 1269.02, 1716.01, and 2114.01, did not reach the requisite threshold for validation and remain planetary candidates. That said, 1436.01 and 2114.01 have low enough FPP/NFPP to be considered likely planets (Giacalone et al. [2021](#page-47-0)). In all cases, the most likely scenario was a planet transiting the target star, followed by a planet transiting the primary star but with unresolved bound companions or unresolved background stars. Two targets, TOI–1279.01 and TOI–1742.01, have FPPs of 26% and 1.8%, respectively, which would place them in the likely planet category. However, given the significant masses obtained for both of these planets (4 $\sigma$  and 5 $\sigma$  constraints, respectively), we consider

<sup>64</sup> https://[triceratops.readthedocs.io](https://triceratops.readthedocs.io/en/latest/tutorials/example.html)/en/latest/tutorials/example.html

<span id="page-11-0"></span>

TOI	No. Sectors	Imaging Used		<b>NFPP</b>	Disposition
260.01	2	Keck/NIRC2 Br $\gamma$	$0.0003 \pm 0.0003$	$<$ 10 <sup><math>-4</math></sup>	<b>Confirmed Planet</b>
1184.01	26	Gemini/NIRI K	$0.0001 \pm 0.0001$	${<}10^{-4}$	Confirmed Planet
1279.01*	6	Palomar/PHARO Br $\gamma$	$0.2631 \pm 0.0855$	${<}10^{-4}$	<b>Confirmed Planet</b>
1472.01	3	Keck/NIRC2 Br $\gamma$	$0.0001 \pm 0.0001$	${<}10^{-4}$	<b>Confirmed Planet</b>
$1742.01*$	26	Keck/NIRC2 Br $\gamma$	$0.0184 \pm 0.0792$	${<}10^{-4}$	<b>Confirmed Planet</b>
1753.01	26	Palomar/PHARO Br $\gamma$	$0.0005 \pm 0.0002$	${<}10^{-4}$	<b>Confirmed Planet</b>
1798.01	5	Gemini/'Alopeke 832 nm	$0.0003 \pm 0.0001$	$<$ 10 <sup>-4</sup>	<b>Confirmed Planet</b>
1180.01	11	WIYN/NESSI 832 nm	$0.0035 \pm 0.0032$	$<\!\!10^{-4}$	Validated Planet
1244.01	26	Keck/NIRC2 K	${<}10^{-4}$	${<}10^{-4}$	<b>Validated Planet</b>
1249.01	$\overline{7}$	Gemini/'Alopeke 832 nm	$0.0083 \pm 0.0016$	$0.0007 \pm 0.0001$	Validated Planet
1269.01	29	Palomar/PHARO Br $\gamma$	$0.0013 \pm 0.0004$	${<}10^{-4}$	Validated Planet
1269.02	29	Palomar/PHARO Br $\gamma$	$0.0091 \pm 0.0012$	$0.0064 \pm 0.001$	Candidate
1438.01	26	Gemini/'Alopeke 832 nm	$0.0047 \pm 0.0027$	$0.0016 \pm 0.0014$	Candidate
1691.01	5	Gemini/'Alopeke 832 nm	$0.0001 \pm 0.0003$	$<\!\!10^{-4}$	Validated Planet
1716.01	3	Gemini/'Alopeke 832 nm	$0.0018 \pm 0.0002$	$0.001 \pm 0.0001$	Candidate
1723.01	3	Shane/ShARCS K-short	$0.0006 \pm 0.0001$	${<}10^{-4}$	Validated Planet
1758.01	$\overline{7}$	Gemini/'Alopeke 832 nm	$0.0002 \pm 0.0002$	$<$ 10 <sup>-4</sup>	Validated Planet
1794.01	5	Palomar/PHARO H-continuous	$0.0002 \pm 0.0001$	${<}10^{-4}$	<b>Validated Planet</b>
1799.01	$\mathbf{2}$	Palomar/PHARO Br $\gamma$	$0.001 \pm 0.0001$	$<$ 10 <sup>-4</sup>	<b>Validated Planet</b>
2128.01	3	Palomar/PHARO H-continuous	$<$ $10^{-4}\,$	$<$ 10 <sup>-4</sup>	Validated Planet
1174.01	$\tau$	Gemini/'Alopeke 832 nm	$0.0012 \pm 0.0004$	${<}10^{-4}$	Validated Planet
1347.02	27	Keck/NIRC2 K'	$0.0097 \pm 0.0007$	${<}10^{-4}$	Validated Planet
1436.01	6	Keck/NIRC2 Br $\gamma$	$0.063 \pm 0.002$	${<}10^{-4}$	<b>Likely Planet</b>
1438.02	26	Gemini/'Alopeke 832 nm	$0.0067 \pm 0.0017$	$0.0045 \pm 0.0014$	Candidate
1443.01	21	Gemini/NIRI K	$0.0003 \pm 0.0002$	$<\!\!10^{-4}$	Validated Planet
1467.01	3	Keck/NIRC2 K	$<$ 10 <sup><math>-4</math></sup>	$<$ 10 <sup>-4</sup>	Validated Planet
1669.01	8	Gemini/'Alopeke 832 nm	$0.0021 \pm 0.0004$	$<$ 10 <sup>-4</sup>	Validated Planet
1776.01	2	Palomar/PHARO Br $\gamma$	$0.0029 \pm 0.0003$	${<}10^{-4}$	<b>Validated Planet</b>
2088.01	28	WIYN/NESSI 832 nm	$0.003 \pm 0.004$	${<}10^{-4}$	<b>Validated Planet</b>
2114.01	$\overline{2}$	Palomar/PHARO Br $\gamma$	$0.019 \pm 0.022$	${<}10^{-4}$	<b>Likely Planet</b>

Table 3 TRICERATOPS Results

Notes. The break in the table separates TOIs with a  $\geqslant 3\sigma$  mass measurement (top), those with a  $1\sigma$ -3 $\sigma$  mass measurement (middle), and those with only an upper limit (bottom). Validated planets satisfy FPP  $\leq 0.015$  and NFPP  $\leq 0.001$ . Likely planets satisfy FPP  $\leq 0.5$  and NFPP  $\leq 0.001$ . \* Indicates a target that would be considered a likely planet but is still confirmed through a significant mass measurement.

them confirmed. Additionally, the low NFPP across all of our targets suggests that none of the planets we targeted are likely to be nearby false positives.

#### 6. Survey Performance

The overarching goal of TKS was to measure precise ( $5\sigma$ ) planet masses (Chontos et al. [2022](#page-47-0)). We use this section to assess the success of TKS as well as reflect on how well we were able to our reach observing quotas. We additionally discuss how effective we were at allocating resources to each individual science case. While detailed information on each science case (SC) can be found in Chontos et al.  $(2022)$  $(2022)$  $(2022)$ , we provide a brief description here: SC1 is dedicated to measuring bulk compositions, SC2 is dedicated to probing system architectures, SC3 is dedicated to measuring masses of possible future atmospheric targets, and SC4 is dedicated to measuring masses of planets orbiting evolved stars.

# 6.1. Observation and Cadence Goals

We intended to get a maximum of 60 RVs for the majority of targets in our survey. However, three notable exceptions include: those designated as interesting multi-planet systems

(100 RVs), those orbiting particularly active hosts (100 RVs), and those in the Distant Giant survey (15 RVs). Figure [6](#page-12-0) shows a breakdown of number of observations per target in the survey, and Figure [7](#page-13-0) (top) shows the number of observations broken down by science case.

In total, 62 stars received more observations than their stated goal. Of the 22 targets that did not meet their observing quota, 14 of the targets had surpassed a  $\geq 5\sigma$  mass and therefore no more observations were needed. At the extreme ends, there were 8 stars that missed their stated goal of observations by greater than 40 observations. TOI–1136 is an activity target, which was intended to get 100 observations. Instead, we opted to pool all the time allocated to the stellar activity case (Technical Outcomes B; TB) into this high-interest target, and so we got 157 total observations. Out of 86 TOIs, only 3 were dropped: $^{65}$  HD 63433 (TOI–1726, Section [4.2](#page-6-0)), TOI–2045, and TOI–2114. Other targets, HD 63935 (TOI–509) and HD 235088 (TOI–1430), received more than 100 observations, despite being initially designated 60, due to increased interest in the target: HD 63935 later presented a second transiting

<sup>65</sup> Targets were dropped for various reasons, including higher-than-expected rotational broadening and stellar rotation periods that coincided with planet orbital periods.

<span id="page-12-0"></span>

Figure 6. The number of observations acquired for each target by instrument. Vertical dashed lines are at the 15, 60, and 100 observation thresholds set for the several TKS science cases.

<span id="page-13-0"></span>

Figure 7. Top: The number of observations obtained broken down by science case. The vertical dashed lines represent the  $N_{\rm obs}$  targets set in Chontos et al. ([2022](#page-47-0)). Bottom: The median time difference between observations (cadence) for our targets broken down by science case.

planet and HD 235088 was a high-priority atmosphere target that displayed increased stellar activity.

To determine our attained cadence, we computed the time difference for every observation of a given target (eliminating differences less than 0.85 days to avoid intra-night cadence), and took the median time difference. In Figure 7 (bottom), we break down the observational cadence of targets by their SC. Note that many targets are represented in multiple science cases. We find that the Bulk Compositions, Atmospheres, and Evolved Hosts have median cadences between 3 and 4 days, while the System Architectures saw a spread of cadences. The third panel from the bottom of the lower plot in Figure 7 shows the lower RV cadence required by the Distant Giants program (∼30 days between observations), which is part of the Architectures SC. The distribution of cadences achieved for different science cases shows what is possible with a multiyear RV survey and may be used in simulating observations for future TKS-like programs.

# 6.2. Mass Significance Achieved

It is similarly straightforward to assess our performance in measuring mass significance. We measure at least a  $5\sigma$  mass



Figure 8. A histogram of the mass significance of each planet in our sample (blue). We highlight TOIs that were part of either the Distant Giants (SC2A) or Stellar Activity (TB) science cases, i.e., programs where the transiting planet was not of primary interest, in the light blue, hashed histogram.

for 58 planets. This increases to 77 planets for  $\geq 3\sigma$  masses; see Figure 8 for a histogram of mass significances of our survey. Breaking this down a bit further, if we filter by planets with masses below 30  $M_{\oplus}$ , then we have measured 30 planets to  $5\sigma$ mass or better and 44 planets to  $3\sigma$  mass. Through these numbers, given the overarching goal of TKS of measuring many precise planet masses, we believe our survey has been very successful.

We can go further by acknowledging that our survey was driven by some science cases that either were not interested in the transiting TOIs (Distant Giants, SC2A) or intentionally made it challenging to measure a precise mass (Stellar Activity, TB). In Figure 8, we highlight this group of targets and find that many of our "underperforming" targets are in either of these categories. For TOIs with radii <4  $R_{\oplus}$ , we were able to reach at least a  $3\sigma$  mass for 38 of them. To put this number in context, the baseline science requirement of the TESS Followup Program (TFOP) was to measure masses for 50 TESSdiscovered planets less than 4  $R_{\oplus}$  radii.<sup>66</sup> Our survey alone was nearly able to achieve this goal.

# 6.3. APF as Survey Support

The use of the Automated Planet Finder in our survey was intended to provide additional RVs to bright  $(V < 10)$  targets. To that end, we conducted complementary APF observations of 28 TOIs, which account for nearly half the RV measurements presented in this work. In order to assess APF's contribution to the goals of the survey, we select a subset of the 28 TOIs and reanalyze these systems excluding APF data. This subset includes massive targets for which a  $5\sigma$  mass was obtained without reaching the requisite 60 observations, targets that had their observations halted at 60 HIRES data points, and targets that received considerable attention from APF, making up  $\sim$ 2/ 3 of their data sets.

A straightforward test of APF's contribution to the survey is its effect on mass precision. In Figure [9](#page-14-0), we show the mass constraints for the TOIs reanalyzed without APF data. Masses measured from the full RV data set were generally larger than

<sup>66</sup> https://[tess.mit.edu](https://tess.mit.edu/followup/)/followup/

<span id="page-14-0"></span>

Figure 9. Comparisons of the mass constraints for a sample of systems observed with both HIRES and APF. Blue error bars denote the mass obtained with both data sets, while black errors bars indicate those without APF data. Generally, APF contributed only slightly to the mass significance for each system, with the exceptions of TOI–1710 b and HIP97166 c.

those obtained with only HIRES data, though the difference is often statistically insignificant. For two systems, TOI–1710 b and HIP97166 c, APF had a significant impact on the mass significance achieved. With TOI-1710 b, a  $3\sigma$  mass was only obtained due to APF RVs, while without APF the non-transiting HIP97166 c would not have been detected in a periodogram of HIRES-only data. In most cases, targets saw slight reductions in their mass error after the inclusion of APF RVs.

For the magnitude range of these stars (V∼8.5-11), APF is a comparably lower-precision instrument, with average per-datum uncertainties being ∼2 times greater than the RV errors attained with HIRES. This means that mass constraints will generally be dominated by HIRES data. However, APF's contribution should not be assessed on mass precision alone. The higher cadence of APF observations (Table [1](#page-3-0)) allowed for more uniform phase coverage, and for HD 332231 b, HIP 97166 b, and HD 156141 b, this resulted in significantly improved constraints on the eccentricity of these planets, with all three cases seeing a reduction in the eccentricity uncertainty by a factor of two.

TOI–1736, HD 146757, and HD 191939 are also systems that benefited greatly from the inclusion of APF data. In TOI– 1736's case, we were able to pick up the acceleration of the system due to an outer companion through the initial 60 observations allocated to the single-TOI system. Additional observations made with APF were able to extend the baseline and critically catch two passages of the companion through periastron (Akana Murphy et al. [2023](#page-47-0), Figure 25). APF provided continual cadence, which was not feasible with HIRES, resulting in strong constraints on both the orbital period and orbital eccentricity of TOI–1736 b's companion.

HD 146757 (Desai et al. [2024](#page-47-0)) benefited from APF data through the identification of an alias in the HIRES data. This case highlights APF's strength as a sanity check when searching for long-period, massive companions with limited primary survey instrument data.

Finally, APF observations of HD 191939 during March through 2020 May (22 observations), at a time when few other observatories were taking data due to COVID-19 shutdowns, were crucial for early detection of planet e, the 101 day Saturnmass planet (Lubin et al. [2022](#page-48-0)). Without these observations, the signal may not have been recognized until much later, and possibly at an aliased period. Subsequent APF data continued to fill the phase space of the companion at a cadence higher than would have been possible with HIRES alone. HD 146757 benefited from APF data through the identification of an alias in the HIRES data.

APF's contribution to the goals of TKS should not be understated. Our understanding of the architecture of exoplanet systems hinges on both a high degree of phase coverage to determine the eccentricity, as well as long baselines to detect companions. For surveys with limited time allocations, these criteria create opposing interests. Supplementing the observations of the primary survey instrument with those made with instruments with less demand not only helps to allocate more time to targets that require a greater density of higher-precision data, but also serves to extend the scientific productivity of previous generation instruments. Future TKS-like surveys should seriously consider using APF-like instruments to extend their observational baseline and obtain denser phase coverage for eccentric planets. Additionally, support instruments can also be used to conduct observations ahead of the primary survey to identify systems that may be more interesting to certain science cases (e.g., Distant Giants) than would be expected based on the presence of a transiting planet alone.

## 7. Expected versus Measured Masses

In Chontos et al. ([2022](#page-47-0)), the expected RV semi-amplitude for each planet in our sample was calculated using theoretical masses derived from mass–radius relations. The relations used were Weiss & Marcy ([2014;](#page-48-0) WM14) for planets with  $R < 4R_{\oplus}$ and Chen & Kipping ([2017](#page-47-0); CK17) for planets with  $4 \le R < 11.3 R_{\oplus}$ , both of which are empirically derived. Figure [10](#page-15-0) shows the masses we measure compared to their expected values, with the error in the predicted mass being propagated from the radius uncertainty. It can immediately be seen that, across the mass range sampled, there is significant scatter around the one-to-one mass comparison line. This scatter is not surprising: we should not expect simple powerlaw mass–radius relationships to be able to capture the compositional diversity of planets at a given radius.

<span id="page-15-0"></span>

Figure 10. Left: A comparison of RV-measured masses to those predicted by the empirical mass–radius relations Weiss & Marcy ([2014](#page-48-0)) and Chen & Kipping ([2017](#page-47-0)). Blue points highlight planets that achieved a  $5\sigma$  mass significance, while the unfilled point represent the full sample. Right: The mass residual (measured minus predicted) as a function of stellar and planetary parameters (from top): RV semi-amplitude, planetary insolation flux, stellar mass, and stellar metallicity. Pearson and Spearman correlation coefficients are given for both the  $5\sigma$  sample and the full sample (in brackets).

Deriving a new mass–radius relation from the TKS sample is beyond the scope of this paper: we are using a limited sample of planets and not using all the RV data available for these systems. But, here we examine how our measured masses compare to what was predicted by these commonly used relations. We restrict this discussion to planets below 11.3  $R_{\oplus}$ since, beyond this threshold, the theoretical mass was assumed to be a Jupiter mass.

#### 7.1. Weiss and Marcy 2014

In their analysis, WM14 report an rms scatter of 2.7  $M_{\oplus}$  and 4.3  $M_{\oplus}$  for planets with R <1.5  $R_{\oplus}$  and  $1.5 \le R < 4 R_{\oplus}$ , respectively, and  $\chi^2$  values that cannot be explained by mass measurement error alone. For the TKS sample, we find the rms scatter to be comparable: 1.2  $M_{\oplus}$  and 5.3  $M_{\oplus}$  for the same radius bins (noting that our sample contains far fewer planets with R <1.5  $R_{\oplus}$ ).

We then check if the WM14 M–R relation appropriately predicted the planet masses, i.e., if the mass residual is centered around zero. For planets with  $1.5 \le R < 4 R_{\oplus}$ , we find that a slight bias toward underpredicted masses of ∼10% while masses predicted from WM14 for the smallest planets (R  $\langle 1.5 R_{\oplus} \rangle$  are in agreement with what we measure from RVs. Exploring this further, we also recalculate mass predictions using the modified version of WM14 that only includes planet masses derived through RVs (as opposed to masses derived from RVs and transit timing variations) and find that this largely mitigates the difference. This suggests that the WM14 M–R relations, which were based primarily on Kepler planets, may still be relied on to approximate the masses of planets coming from TESS.

## 7.2. Chen and Kipping 2017

For planets with masses predicted from CK17  $(4 \leq$  $R < 11.3 R_{\oplus}$ , we report a higher rms scatter of 32  $M_{\oplus}$ , but

also find this M–R relation produces the largest mass discrepancies, with an average percent difference of 24%. Considering planets in this category had an average measurement uncertainty of 11%, this represents an appreciable difference. We also use the complete CK17 M–R relation to include planets below  $4 R_{\oplus}$  in order to check whether this underestimation persists for smaller worlds. This extends the Neptunian relation down to ~1.24  $R_{\oplus}$ , at which point it transitions to the Terran world relation. The predicted masses from these relations show a percentage difference comparable to the RV+TTV WM14 relation, with the exception that this underestimation in mass is more prevalent for the smallest planets. Below 1.5  $R_{\oplus}$ , the mean percentage difference increases to 28%, although we note that we have few planets in this radius space and the average measurement uncertainty for these planets is 37%.

# 7.3. Comparison to Non-TKS TESS Planets

We also perform a similar analysis for both M–R relations with the 169 TESS-discovered exoplanets with a mass measurement. $67$  For WM14, we find a similar rms scatter, and again a slight positive bias in the mass residual that disappears once the RV-only relation is used.

With CK17, the underestimation of mass persists for the greater TESS sample. Between 1.5 and 4  $R_{\oplus}$ , the percent difference nearly doubles from 13% in just the TKS sample to 20% for all TESS-discovered planets. Beyond 4  $R_{\oplus}$ , the difference is comparable to the TKS sample at 27%, suggesting that, on average, CK17 underestimates planet masses by around a quarter, and in extreme cases, by one half. However, we do note that, below R  $\sim$ 1.5  $R_{\oplus}$ , the mass residual reduces to 0.5  $M_{\oplus}$  for the entire TESS-discovered sample. In contrast to WM14, CK17 has fewer "Neptunian" worlds, with the bulk of their sample having radii larger than  $\sim 10 R_{\oplus}$ . Therefore, this

<sup>67</sup> From the NASA Exoplanet Archive as of Sept. 2023

Table 4 Masses and Derived Parameters

<span id="page-16-0"></span>

TOI	TKS ID	Mass $(M_{\oplus})$	Mass Provenance	Density $(g \text{ cm}^{-3})$	$\boldsymbol{a}$ (au)	Literature Mass $(M_{\oplus})$	Reference
260.01	HIP1532	$3.3^{+1.1}_{-1.0}$	Mass	$5.6^{+2.02}_{-2.83}$	$0.091_{-0.002}^{+0.002}$	$\cdots$	
					$0.093\substack{+0.002\\-0.002}$		
266.01	HIP8152	$8.9^{+1.7}_{-1.7}$	Mass	$2.97\substack{+0.81 \\ -1.08}$		$7.8^{+1.8}_{-1.8}$	Akana Murphy et al. (2023) <sup>*</sup>
266.02	HIP8152	$10.7^{+2.1}_{-2.2}$	Mass	$3.63_{-1.38}^{+1.02}$	$0.139_{-0.002}^{+0.003}$	9.4 <sup>+2.2</sup> 40.7 <sup>+3.7</sup> 40.7 <sup>+3.7</sup>	Akana Murphy et al. (2023) <sup>*</sup>
329.01	T000329	$40.4^{+3.8}_{-3.7}$	Mass	$1.84\substack{+0.41 \\ -0.55}$	$0.0638_{-0.0013}^{+0.0012}$		Chontos et al. $(2024)^*$
465.01	WASP156	$41.8^{+2.9}_{-2.9}$	Mass	$1.05_{-0.23}^{+0.18}$ $0.8_{-0.35}^{+0.29}$	$0.0458_{-0.0009}^{+0.0008}$	$40.7^{+3.2}_{-2.9}$	Demangeon et al. (2018)
469.01	42813	$5.9^{+2.0}_{-2.0}$	Mass		$0.11^{+0.002}_{-0.002}$	$5.8^{+2.4}_{-2.4}$	Akana Murphy et al. (2023) <sup>*</sup>
480.01	39688	$20.8^{+4.4}_{-4.2}$	Mass	$4.91^{+1.3}_{-1.65}$	$0.077^{+0.001}_{-0.001}$	$15.8^{+2.2}_{-2.1}$	Chontos et al. $(2024)^*$
509.01	63935	$10.5^{+1.6}_{-1.5}$	Mass	$2.32^{+0.5}_{-0.65}$	$0.083_{-0.002}^{+0.002}$	$10.8^{+1.8}_{-1.8}$	Scarsdale et al. $(2021)^*$
509.02	63935	$11.1^{+2.0}_{-1.9}$	Mass	$2.11\substack{+0.49 \\ -0.62}$	$0.148^{+0.003}_{-0.003}$	$11.1^{+2.4}_{-2.4}$	Scarsdale et al. $(2021)^*$
554.01	25463	$8.7^{+2.7}_{-2.6}$	Mass	$2.87\substack{+0.95 \\ -1.2}$	$0.076_{-0.001}^{+0.001}$	$8.5^{+3.1}_{-3.1}$	Akana Murphy et al. (2023) <sup>*</sup>
554.02	25463	$4.3^{+1.9}_{-1.9}$	Mass	$9.09_{-5.87}^{+4.32}$	$0.0443_{-0.0008}^{+0.0008}$	< 4.1	Akana Murphy et al. $(2023)^*$
561.01	T000561	$3.0_{-0.45}^{+0.45}$	Mass	$5.3^{+1.2}_{-1.51}$	$0.0104^{+0.0002}_{-0.0002}$	$2.24_{-0.2}^{+0.2}$	Brinkman et al. $(2023)^*$
561.02	T000561	$6.9^{+1.3}_{-1.2}$	Mass	$1.5^{+0.39}_{-0.51}$	$0.087^{+0.002}_{-0.002}$	$6.6^{+0.73}_{-0.73}$	Brinkman et al. $(2023)^*$
561.03	T000561	$11.0^{+1.8}_{-1.8}$	Mass	$2.97_{-0.91}^{+0.7}$	$0.156_{-0.003}^{+0.003}$	$12.15^{+1.1}_{-1.1}$	Brinkman et al. $(2023)^*$
561.04	T000561	$8.2^{+3.0}_{-3.2}$	Mass	.	$0.327_{-0.008}^{+0.008}$	$13.6^{+1.4}_{-1.4}$	Brinkman et al. $(2023)^*$
669.01	T000669	$10.0^{+1.4}_{-1.4}$	Mass	$3.17^{+0.73}_{-0.95}$	$0.048\substack{+0.001\\-0.001}$	$9.8^{+1.5}_{-1.5}$	Akana Murphy et al. (2023) <sup>*</sup>
849.01	T000849	$41.8^{+2.5}_{-2.4}$	Mass	$\underset{1.93^{+1.29}{5.6^{+0.76}_{-0.92}}}{5.6^{+0.76}_{-0.92}}$	$0.0155_{-0.0003}^{+0.0003}$	39.09 <sup>+2.66</sup> 6.0 <sup>+1.3</sup> 6.0 <sup>+1.3</sup>	Armstrong et al. (2020)
1136.01	T001136	$7.6^{+5.1}_{-5.0}$	Mass		$0.067^{+0.001}_{-0.001}$		Beard et al. $(2024)^*$
1136.02	T001136	$12.3^{+6.8}_{-7.0}$	Mass	$0.68^{+0.38}_{-0.43}$	$0.107\substack{+0.002\\-0.002}$	$8.35_{-1.6}^{+1.8}$	Beard et al. $(2024)^*$
1136.03	T001136	$<$ 24.0	Mass	$<\!8.0$	$0.14_{-0.002}^{+0.002}$	$6.07^{+1.09}_{-1.01}$	Beard et al. $(2024)^*$
1136.04	T001136	$<$ 32.0	Mass	<3.0	$0.175^{+0.003}_{-0.003}$	$9.7^{+3.9}_{-3.7}$	Beard et al. $(2024)^*$
1136.05	T001136	$<$ 31.0	Mass	$<$ 25.0	$0.0514^{+0.0009}_{-0.0009}$	$3.5^{+0.8}_{-0.7}$	Beard et al. $(2024)^*$
1136.06	T001136	$12.2^{+9.3}_{-9.6}$	Mass	$4.11_{-3.37}^{+3.18}$	$0.23_{-0.004}^{+0.004}$	$5.6^{+4.1}_{-3.2}$	Beard et al. $(2024)^*$
1173.01	T001173	$28.3^{+4.1}_{-4.0}$	Mass	$0.21_{-0.05}^{+0.04}$	$0.071^{+0.001}_{-0.001}$	$\cdots$	
1174.01	T001174	$<$ 32.0	Mass	< 11.0	$0.08^{+0.001}_{-0.001}$	.	
1180.01	T001180	$10.0^{+5.3}_{-5.4}$	Mass	$1.93_{-1.25}^{+1.07}$	$0.082^{+0.002}_{-0.002}$	.	.
1181.01	T001181	$377.0^{+17.0}_{-17.0}$	Mass	$0.49^{+0.08}_{-0.11}$	$0.0364^{+0.0006}_{-0.007}$	$375.0^{+44.0}_{-44.0}$	Kabáth et al. (2022)
1184.01	T001184	$6.8^{+2.3}_{-2.3}$	Mass	$2.63_{-1.24}^{+0.98}$	$0.056_{-0.001}^{+0.001}$		
1194.01	T001194	$120.0^{+5.6}_{-5.7}$	Mass	$0.99\substack{+0.16 \\ -0.21}$	$0.034^{+0.0006}_{-0.0006}$		
1244.01	T001244	$6.6^{+3.8}_{-3.7}$	Mass	$2.66_{-1.78}^{+1.52}$	$0.061_{-0.001}^{+0.001}$	$\cdots$	.
1246.01	T001246	$5.4^{+1.9}_{-1.9}$	Mass	$0.73^{+0.28}_{-0.33}$	$0.132^{+0.003}_{-0.003}$	$5.3^{+1.7}_{-1.7}$	Turtelboom et al. $(2022)^*$
1246.02	T001246	$8.0^{+1.2}_{-1.1}$	Mass	$1.84\substack{+0.42 \\ -0.54}$	$0.049^{+0.001}_{-0.001}$	$8.1^{+1.1}_{-1.1}$	Turtelboom et al. $(2022)^*$
1246.03	T001246	$9.1_{-1.2}^{+1.3}$	Mass	$3.39_{-1.02}^{+0.78}$	$0.061^{+0.001}_{-0.001}$	$8.8^{+1.2}_{-1.2}$	Turtelboom et al. $(2022)^*$
1246.04	T001246	$14.5^{+2.3}_{-2.3}$	Mass	$1.83\substack{+0.47 \\ -0.65}$	$0.212_{-0.004}^{+0.004}$	$14.8^{+2.2}_{-2.2}$	Turtelboom et al. $(2022)^*$
1246.05	T001246	$31.4^{+4.5}_{-4.3}$	$m_p$ sin( <i>i</i> )	$\ldots$	$0.389^{+0.008}_{-0.008}$	$25.6^{+3.6}_{-3.6}$	Turtelboom et al. (Tess-Keck Survey 2022)*
1247.01	135694	$6.1_{-1.8}^{+1.8}$	Mass	$2.05\substack{+0.68 \\ -0.8}$	$0.12_{-0.002}^{+0.002}$	$5.7^{+2.1}_{-2.1}$	Akana Murphy et al. (2023) <sup>*</sup>
1248.01	T001248	$27.4^{+3.6}_{-3.5}$	Mass	$0.47^{+0.1}_{-0.12}$	$0.051^{+0.001}_{-0.001}$	$\cdots$	.
1249.01	T001249	$11.9^{+5.2}_{-5.1}$	Mass	$1.84^{+0.85}_{-1.11}$	$0.109_{-0.002}^{+0.002}$	$\cdots$	$\cdots$
1255.01	HIP97166	$19.1^{+1.6}_{-1.5}$	Mass	$6.86\substack{+1.26 \\ -1.58}$	$0.091\substack{+0.002\\-0.002}$	$20.0^{+1.5}_{-1.5}$	MacDougall et al. $(2021)^*$
1255.02	HIP97166	$8.0^{+1.6}_{-1.6}$	$m_p$ sin(i)	$\cdots$	$0.124_{-0.002}^{+0.002}$	$9.9^{+1.8}_{-1.8}$	MacDougall et al. $(2021)^*$
1269.01	T001269	$6.4^{+3.0}_{-2.9}$	Mass	$2.53_{-1.42}^{+1.22}$	$0.0496^{+0.0009}_{-0.001}$	$\cdots$	
1269.02	T001269	$6.8^{+4.4}_{-4.3}$	Mass	$3.15_{-2.29}^{+2.05}$	$0.083^{+0.001}_{-0.002}$	$\cdots$	
1272.01	T001272	$27.0^{+2.7}_{-2.5}$	Mass	$2.09_{-0.59}^{+0.44}$	$0.0417\substack{+0.0007\\ -0.0007}$	$24.6^{+2.3}_{-2.3}$	MacDougall et al. $(2022)^*$
1272.02	T001272	$21.7^{+3.6}_{-3.6}$	$m_p$ sin( <i>i</i> )	$\cdots$	$0.079^{+0.001}_{-0.001}$	$26.7^{+3.1}_{-3.1}$	MacDougall et al. $(2022)^*$
1279.01	T001279	$10.6^{+2.5}_{-2.5}$	Mass	$3.05_{-1.14}^{+0.88}$	$0.085_{-0.002}^{+0.002}$	$\cdots$	$\ldots$
1288.01	T001288	$44.1^{+2.7}_{-2.6}$	Mass	$1.97^{+0.34}_{-0.44}$	$0.0374^{+0.0007}_{-0.0008}$	$42.0_{-3.0}^{+3.0}$	Knudstrup et al. (2023)
1288.02	T001288	$85.7^{+12.1}_{-11.7}$	$m_p$ sin( <i>i</i> )	$\cdots$	$1.07^{+0.03}_{-0.03}$	$84.0^{+7.0}_{-7.0}$	Knudstrup et al. (2023)
1294.01	T001294	$63.9^{+4.7}_{-4.5}$	Mass	$0.34_{-0.08}^{+0.06}$	$0.051_{-0.001}^{+0.001}$	$62.1_{-4.6}^{+4.8}$	Chontos et al. $(2024)^*$
1294.02	T001294	$147.0^{+15.0}_{-14.0}$	$m_p$ sin( <i>i</i> )	$\cdots$	$0.61_{-0.02}^{+0.02}$	$148.5^{+17.0}_{-15.8}$	Chontos et al. $(2024)^*$
1296.01	T001296	$95.3^{+4.7}_{-4.6}$	Mass	$0.19_{-0.04}^{+0.03}$	$0.051\substack{+0.001\\-0.001}$	$94.7^{+1.4}_{-1.4}$	Moutou et al. $(2021)$
1298.01	T001298	$99.0_{-6.0}^{+6.0}$	Mass	$0.6_{-0.13}^{+0.1}$	$0.057^{+0.001}_{-0.001}$	$113.0^{+10.0}_{-10.0}$	Moutou et al. $(2021)$
1339.01	191939	$9.6^{+0.9}_{-0.9}$	Mass	$1.34_{-0.33}^{+0.25}$	$0.079^{+0.002}_{-0.002}$	$10.0^{+0.7}_{-0.7}$	Orell-Miquel et al. $(2023)^*$
1339.02	191939	$6.1^{+1.3}_{-1.3}$	Mass	$1.03_{-0.34}^{+0.27}$	$0.171\substack{+0.003 \\ -0.003}$	$8.0\substack{+1.0\\-1.0}$	Orell-Miquel et al. $(2023)^*$
1339.03	191939	$2.2^{+1.3}_{-1.3}$	Mass	$0.4_{-0.26}^{+0.23}$	$0.209_{-0.004}^{+0.004}$	$2.8^{+0.6}_{-0.6}$	Orell-Miquel et al. $(2023)^*$
1339.04	191939	$116.0^{+5.0}_{-5.0}$	$m_p$ sin( <i>i</i> )	$\cdots$	$0.4\substack{+0.008\\-0.008}$	$112.2^{+4.0}_{-4.0}$	Orell-Miquel et al. $(2023)^*$
1339.05	191939	$\cdots$	$m_p$ sin( <i>i</i> )	$\cdots$	$\cdots$	$<$ 660.0	Orell-Miquel et al. $(2023)^*$
1339.06	191939	$\ldots$	$m_p$ sin( <i>i</i> )	$\cdots$	$\cdots$	<13.5	Orell-Miquel et al. $(2023)^*$

Table 4 (Continued)

TOI	TKS ID	Mass	Mass Provenance	Density	$\boldsymbol{a}$	Literature Mass	Reference
		$(M_{\oplus})$		$(g \text{ cm}^{-3})$	(au)	$(M_{\oplus})$	
1347.01	T001347	$12.2^{+1.6}_{-1.5}$	Mass	$11.17^{+2.46}_{-3.2}$	$0.0171^{\mathrm {+0.0003}}_{\mathrm {-0.0003}}$	$11.1^{+1.2}_{-1.2}$	Rubenzahl et al. $(2024)^*$
1347.02	T001347	< 9.0	Mass	< 11.0	$0.0546_{-0.001}^{+0.0009}$	${<}6.4$	Rubenzahl et al. $(2024)^*$
1386.01	T001386	$45.8^{+5.8}_{-5.7}$	Mass	$1.04_{-0.27}^{+0.21}$	$0.173^{+0.003}_{-0.003}$	$\ldots$	
1386.02	T001386	$75.2^{+15.3}_{-15.0}$	$m_p$ sin(i)	$\cdots$	$0.75_{-0.02}^{+0.03}$	$\cdots$	
1410.01	T001410	$12.5^{+1.1}_{-1.0}$	Mass	$2.3_{-0.74}^{+0.53}$	$0.0207^{+0.0004}_{-0.0004}$	$12.4^{+0.5}_{-0.5}$	Livingston et al. in review
1410.02	T001410	$27.0^{+3.3}_{-3.2}$	$m_p$ sin( <i>i</i> )	$\ldots$	$0.239_{-0.005}^{+0.004}$	$27.2^{+1.7}_{-1.7}$	Livingston et al. in review
1411.01	GJ9522A	$2.0^{+1.2}_{-1.1}$	Mass	$6.36_{-4.14}^{+3.71}$	$0.02122^{+0.0004}_{-0.0004}$	$\cdots$	
1422.01	T001422	$12.1^{+4.5}_{-4.5}$	Mass	$1.08^{+0.42}_{-0.5}$	$0.108^{+0.002}_{-0.002}$	$9.0^{+2.3}_{-2.0}$	Naponiello et al. (2022)
1430.01	235088	$4.2^{+2.4}_{-2.4}$	Mass	$2.58^{+1.47}_{-1.66}$	$0.071\substack{+0.001\\-0.001}$	$\cdots$ $\ldots$	
1436.01 1437.01	T001436 154840	$<\!8.0$ $10.4^{+2.6}_{-2.6}$	Mass Mass	< 12.0 $3.91^{+1.11}_{-1.4}$	$0.0164^{+0.0003}_{-0.0003}$ $0.139_{-0.003}^{+0.003}$	$9.6^{+3.9}_{-3.3}$	Pidhorodetska et al. in review*
1438.01	T001438	$10.9^{+5.9}_{-5.9}$	Mass	$2.47^{+1.4}_{-1.85}$	$0.056\substack{+0.001\\-0.001}$	$\cdots$	
1438.02	T001438	$5.6^{+5.0}_{-5.1}$	Mass	$1.7^{+1.51}_{-1.85}$	$0.084^{+0.002}_{-0.002}$	.	
1439.01	T001439	$38.5^{+5.7}_{-5.6}$	Mass	$2.76^{+0.64}_{-0.85}$	$0.192^{+0.004}_{-0.004}$	$38.2^{+5.6}_{-5.7}$	Chontos et al. $(2024)^*$
1443.01	T001443	$<$ 30.0	Mass	$<$ 14.0	$0.147^{+0.003}_{-0.003}$	$\cdots$	.
1444.01	T001444	$3.58_{-0.75}^{+0.76}$	Mass	$6.84^{+1.78}_{-2.27}$	$0.0116_{-0.0002}^{+0.0002}$	$3.87^{+0.71}_{-0.71}$	Dai et al. $(2021b)^*$
1444.02	T001444	$10.2^{+2.7}_{-2.9}$	$m_p$ sin( <i>i</i> )	$\cdots$	$0.123\substack{+0.002\\-0.002}$	$11.8^{+2.9}_{-2.9}$	Dai et al. $(2021b)^*$
1451.01	T001451	$15.2^{+2.8}_{-2.8}$	Mass	$4.67^{+1.18}_{-1.55}$	$0.127^{+0.003}_{-0.003}$	$\cdots$	
1456.01	332231	$74.1^{+5.6}_{-5.4}$	Mass	$0.55^{+0.1}_{-0.12}$	$0.145^{+0.003}_{-0.003}$	$77.6^{+6.7}_{-6.7}$	Dalba et al. $(2020)^*$
1467.01	T001467	< 9.0	Mass	< 11.0	$0.0495_{-0.0009}^{+0.0009}$	$\cdots$	
1471.01	12572	$8.3^{+1.7}_{-1.7}$	Mass	$0.81_{-0.25}^{+0.21}$	$0.145^{+0.003}_{-0.003}$	$8.4^{+2.0}_{-2.0}$	Akana Murphy et al. $(2023)^*$
1471.02	12572	${<}7.0$	Mass	< 1.0	$0.269\substack{+0.005 \\ -0.005}$	${<}7.9$	Akana Murphy et al. $(2023)^*$
1472.01	T001472	$16.5^{+5.1}_{-5.1}$	Mass	$1.24_{-0.51}^{+0.41}$	$0.065_{-0.001}^{+0.001}$	$\cdots$	
1473.01	6061	$10.0^{+2.5}_{-2.5}$	Mass	$3.82^{+1.1}_{-1.39}$	$0.06\substack{+0.001\\-0.001}$	$10.8^{+2.7}_{-2.7}$	Akana Murphy et al. $(2023)^*$
1601.01	T001601	$387.0^{+16.0}_{-17.0}$	Mass	$0.74^{+0.14}_{-0.18}$	$0.069\substack{+0.001\\-0.001}$	$388.0^{+10.0}_{-10.0}$	Chontos et al. $(2024)^*$
1611.01	207897	$14.8^{+1.5}_{-1.4}$	Mass	$6.29_{-1.78}^{+1.32}$	$0.117_{-0.002}^{+0.002}$	$14.4^{+1.6}_{-1.6}$	Heidari et al. (2022)
1669.01	T001669	< 13.0	Mass	< 6.0	$0.0376^{+0.0007}_{-0.0007}$		
1669.02	T001669	$191.0^{+25.0}_{-24.0}$	$m_p$ sin( <i>i</i> )	$\ldots$	$1.27^{+0.03}_{-0.03}$	$182.0^{+24.0}_{-24.0}$	Van Zandt et al. $(2023)^*$
1691.01	T001691	$14.6^{+5.4}_{-5.3}$	Mass	$1.76_{-0.8}^{+0.68}$	$0.126_{-0.003}^{+0.003}$	$\cdots$	
1694.01 1694.02	T001694 T001694	$31.3^{+2.7}_{-2.5}$ $297.0^{+16.0}_{-16.0}$	Mass	$1.13_{-0.26}^{+0.21}$ $\cdots$	$0.045\substack{+0.0008\\-0.0008}$ $0.1^{+0.02}_{-0.02}$	$26.1^{+1.1}_{-1.1}$ $334.0^{+16.0}_{-16.0}$	Van Zandt et al. $(2023)^*$ Van Zandt et al. $(2023)^*$
1710.01	T001710	$22.4^{+4.1}_{-4.0}$	$m_p$ sin(i) Mass	$0.87^{+0.2}_{-0.25}$	$0.166^{+0.003}_{-0.0013}$	.	
1716.01	237566	$4.8^{+2.0}_{-2.1}$	Mass	$1.09_{-0.59}^{+0.49}$	$0.081^{+0.002}_{-0.002}$		
1723.01	T001723	$10.4^{+5.2}_{-4.8}$	Mass	$1.6^{+0.79}_{-0.91}$	$0.114_{-0.002}^{+0.002}$		
1726.01	63433	$37.3^{+9.6}_{-9.6}$	Mass	$20.15_{-7.15}^{+5.89}$	$0.072^{+0.001}_{-0.001}$	$2.4^{+3.0}_{-1.8}$	Damasso et al. $(2023)$
1726.02	63433	$17.2_{-9.3}^{+10.3}$	Mass	$5.43_{-3.47}^{+3.13}$	$0.146_{-0.003}^{+0.003}$	$18.9^{+7.0}_{-6.9}$	Damasso et al. $(2023)$
1736.01	T001736	$12.3^{+1.5}_{-1.4}$	Mass	$2.09_{-0.56}^{+0.43}$	$0.073^{+0.001}_{-0.001}$	$11.9^{+1.6}_{-1.6}$	Akana Murphy et al. $(2023)^*$
1736.02	T001736	$2441.0^{+97.0}_{-99.0}$	$m_p$ sin( <i>i</i> )		$1.37^{+0.03}_{-0.03}$	$2477.0^{+118.0}_{-118.0}$	Akana Murphy et al. (2023) <sup>*</sup>
1742.01	156141	$9.7^{+1.9}_{-1.9}$	Mass	$4.0^{+0.98}_{-1.21}$	$0.154^{+0.003}_{-0.003}$	$\cdots$	
1751.01	146757	$19.5^{+3.5}_{-3.4}$	Mass	$4.13_{-1.26}^{+0.99}$	$0.215_{-0.004}^{+0.004}$	$20.9^{+4.5}_{-4.2}$	Desai et al. $(2024)^*$
1753.01	T001753	$16.6^{+5.2}_{-5.2}$	Mass	$5.95_{-2.49}^{+2.06}$	$0.059\substack{+0.001\\-0.001}$	$\cdots$	.
1758.01	T001758	$6.9^{+6.2}_{-6.3}$	Mass	$0.83_{-0.83}^{+0.75}$	$0.138_{-0.003}^{+0.003}$	$\ldots$	
1759.01	T001759	$<$ 15.0	Mass	$<$ 3.0	$0.115_{-0.002}^{+0.002}$	$6.8^{+2.0}_{-2.0}$	Martioli et al. (2022b)
1775.01	T001775	$96.0^{+15.0}_{-15.0}$	Mass	$1.01^{+0.22}_{-0.27}$	$0.08\substack{+0.002\\-0.002}$	$\cdots$	.
1776.01	95072	$1.4^{+1.1}_{-1.0}$	Mass	$4.13_{-3.79}^{+3.13}$	$0.0378_{-0.0008}^{+0.0007}$	$\cdots$	
1778.01	77946	$10.9^{+1.8}_{-1.7}$	Mass	$2.44^{+0.61}_{-0.83}$	$0.073^{+0.001}_{-0.01}$		
1794.01	T001794	$8.7^{+4.1}_{-4.0}$	Mass	$1.31_{-0.76}^{+0.62}$	$0.082\substack{+0.002\\-0.002}$	$\cdots$	
1797.01	93963	$18.4^{+3.2}_{-3.1}$	Mass	$3.31^{+0.8}_{-1.01}$	$0.04726^{\mathrm {+0.0008}}_{\mathrm {-0.0009}}$	$19.2^{+4.1}_{-4.1}$	Serrano et al. (2022)
1797.02	93963	$2.8^{+2.3}_{-2.3}$	Mass	$4.81^{+3.91}_{-4.36}$	$0.02061^{+0.0004}_{-0.0004}$	$7.8^{+3.2}_{-3.2}$	Serrano et al. (2022)
1798.01	T001798	$6.5^{+2.0}_{-2.0}$	Mass	$2.61^{+0.9}_{-1.11}$	$0.074^{+0.001}_{-0.001}$	$\cdots$	.
1798.02 1799.01	T001798 96735	$5.6^{+0.8}_{-0.7}$ $4.0^{+1.7}_{-1.8}$	Mass	$11.17^{+2.5}_{-3.39}$ $7.51_{-4.4}^{+3.45}$	$0.0107\substack{+0.0002\\-0.0002}$ $0.071_{-0.001}^{+0.001}$	$\cdots$ $\cdots$	.
1801.01	HIP57099	$<$ 16.0	Mass Mass	< 12.0	$0.075_{-0.001}^{+0.001}$	$\ldots$	. .
1807.01	HIP65469	$2.44\substack{+0.6 \\ -0.57}$	Mass	$3.98^{+1.18}_{-1.53}$	$0.0122\substack{+0.0002\\-0.0002}$	$2.57^{+0.5}_{-0.5}$	Nardiello et al. (2022)
1823.01	TIC142381532	$67.4^{+8.1}_{-8.3}$	Mass	$0.86_{-0.24}^{+0.18}$	$0.212_{-0.004}^{+0.004}$	$\ldots$	$\cdots$
1824.01	T001824	$16.0^{+2.6}_{-2.6}$	Mass	$4.24_{-1.21}^{+0.96}$	$0.151^{+0.003}_{-0.003}$	$18.5^{+3.2}_{-3.2}$	S. Lange et al. (2024, in preparation)*
1836.01	148193	$24.7^{+6.1}_{-5.9}$	Mass	$0.24\substack{+0.07 \\ -0.08}$	$0.157\substack{+0.003\\-0.003}$	$29.2^{+4.7}_{-4.8}$	Chontos et al. $(2024)^*$

TOI	TKS ID	<b>Mass</b> $(M_{\oplus})$	Mass Provenance	Density $(g \text{ cm}^{-3})$	$\boldsymbol{a}$ (au)	<b>Literature Mass</b> $(M_{\oplus})$	Reference
1836.02	148193	${<}8.0$	<b>Mass</b>	<3.0	$0.03089^{+0.0005}_{-0.0005}$	$\cdots$	$\cdots$
1842.01	T001842	$89.4^{+7.8}_{-7.8}$	Mass	$0.26_{-0.06}^{+0.05}$	$0.1^{+0.002}_{-0.002}$	$89.8^{+7.4}_{-7.4}$	Chontos et al. $(2024)^*$
1898.01	83342	$129.0^{+7.0}_{-8.0}$	Mass	$0.85^{+0.15}_{-0.2}$	$0.269_{-0.005}^{+0.005}$	$127.0^{+7.0}_{-7.0}$	Chontos et al. $(2024)^*$
2019.01	T002019	$34.4^{+4.4}_{-4.4}$	Mass	$1.25_{-0.37}^{+0.28}$	$0.128_{-0.003}^{+0.003}$	$34.7^{+4.2}_{-4.1}$	Chontos et al. $(2024)^*$
2045.01	T002045	$211.0^{+21.0}_{-20.0}$	Mass	$0.55_{-0.14}^{+0.11}$	$0.093_{-0.002}^{+0.002}$	$\cdots$	$\cdots$
2076.01	T002076	$16.1_{-8.2}^{+8.3}$	Mass	$4.1^{+2.14}_{-2.52}$	$0.088^{+0.002}_{-0.002}$	$\cdots$	$\cdots$
2076.02	T002076	$<$ 55.0	Mass	< 6.0	$0.142_{-0.003}^{+0.003}$	$\cdots$	$\cdots$
2076.03	T002076	${<}29.0$	<b>Mass</b>	$<$ 4.0	$0.199^{+0.004}_{-0.004}$	$\cdots$	$\cdots$
2088.01	T002088	<37.0	<b>Mass</b>	< 2.0	$0.472^{+0.009}_{-0.009}$	$\cdots$	$\cdots$
2114.01	T002114	$<$ 199.0	<b>Mass</b>	< 0.0	$0.075_{-0.001}^{+0.001}$	$\cdots$	$\cdots$
2128.01	155060	$4.5^{+2.0}_{-2.0}$	Mass	$2.68^{+1.24}_{-1.48}$	$0.127_{-0.003}^{+0.003}$	$\cdots$	$\cdots$
2145.01	HIP86040	$1767.0^{+71.0}_{-72.0}$	Mass	$5.06_{-1.11}^{+0.86}$	$0.111_{-0.002}^{+0.002}$	$1672.0^{+121.0}_{-118.0}$	Rodriguez et al. (2023)

Table 4 (Continued)

Note. \* Denotes a TKS publication. Upper limits are given as  $3\sigma$  upper limits. Note that masses marked as  $m_p \sin(i)$  lack an inclination measurement and should be treated as a lower limit.

sample may not be appropriately describing the lower-mass population of planets, which comprises the majority of worlds discovered to date.

#### 7.4. Sources of M–R Relation Bias and Implications

Empirical M–R relations tend to rely on planets that are well characterized in both radius and mass space. This is especially true for CK17, who instituted a  $3\sigma$  mass significance threshold for their sample, thereby not taking into account selection effects or survey bias. Owing to their larger K-amplitudes, a statistically significant mass measurement is often easier to obtain for more massive planets, and subsequent M–R relations based on these samples may be biased to higher masses (known as publication bias; see Burt et al. [2018;](#page-47-0) Montet [2018](#page-48-0)).

Comparing measured masses to their predicted values, we see the opposite trend in parts of the TKS sample. Restricting ourselves to planets that reached a  $5\sigma$  mass (blue points in Figure [5](#page-9-0)), we see that they are mainly the higher-mass planets for a given radius bin, showing that M–R relations based on "well-characterized" samples will indeed be biased toward higher mass. Weiss & Marcy ([2014](#page-48-0)) mitigated this bias by allowing for "negative" masses. It is possible that, if we were able to obtain greater mass significance on the remaining underperforming planets, the average mass residual for the well-characterized sample would move toward zero for the WM14 relation.

We may also expect the affect of publication bias to grow with decreasing RV semi-amplitude, as the noise floor of the instrument is approached. The right panel of Figure [5](#page-9-0) shows the mass residual of our sample as a function of select planetary and stellar parameters. Interestingly, the mass residual increases with larger RV semi-amplitude. HIRES is more than capable of measuring semi-amplitudes above  $5 \text{ m s}^{-1}$ , and planet masses in this range, primarily predicted with CK17, should not be affected by a stopping criteria based on mass significance. This implies that the CK17 M–R relation, at least for planets between 1.5 and  $\sim$ 11  $R_{\oplus}$ , is not necessarily affected by publication bias.

Another explanation is that the CK17 relations may simply be outdated. The number of planets below 100  $M_{\oplus}$  has doubled since CK17's study, providing a far more diverse data set from which to derive relationships between planet properties. The vast majority of CK17's sample were Jovian worlds with few planets between 1 and 10  $R_{\oplus}$ . A consequence of this, perhaps resulting from the  $3\sigma$  threshold criteria, is that the transition point between Terran and Neptunian worlds is placed at 2  $M_{\oplus}$ . In a recent reassessment of the M–R relation, Baron et al. ([2023](#page-47-0)) find that a transition nearer to ∼8  $M_{\oplus}$  more appropriately describes the current sample of exoplanets. A higher mass transition point results in a steeper M–R slope predicting higher masses for a given radius. While Baron et al. ([2023](#page-47-0)) also institute a quality cut on planet mass and radius, their analysis still benefits from the significant increase in well-characterized small planets.

Similarly, Edmondson et al. ([2023](#page-47-0)) also find discrepancies between measurements and CK17's predictions. When used to predict planetary radius, a scenario that will become increasingly important with the discovery of more non-transiting planets, CK17 tends to overestimate planet radius in agreement with our underestimate in mass.

M–R relations serve an important role in not only helping us to understand the population of exoplanets we are presented with today but also for highlighting targets that are particularly valuable for future follow-up. The transit spectroscopy metric (TSM, Kempton et al. [2018](#page-48-0)) estimates the expected signal-tonoise ratio of an atmosphere signal obtained with JWST, and it depends on the mass of the planet. Studies that use M–R relations to calculate the TSM can be crucial for bringing planets with high potential to be well characterized to the forefront (Hord et al. [2024](#page-47-0)). However, care should be employed with which M–R relations to use, and these empirical trends should not be treated as a substitute for obtaining actual mass measurements. In particular, TSMs calculated with masses predicted from CK17 may be higher than expected, especially for planets larger than  $\sim$ 1.5  $R_{\oplus}$ , producing exaggerated scale heights and overestimated TSM values.

#### 8. Conclusions

In this work, we present a uniform analysis of 9204 RV measurements taken as part of the TESS-Keck Survey.

<span id="page-19-0"></span>

Table 5





Ultimately, we obtain masses or mass upper limits on a total of 112 transiting and 15 non-transiting planets marking the culmination of 3 yr of observations from Keck/HIRES and the Automated Planet Finder. We conclude with our results and main findings:

- 1. The TESS-Keck Survey has generated a wealth of RV data that, with this paper, is now public (Section [2.1](#page-2-0)). This includes 5110 Keck/HIRES RVs and activity indicators in addition to 4534 RVs from the Automated Planet Finder covering an average baseline of nearly 4 yr (inclusive of archival data).
- 2. We uniformly fit all the data in our survey (Section [3](#page-4-0)) to produce the largest homogeneously analyzed RV data set of TESS-discovered planets to date. We also incorporate updated stellar and planetary parameters from the TKS System Properties Catalog (MacDougall et al. [2023](#page-48-0)).
- 3. We confirm or validate, through RV measurements and statistical validation, 32 new planets of varying size, adding to the many previously published systems in our survey (Section  $5.2 \& 5.3$  $5.2 \& 5.3$  $5.2 \& 5.3$ ). This ensures that future mass–radius relations derived with the TKS sample can incorporate mass upper limits on bona fide planets.
- 4. We assess the performance of our survey (Section [6](#page-11-0)) and find that we were able to achieve our observing goals (Chontos et al. [2022](#page-47-0)) for the majority of our targets. In total, we were able to achieve a  $\geqslant 5\sigma$  mass significance for 58 planets in our sample, increasing to 77 planets at  $\geq 3\sigma$  significance. We measured masses for 38 small planets below 4  $R_{\oplus}$  to 3 $\sigma$ , nearly achieving the TESS Follow-up Program's base science requirement of 50 small planets with masses in a single survey.
- 5. We gauged the benefit of using the Automated Planet Finder as survey support (Section [6.3](#page-13-0)). APF data played a crucial role in filling out the phase space for many

planets. This resulted in longer baselines for these systems, more precise constraints in the eccentricity of 3 TOIs, and the identification of an observing alias that would have been mistaken as a distant companion.

6. We compare our masses to those predicted using empirical mass–radius relations (Section [7](#page-14-0)). We find that the Weiss & Marcy ([2014](#page-48-0)) mass–radius relationship underpredicts the planet masses when compared to our data set. However, this discrepancy largely disappears when using the version of their MR relation that excludes masses derived from TTVs. Similarly, masses from Chen & Kipping  $(2017)$  $(2017)$  $(2017)$  seem to be underestimated by factors of  $1/4$  to  $1/2$  compared to our measurements, suggesting this mass–radius relation may be outdated and is not able to accurately describe the planet population we observe today.

The derived masses for our sample can be found in Table [4](#page-16-0) and the fitted parameters for each system in Table [6](#page-22-0). This work represents another set of uniformly analyzed planet masses for TESS-discovered planets, making this the second work to do so (Teske et al. [2021](#page-48-0)). Our survey contributes to the legacy of HIRES as a prolific instrument for planet detection and mass measurement as we enter an era of sub-m/s RV precision pushing down to smaller planet masses and longer orbital periods.

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Facilities: Keck:I (HIRES), APF (Levy), TESS, Exoplanet Archive

Software: RadVel (Fulton & Petigura [2018](#page-47-0)), Lightkurve (Lightkurve Collaboration et al. [2018](#page-48-0)), triceratops (Giacalone et al. [2021](#page-47-0))

# Appendix Radial Velocity Fits

In this appendix, we provide additional tables and figures mentioned earlier the text. Table [5](#page-19-0). gives a summary of the HRI observations taken for unpublished systems included in this work. Table [6](#page-22-0) provides the parameters that were fit for in RV model. Table [7](#page-32-0) provides the hyperparameter estimates for the relevant GP models. Figures [11](#page-33-0)–[24](#page-46-0) show the periodogram analysis of our RV, activity indicators, and window function for each of the systems we were able to confirm through mass measurements (Section [5.2](#page-7-0)). We also show the fits to the RV data points.

<span id="page-22-0"></span>



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Table 6

Note. The first line for each target displays fitted parameters for the system as a whole, such as RV offset for each instrument, linear and quadratic terms, etc. Values in parenthesis represent the 1 $\sigma$  errors for each BJD is given relative to the TESS standard  $T_0$ -2457000.

<span id="page-32-0"></span>

OF FILLE Parameters							
TOI (TKS ID)	$\eta_1$ (HIRES)	$\eta_1$ (APF)	$\eta_2$	$\eta_3$	$\eta_4$		
1759 (T001759)	$5.1^{+1.6}_{-1.1}$	$\cdots$	$178^{+140}_{-171}$	$36.63_{-0.53}^{+0.9}$	$0.487^{+0.049}_{-0.049}$		
260 (HIP 1532)	$3.12^{+0.72}_{-0.6}$	$\cdots$	$274_{-170}^{+170}$	$36.6^{+0.8}_{-0.44}$	$0.474_{-0.048}^{+0.051}$		
1174 (T001174)	$6.4^{+5.2}_{-6.3}$	$\cdots$	$145^{+100}_{-99}$	$11.49^{+0.35}_{-0.35}$	$0.501_{-0.05}^{+0.051}$		
1272 (T001272)	$0.032_{-0.018}^{+0.094}$	$\cdots$	$81^{+210}_{-42}$	$28.31_{-0.62}^{+0.63}$	$0.5_{-0.051}^{+0.052}$		
2076 (T002076)	$36.7^{+8.7}_{-7.4}$	$55^{+14}_{-11}$	$161_{-37}^{+77}$	$7.367^{+0.02}_{-0.022}$	$0.381^{+0.057}_{-0.049}$		
1473 (6061)	$3.7^{+2.5}_{-3.1}$	$9.6^{+2.2}_{-2.2}$	$27^{+30}_{-11}$	$13.87^{+0.81}_{-0.87}$	$0.488^{+0.052}_{-0.05}$		
1430 (235088)	$2.1^{+3.0}_{-1.9}$	$0.95_{-0.76}^{+4.8}$	$15.5^{+300}_{-7.7}$	$7.6^{+1.6}_{-2.3}$	$0.498_{-0.05}^{+0.051}$		
1726 (63433)	$15^{+16}_{-15}$	$44.7^{+7.4}_{-5.9}$	$91^{33}_{-40}$	$6.382^{+0.021}_{-0.016}$	$0.367^{+0.045}_{-0.05}$		
1807 (HIP 65469)	$33^{+5}_{-4}$	$\cdots$	$64.6^{+7.9}_{-7.6}$	$8.728_{-0.022}^{+0.021}$	$0.333_{-0.027}^{+0.028}$		
1347 (T001347)	$12.4_{-1.9}^{+2.9}$	$\cdots$	$111^{+76}_{-40}$	$16.47^{+0.12}_{-0.17}$	$0.460_{-0.05}^{+0.052}$		
1824 (T001824)	$5.16_{-0.58}^{+0.68}$	$4.68_{-0.68}^{+0.73}$	$10.3^{2.9}_{-2.3}$	$8.36_{-0.51}^{+0.47}$	$0.518_{-0.046}^{+0.048}$		

Table 7 GP Fitted Parameters

Note. Not all systems have APF data.  $\eta_1$ —GP amplitude for given instrument.  $\eta_2$ —GP exponential scaling.  $\eta_3$ —GP period.  $\eta_4$ —GP periodic scaling

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<span id="page-33-0"></span>

Figure 11. Top: periodogram analysis of all RVs for HIP 1532 b in our data set, first subtracting any offsets and the Gaussian process, then removing the best-fit planet signal (first three panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the rotation period of the star is in purple. The 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital plus Gaussian process model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-34-0"></span>

Figure 12. Top: periodogram analysis of all RVs for TOI–1173 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-35-0"></span>

Figure 13. Top: periodogram analysis of all RVs for TOI–1184 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-36-0"></span>

Figure 14. Top: periodogram analysis of all RVs for TOI–1194 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-37-0"></span>

Figure 15. Top: periodogram analysis of all RVs for TOI–1248 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-38-0"></span>

Figure 16. Top: periodogram analysis of all RVs for TOI-1279 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-39-0"></span>

Figure 17. Top: periodogram analysis of all RVs for TOI 1451 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. The vertical red line denotes the location of a known HIRES alias at ∼90 days. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-40-0"></span>

Figure 18. Top: periodogram analysis of all RVs for TOI–1472 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-41-0"></span>

Figure 19. Top: periodogram analysis of all RVs for HD 156141 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-42-0"></span>

Figure 20. Top: periodogram analysis of all RVs for TOI-1753 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-43-0"></span>

Figure 21. Top: periodogram analysis of all RVs for TOI-1775 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-44-0"></span>

Figure 22. Top: periodogram analysis of all RVs for HD 77946 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

<span id="page-45-0"></span>

Figure 23. Top: periodogram analysis of all RVs for TOI-1798 b and c in our data set, first subtracting any offsets, then removing the best-fit planet signals (first three panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The periods of planet b and c are shown with the blue and orange vertical lines respectively. The 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the two-planet fit are shown in the bottom panel. Bottom: RVs phasefolded to the orbital ephemeris of planet b and planet c.

<span id="page-46-0"></span>

Figure 24. Top: periodogram analysis of all RVs for TOI–1823 b in our data set, first subtracting any offsets, then removing the best-fit planet signal (first two panels). We also show a periodogram of the S-indices and the RV window function (last two panels). The planet period is given is shown with the blue vertical line, while the 1% false-alarm probability (FAP) is the horizontal blue dashed line. Middle: the best-fit Keplerian orbital model (blue line). We add in quadrature the jitter terms in Table [6](#page-22-0) with the measurement errors for each RV. Residuals to the one-planet fit are shown in the bottom panel. Bottom: RVs phase-folded to the orbital ephemeris of planet b.

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